# Supercritical self-avoiding walks are space-filling

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#### 1 Introduction

In 1953, Paul Flory [Flo53] proposed considering self-avoiding walks (*i.e.* visiting every vertex at most once) on a lattice as a model for polymer chains. Self-avoiding walks have turned out to be a very interesting object, leading to rich mathematical theories and challenging questions; see [MS93, BDGS12].

Denote by  $c_n$  the number of n-step self-avoiding walks on the hypercubic lattice ( $\mathbb{Z}^d$  with edges between nearest neighbors) started from some fixed vertex, e.g. the origin. Elementary bounds on  $c_n$  (for instance  $d^n \leq c_n \leq 2d(2d-1)^{n-1}$ ) guarantee that  $c_n$  grows exponentially fast. Since an (n+m)-step self-avoiding walk can be uniquely cut into an n-step self-avoiding walk and a parallel translation of an m-step self-avoiding walk, we infer that  $c_{n+m} \leq c_n c_m$ , from which it follows that there exists  $\mu = \mu(d) \in (0, +\infty)$  such that  $\mu := \lim_{n\to\infty} c_n^{1/n}$ . The positive real number  $\mu$  is called the connective constant of the lattice. The connective constant can be approximated in a number of ways, yet no closed formula exists in general. In the case of the hexagonal lattice, it was recently proved to be equal to  $\sqrt{2+\sqrt{2}}$  in [DCS12].

As it stands, the model has a strong combinatorial flavor. A more geometric variation was suggested by Lawler, Schramm and Werner [LSW04]. Let us describe their construction now — they were interested in the two-dimensional case, but here we will not make this restriction. Let  $\Omega$  be a simply connected domain in  $\mathbb{R}^d$  with two points a, b on the boundary. For  $\delta > 0$ , let  $\Omega_{\delta}$  be the largest connected component of  $\Omega \cap \delta \mathbb{Z}^d$  and let  $a_{\delta}$ ,  $b_{\delta}$  be the two sites of  $\Omega_{\delta}$  closest to a and b respectively. We think of  $(\Omega_{\delta}, a_{\delta}, b_{\delta})$  as being an approximation of  $(\Omega, a, b)$ . See figure 1.

Let x > 0. On  $(\Omega_{\delta}, a_{\delta}, b_{\delta})$ , define a probability measure on the finite set of self-avoiding walks in  $\Omega_{\delta}$  from  $a_{\delta}$  to  $b_{\delta}$  by the formula

$$\mathbb{P}_{(\Omega_{\delta}, a_{\delta}, b_{\delta}, x)}(\gamma) = \frac{x^{|\gamma|}}{Z_{(\Omega_{\delta}, a_{\delta}, b_{\delta})}(x)} \tag{1}$$

where  $|\gamma|$  is the length of  $\gamma$  (i.e. the number of edges), and  $Z_{(\Omega_{\delta},a_{\delta},b_{\delta})}(x)$  is a normalizing factor. A random curve  $\gamma_{\delta}$  with law  $\mathbb{P}_{(\Omega_{\delta},a_{\delta},b_{\delta},x)}$  is called the self-avoiding walk with parameter x in  $(\Omega_{\delta},a_{\delta},b_{\delta})$ . The sum  $Z_{(\Omega_{\delta},a_{\delta},b_{\delta})}(x) = \sum_{\gamma} x^{|\gamma|}$  (with the sum taken over all self-avoiding walks in  $\Omega_{\delta}$  from  $a_{\delta}$  to  $b_{\delta}$ ) is sometimes called the partition function (or generating function) of self-avoiding walks from  $a_{\delta}$  to  $b_{\delta}$  in the domain  $\Omega_{\delta}$ .

When the domain  $(\Omega, a, b)$  is fixed, we are interested in the scaling limit of the family  $(\gamma_{\delta})$ , *i.e.* its geometric behavior when  $\delta$  goes to 0. The qualitative behavior is expected

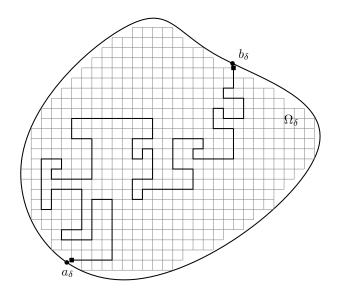


Figure 1: A domain  $\Omega$  with two points a and b on the boundary (circles) and the graph  $\Omega_{\delta}$ . The points  $a_{\delta}$  and  $b_{\delta}$  are depicted by squares. An example of a possible walk from  $a_{\delta}$  to  $b_{\delta}$  is presented. Note that there is a finite number of them.

to differ drastically depending on the value of x. A phase transition occurs at the value  $x_c = 1/\mu$ , where  $\mu$  is the connective constant:

When  $x < 1/\mu$ :  $\gamma_{\delta}$  converges to a deterministic curve corresponding to the geodesic between a and b in  $\Omega$  (assuming it is unique — otherwise some adaptations need to be done). When rescaled,  $\gamma_{\delta}$  should have Gaussian fluctuation of order  $\sqrt{\delta}$  around the geodesic. These results follow from the strong results of Ioffe [Iof98] on the unrestricted self-avoiding walk, though we are not aware of a reference for the details.

When  $x = 1/\mu$ :  $\gamma_{\delta}$  should converge to a random simple curve. In dimensions four and higher, the limit is believed to be a Brownian excursion from a to b in the domain  $\Omega$ . This is heuristically related to a number of rigorously proved results: in dimensions five and above to the work of Brydges and Spencer [BS85] and Hara and Slade [HS91, HS92] who showed that unrestricted self-avoiding walk converges to Brownian motion (see also the book [MS93]). Dimension four, the so-called *upper critical dimension*, is much harder, but recently some impressive results have been achieved using a supersymmetric renormalization group approach. These results are limited to continuous time weakly self-avoiding walk, see [BIS09, BS, BDS] and references within.

In dimension two, the scaling limit is conjectured to be the Schramm-Löwner Evolution of parameter 8/3, and in fact it was proved that this is true *if* the scaling limit exists as a continuous curve and is conformally invariant [LSW04].

Finally, dimension three remains a mystery, and there is no clear candidate for the scaling limit of self-avoiding walk.

When  $x > 1/\mu$ :  $\gamma_{\delta}$  is expected to become space-filling in the following sense: for any open set  $U \subset \Omega$ ,

$$\mathbb{P}_{(\Omega_{\delta}, a_{\delta}, b_{\delta}, x)} [\gamma_{\delta} \cap U = \varnothing] \to 0$$

when  $\delta$  goes to 0. On the one hand, let us mention that it is not clear in which sense (if any) ( $\gamma_{\delta}$ ) has a scaling limit when  $d \geq 3$ . On the other hand, the scaling limit is predicted [Smi06, Conjecture 3] to exist in dimension two (in the case of the hexagonal lattice at least). It should be the Schramm-Löwner Evolution of parameter 8, which is conformally invariant.

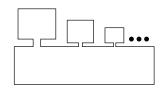
One cannot hope that  $\gamma_{\delta}$  would be space-filling in the strictest possible sense, namely that every vertex is visited. Nevertheless, one can quantify the size of the biggest hole not visited by the walk. The subject of this paper is the proof of a result which quantifies how  $\gamma_{\delta}$  becomes space filling. Here is a precise formulation.

**Theorem 1.** Let  $\mathbb{D}$  be the unit disk and let a and b be two points on its boundary. For every  $x > 1/\mu$ , there exist  $\xi = \xi(x) > 0$  and c = c(x) > 0 such that

 $\mathbb{P}_{(\mathbb{D}_{\delta}, a_{\delta}, b_{\delta}, x)}[there \ exists \ a \ component \ of \ \mathbb{D}_{\delta} \setminus \Gamma_{\delta}^{\xi} \ with \ cardinality \ larger \ than \ c \log(1/\delta)] \to 0$ 

when  $\delta \to 0$ , where  $\Gamma_{\delta}^{\xi}$  is the set of sites in  $\mathbb{D}_{\delta}$  at graph distance less than  $\xi$  from  $\gamma_{\delta}$ .

The theorem is stated for the unit disk to avoid various connectivity problems. Indeed, assume that at some given scale  $\delta$  our domain  $\Omega$  has a part which is connected by a "bridge" of width  $\delta$ . Then the graph  $\Omega_{\delta}$  will have a large part connected by a single edge, which does not leave the self-avoiding walk enough space to



enter and exit. Thus an analog of theorem 1 will not hold for this  $\Omega$ . It is not difficult to construct a single domain  $\Omega$  with such "mushrooms" in many scales. See the figure on the right. In order to solve this issue, one can start from an arbitrary domain and expand it microscopically. This gives rise to the following formulation:

**Theorem 2.** Let  $(\Omega, a, b)$  be a bounded domain with two points on the boundary. For every  $x > 1/\mu$ , there exist  $\xi = \xi(x) > 0$  and c = c(x) > 0 such that

$$\mathbb{P}_{((\Omega+B(\xi\delta))_{\delta},a_{\delta},b_{\delta},x)}\big[\exists \ a \ component \ of \ (\Omega+B(\xi\delta))_{\delta} \setminus \Gamma_{\delta}^{\xi} \ larger \ than \ c\log(1/\delta)\big] \to 0$$

when  $\delta \to 0$ , where  $\Gamma_{\delta}^{\xi}$  is the set of sites in  $(\Omega + B(\xi \delta))_{\delta}$  at graph distance less than  $\xi$  from  $\gamma_{\delta}$ .

Here  $\Omega + B(\xi \delta)$  is  $\Omega$  expanded by  $\xi \delta$  i.e.

$$\Omega + B(\xi \delta) = \{z : \operatorname{dist}(z, \Omega) < \xi \delta\}.$$

Since  $\xi$  depends only on x, and  $\delta \to 0$ , this is a microscopic expansion.

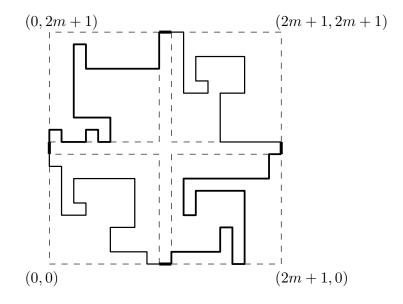


Figure 2: by concatenating four walks in squares of size m (plus four edges), one obtains a loop in the square of size 2m + 1 going through the middle of the sides.

The strategy of the proof is fairly natural. We first prove that in the supercritical phase, one can construct a lot (compared to their energy) of self-avoiding polygons in a prescribed box. Then, we show that the self-avoiding walk cannot leave holes that are too large, since adding polygons in the big holes to the self-avoiding walk would increase the entropy drastically while decreasing the energy in a reasonable way. In particular, a comparison energy/entropy shows that self-avoiding walks leaving big holes are unlikely. We present the proof **only in the case** d = 2, even though the reasoning carries over to all dimensions without difficulty (see Remark 8). One can also extend the result to other lattices with sufficient symmetry in a straightforward way (for instance to the hexagonal lattice).

## 2 Self-avoiding polygons in a square

In this section, we think of a walk as being indexed by (discrete) time t from 0 to n. For m > 0, let  $P_m$  be the set of self-avoiding polygons in  $[0, 2m + 1]^2$  that touch the middle of every face of the square: more formally, such that the edges [(m,0),(m+1,0)], [(2m+1,m),(2m+1,m+1)], [(m,2m+1),(m+1,2m+1)] and [(0,m),(0,m+1)] belong to the polygon, see figure 2. For x > 0, let  $Z_m(x)$  be the partition function (with parameter x) of  $P_m$ , i.e.

$$Z_m(x) = \sum_{\gamma \in P_m} x^{|\gamma|}.$$

**Proposition 3.** For  $x > 1/\mu$ , we have  $\limsup_{m\to\infty} Z_m(x) = \infty$ .

It is classical that the number of self-avoiding walks with certain constraints grows at the same exponential rate as the number of self-avoiding walks without constraints

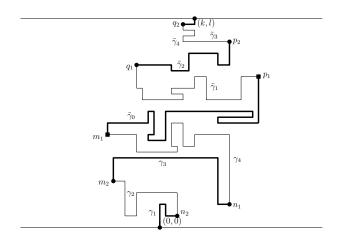


Figure 3: The decomposition of a bridge into walks. One can construct a squared walk in a rectangle by reflecting non-bold walks and then concatenate all the walks together.

(we will show it in our context in the proof of lemma 5 below). For instance, let x(v) and y(v) be the first and the second coordinates of the vertex v. The number  $b_n$  of self-avoiding bridges of length n, meaning self-avoiding walks  $\gamma$  of length n such that  $y(\gamma_0) = \min_{t \in [0,n]} y(\gamma_t)$  and  $y(\gamma_n) = \max_{t \in [0,n]} y(\gamma_t)$  satisfies

$$e^{-c\sqrt{n}}\mu^n \le b_n \le \mu^n \tag{2}$$

for every n [HW62] (see also [MS93] for a modern exposition). This result harnesses the following theorem on integer partitions which dates back to 1917.

**Theorem 4** (Hardy & Ramanujan [HR17]). For an integer  $A \ge 1$ , let  $P_D(A)$  denote the number of ways of writing  $A = A_1 + \cdots + A_k$  with  $A_1 > \cdots > A_k \ge 1$ , for any  $k \ge 1$ . Then

$$\log P_D(A) \sim \pi \sqrt{\frac{A}{3}}$$

as  $A \to \infty$ .

In the following, we need a class of walks with even more restrictive constraints. A squared walk (of span k) is a self-avoiding walk such that  $\gamma_0 = (0,0)$ ,  $\gamma_n = (k,k)$  and  $\gamma \in [0,k]^2$ .

**Lemma 5.** For c sufficiently large and n even, the number  $a_n$  of squared walks of length n satisfies

$$a_n \ge \mu^n e^{-c\sqrt{n}}.$$

**Proof step 1: rectangles.** Let  $\Lambda_n$  be the set of self-avoiding bridges of length n starting at the origin. Let  $\Sigma_n$  be the set of n-step self-avoiding walks for which there exists (k, l) such that  $\gamma_0 = (0, 0)$ ,  $\gamma_n = (k, l)$  and  $\gamma \in [0, k] \times [0, l]$ . We construct a map from  $\Lambda_n$  to  $\Sigma_n$ .

Fix  $\gamma \in \Lambda_n$  and denote by  $m_1$  the first time at which  $x(\gamma_{m_1}) = \min_{t \in [0,n]} x(\gamma_t)$ , see figure 3. Then, define  $n_1$  to be the first time at which  $x(\gamma_{n_1}) = \max_{t \in [0,m_1]} x(\gamma_t)$ . One can then define recursively  $m_k$ ,  $n_k$ , by the formulæ

$$m_{k} = \min\{r \le n_{k-1} : x(\gamma_{r}) = \min_{t \in [0, n_{k-1}]} x(\gamma_{t})\}$$

$$n_{k} = \min\{r \le m_{k} : x(\gamma_{r}) = \max_{t \in [0, m_{k}]} x(\gamma_{t})\}$$

We stop the recursion the first time  $m_k$  or  $n_k$  equals 0. For convenience, if the first time is  $n_k$ , we add a further step  $m_{k+1} = 0$ . We are then in the possession of a sequence of integers  $m_1 > n_1 > m_2 > \cdots > m_r \ge n_r \ge 0$  and a sequence of walks  $\gamma_{2r-1} = \gamma[n_1, m_1]$ ,  $\gamma_{2r-2} = \gamma[m_2, n_1], \ldots, \gamma_1 = \gamma[0, m_r]$ . Note that the width of the walks  $\gamma_i$  is strictly increasing (see figure 3 again).

Similarly, let  $p_1$  be the last time at which  $x(\gamma_{p_1}) = \max_{t \in [m_1, n]} x(\gamma_t)$  and  $q_1$  the last time at which  $x(\gamma_{q_1}) = \min_{t \in [p_1, n]} x(\gamma_t)$ . Then define recursively  $p_k$  and  $q_k$  by the following formula

$$p_{k} = \max\{r \ge q_{k-1} : x(\gamma_{r}) = \max_{t \in [q_{k-1}, n]} x(\gamma_{t})\}$$

$$q_{k} = \max\{r \le p_{k} : x(\gamma_{r}) = \min_{t \in [p_{k}, n]} x(\gamma_{t})\}$$

This procedure stops eventually and we obtain another sequence of walks  $\tilde{\gamma}_0 = \gamma[m_1, p_1]$ ,  $\tilde{\gamma}_1 = \gamma[p_1, q_1]$ , etc. This time, the width of the walks is strictly decreasing, see figure 3 one more time.

For a walk  $\omega$ , we set  $\sigma(\omega)$  to be its reflexion with respect to the vertical line passing through its starting point. Let  $f(\gamma)$  be the concatenation of  $\gamma_1$ ,  $\sigma(\gamma_2)$ ,  $\gamma_3$ ,..., $\sigma(\gamma_r)$ ,  $\tilde{\gamma}_0$ ,  $\sigma(\tilde{\gamma}_1)$ ,  $\tilde{\gamma}_2$  and so on. This walk is contained in the rectangle with corners being its endpoints so that f maps  $\Lambda_n$  on  $\Sigma_n$ .

In order to estimate the cardinality of  $\Sigma_n$ , we remark that each element of  $\Sigma_n$  has a limited number of possible pre-images under f. More precisely, the map which gives  $f(\gamma)$  and the widths of the walks  $(\gamma_i)$  and  $(\tilde{\gamma}_i)$  is one-to-one (the reverse procedure is easy to identify). The number of possible widths for  $\gamma_i$  and  $\tilde{\gamma}_i$  is the number of pairs of decreasing sequences partitioning an integer  $l \leq n$ . This number is bounded by  $e^{c\sqrt{n}}$  (Theorem 4). Therefore, the number of possible pre-images under f is bounded by  $e^{c\sqrt{n}}$ . The cardinality of  $\Sigma_n$  is thus larger than  $e^{-c\sqrt{n}}b_n \geq e^{-2c\sqrt{n}}\mu^n$ .

Step 2: Squares We have bounded from below the number of n-step self-avoiding walks 'contained in a rectangle'. We now extend this bound to the case of squares. There exist  $k, l \leq n$  such that the number of elements of  $\Sigma_n$  with (k, l) as an ending point is larger than  $e^{-2c\sqrt{n}}\mu^n/n^2$ . By taking two arbitrary walks of  $\Sigma_n$  ending at (k, l), one can construct a 2n-step self-avoiding walk with  $\gamma_0 = (0, 0)$  and  $\gamma_{2n} = (k + l, k + l)$  contained in  $[0, k + l]^2$  by reflecting orthogonally to  $e^{i\pi/4}\mathbb{R}$  the first walk, and then concatenating the two, see figure 4. We deduce that  $a_{2n} \geq \mu^{2n}e^{-4c\sqrt{n}}/n^4$ . This shows the lemma for n sufficiently large, and one can increase c if necessary to handle all even n.

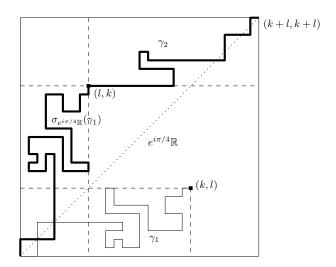


Figure 4: This figure depicts the passage of two walks in the rectangle  $[0, k] \times [0, l]$  to a walk in the square  $[0, k + l]^2$ .

**Proof of Proposition 3** Squared walks with length n were defined as walks between corners of some  $m \times m$  square, but m was not fixed. Fix now m to be such that the number of such walks is maximized (and then it is at least  $a_n/n$  where  $a_n$  is the total number of squared walks). It is interesting to remark that finding the maximal m as an explicit function of n, even asymptotically, seems difficult, probably no easier than the  $SLE_{8/3}$  conjecture. But we do not need to know its value. From any quadruplet  $(\gamma_1, \gamma_2, \gamma_3, \gamma_4)$  of such squared self-avoiding walks, one can construct a self-avoiding polygon of  $P_m$  as follows (see figure 2):

- translate  $\gamma_1$  and  $\gamma_3$  by (m+1,0) and (0,m+1) respectively,
- rotate  $\gamma_2$  and  $\gamma_4$  by an angle  $\pi/2$ , and then translate them by (m,0) and (2m+1,m+1) respectively,
- add the four edges [(m,0),(m+1,0)], [(2m+1,m),(2m+1,m+1)], [(m,2m+1),(m+1,2m+1)] and [(0,m),(0,m+1)].

Since each walk is contained in a square, one can easily check that we obtain a (4n+4)-step polygon in  $P_m$ . Using Lemma 5, we obtain

$$Z_m(x) \ge x^{4n+4} \left(\frac{a_n}{n}\right)^4 \ge \left(\frac{x^{n+1} \mu^n e^{-c\sqrt{n}}}{n}\right)^4.$$

When n goes to infinity, the right-hand side goes to infinity and the claim follows readily.

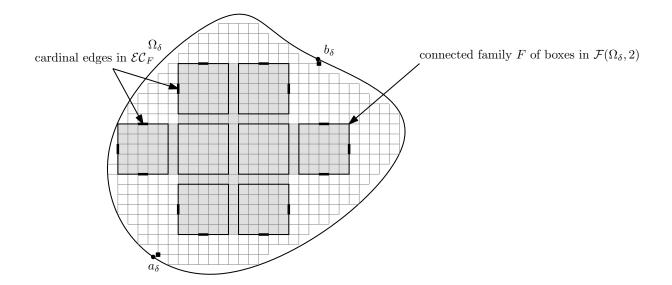


Figure 5: A discrete domain with a connected component of adjacent boxes of size 5 (m = 2). Edges of  $\mathscr{E}_F$  lie in the gray area.

#### 3 Proof of the main results

The strategy is the following. We first show that for some hole (namely it will be a connected union of boxes of some size m), the probability that the self-avoiding walk gets close to it without intersecting it can be estimated in terms of  $Z_m(x)$ . This claim is the core of the argument, and is presented in Proposition 7. Next, we show that choosing m large enough (or equivalently  $Z_m(x)$  large enough), the probability to avoid some connected union of k boxes decays exponentially fast in k, thus implying Theorem 1.

Let m > 0. A cardinal edge of a (square) box B of side length 2m + 1 is an edge of the lattice in the middle of one of the sides of B. For  $m \in \mathbb{N}$ , two boxes B and B' of side length 2m + 1 are said to be adjacent if they are disjoint and each has a cardinal edge, [xy] and [zt] respectively, such that  $x \sim z$ ,  $y \sim t$  (see figure 5). A family F of boxes is called connected if every two boxes can be connected by a path of adjacent boxes in F. To simplify the picture, we will assume that all our boxes have their lower left corner in  $(2m+2)\delta\mathbb{Z}^2$ .

Let us return to the issue of domain regularity discussed after Theorem 1. With the definitions above we can now explain that, in fact, our only requirement from the domain is that the family of all boxes in  $\Omega_{\delta}$  is connected. Let us state this formally.

**Theorem 6.** For every  $x > 1/\mu$ , there exists m = m(x) and c(x) > 0 such that for every domain  $\Omega$  and every  $\delta > 0$  such that the set of all boxes in  $\Omega_{\delta}$  is connected, one has, for every a and b in the boundary of  $\Omega$ , and every  $\lambda > 0$ ,

$$\mathbb{P}_{\Omega_{\delta}, a_{\delta}, b_{\delta}, x}(\exists \ a \ component \ of \ \Omega_{\delta} \setminus \Gamma_{\delta}^{6m} \ of \ size \ > \lambda) \leq \frac{C(x, \Omega)}{\delta^{2}} e^{-c(x)\lambda}.$$

The boxes in theorem 6 have side length 2m(x) + 1, and are aligned to have their lower-left corner in  $(2m+2)\delta\mathbb{Z}^2$ .  $\Gamma^{\xi}_{\delta}$  is as in Theorem 1.

**Proof of Theorem 1 given Theorem 6** Here our domain is  $\mathbb{D}$ . Clearly, the family of all boxes in  $\mathbb{D}_{\delta}$  is connected (it is an interval in every row and every column), hence Theorem 6 applies. Taking  $\lambda = C_1 \log(1/\delta)$  for  $C_1$  sufficiently large gives the result.

**Proof of Theorem 2 given Theorem 6** Again, all we have to show is that the family of boxes in  $(\Omega + B(\xi \delta))_{\delta}$  is connected for  $\xi$  sufficiently large. Taking  $\xi = 6m$  we see that every box in  $\Omega + B(6m\delta)$  can be connected to a box in  $\Omega$ , and any two boxes in  $\Omega$  can be connected by taking a path  $\gamma$  in  $\Omega$  between them (here is where we use that  $\Omega$  is connected) and checking that  $\gamma + B(6m\delta)$  contains a path of connected boxes.

Hence we need to prove Theorem 6. Let  $\delta > 0$ . Let  $\mathscr{F}(\Omega_{\delta}, m)$  be the set of connected families of boxes included in  $\Omega_{\delta}$ , again with the lower left corner in  $(2m+2)\delta\mathbb{Z}^2$ . For  $F \in \mathscr{F}(\Omega_{\delta}, m)$ , let  $\mathscr{V}_F$  be the set of vertices in boxes of F, and let  $\mathscr{E}_F$  be the set of edges with both end-points in  $\mathscr{V}_F$ . For two subsets A and B of the vertices of  $\Omega_{\delta}$  define the box distance between them as the size of the smallest set of connected boxes containing one box in A and one box in B, minus 1. The boxes do not have to be different, but then the distance is 0 — if no such connected set exists, then the distance is  $\infty$ . Denote it by boxdist(A, B).

**Proposition 7.** Let  $(\Omega, a, b)$  be a domain with two points on the boundary. Fix  $\delta > 0$  and  $m \in \mathbb{N}$  and assume the collection of all boxes is connected. Then there exists  $C(x, m) < \infty$  such that for every  $F \in \mathscr{F}(\Omega_{\delta}, m)$ ,

$$\mathbb{P}_{\delta,x}(\text{boxdist}(\gamma_{\delta}, \mathscr{V}_F) = 1) \leq C(x,m) Z_m(x)^{-|F|}.$$

Recall that  $\mathbb{P}_{\delta,x}$  is our measure on self-avoiding walks with parameter x, on the discretized domain  $\Omega_{\delta}$ .

**Proof** For  $F \in \mathcal{F}(\Omega_{\delta}, m)$ , let  $\mathscr{E}\mathscr{C}_F$  be the set of external cardinal edges of F i.e. all cardinal edges in boxes of F which have neighbors outside of F. Let  $S_F$  be the set of self-avoiding polygons included in  $\mathscr{E}_F$  visiting all the edges in  $\mathscr{E}\mathscr{C}_F$ . Let  $Z_F(x)$  be the partition function of polygons in  $S_F$ . We have:

Claim. For  $F \in \mathscr{F}(\Omega_{\delta}, m)$ ,  $Z_F(x) \geq Z_m(x)^{|F|}$ .

**Proof of the claim.** We prove the result by induction on the cardinality of  $F \in \mathcal{F}(\Omega_{\delta}, \xi)$ . If the cardinality of F is 1,  $Z_F(x) = Z_m(x)$  by definition. Consider  $F_0 \in \mathcal{F}(\Omega_{\delta}, \xi)$  and assume the statement true for every  $F \in \mathcal{F}(\Omega_{\delta}, \xi)$  with  $|F| < |F_0|$ . There exists a box B in  $F_0$  such that  $F_0 \setminus \{B\}$  is still connected. Therefore, for every pair  $(\gamma, \gamma') \in S_{\{B\}} \times S_{F_0 \setminus \{B\}}$ , one can associate a polygon in  $S_F$  in a one-to-one fashion. Indeed, B is adjacent to a box  $B' \in F_0 \setminus \{B\}$  so that one of the four cardinal edges (called [ab]) of B is adjacent to a cardinal edge [cd] of B'. Note that [cd] belongs to  $\mathscr{EC}_{F_0 \setminus \{B\}}$ . Then, by changing the edges [cd] and [ab] of  $\gamma$  and  $\gamma'$  into the edges [ac] and [bd], one obtains a polygon in  $S_F$ . Furthermore, the construction is one-to-one and we deduce

$$Z_{F_0}(x) \ge Z_{F_0 \setminus \{B\}}(x) Z_B(x) \ge Z_m^{|F_0 \setminus \{B\}|} Z_m(x) = Z_m(x)^{|F_0|}.$$

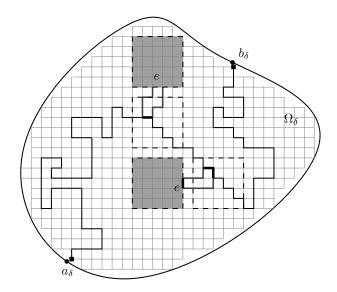


Figure 6: Two examples of the polygon  $\ell$  (the small N- and L-shaped loops). It overlaps the curve in one edge exactly, except in the second case, where we have no choice but overlapping the walk on two edges.

Consider the set  $\Theta_F$  of walks not intersecting F yet reaching to a neighboring box. Let e be a cardinal edge of F in distance 1 from a box visited by  $\gamma$ . For each  $\gamma \in \Theta_F$ , consider a self-avoiding polygon  $\ell = \ell(\gamma)$  ( $\ell$  standing for "link") satisfying the following three properties:

- it contains e and is included in  $(\Omega_{\delta} \setminus \mathscr{E}_F) \cup \{e\}$ ,
- it intersects  $\gamma$  at one or two adjacent edges only,
- it has length smaller than 10m + 10 (for simplicity, we will bound the length by 100m).

One can easily check that such a polygon always exists, see figure 6.

Now, consider the map that associates to  $(\gamma_1, \gamma_2) \in \Theta_F \times S_F$  the symmetric difference  $\gamma = f(\gamma_1, \gamma_2)$  of  $\gamma_1$ ,  $\ell(\gamma_1)$  and  $\gamma_2$  (symmetric difference here meaning as sets of edges). Note that the object that we obtain is a walk from  $a_\delta$  to  $b_\delta$  in  $\Omega_\delta$ , which can be verified to be self-avoiding by noting that each vertex has degree 0 or 2 and that the set is connected. Further, its length is equal to  $|\gamma_1| + |\ell(\gamma_1)| + |\gamma_2| - 4$  or  $|\gamma_1| + |\ell(\gamma_1)| + |\gamma_2| - 6$  (this being due to the fact that  $\gamma_1$  and  $\ell(\gamma_1)$  intersect at one or two adjacent edges — each intersection reduces the length by 2 edges). Now, given a path  $\gamma$  there is a limited number of ways it may be written as  $f(\gamma_1, \gamma_2)$ . Indeed, e can be located by the only two paths that exit F, and that gives  $\gamma_2$ . Given e,  $\ell$  has only a limited number of possibilities, say  $100m^2$ , and

once one knows  $\ell$  this gives  $\gamma_1$ . We can now write

$$Z_{\Theta_{F}}(x) \cdot Z_{F}(x) = \left(\sum_{\gamma_{1} \in \Theta_{F}} x^{|\gamma_{1}|}\right) \left(\sum_{\gamma_{2} \in S_{F}} x^{|\gamma_{2}|}\right)$$

$$\leq \max(1, x^{-100m}) \sum_{\gamma_{1} \in \Theta_{F}, \gamma_{2} \in S_{F}} x^{|\gamma_{1}| + |\ell(\gamma_{1})| + |\gamma_{2}|}$$

$$\leq \max(1, x^{-100m}) \max(x^{4}, x^{6}) \sum_{\gamma_{1} \in \Theta_{F}, \gamma_{2} \in S_{F}} x^{|f(\gamma_{1}, \gamma_{2})|}$$

$$\leq 100m^{2} \max(x^{6}, x^{-100m+4}) \sum_{\gamma \in f(\Theta_{F} \times S_{F})} x^{|\gamma|}$$

$$\leq 100m^{2} \max(x^{6}, x^{-100m+4}) Z_{(\Omega_{\delta}, a_{\delta}, b_{\delta})}(x),$$

where in the first inequality we used the fact that  $\ell(\gamma_1)$  has length smaller than 100m, in the second the fact that  $|f(\gamma_1, \gamma_2)|$  equals  $|\gamma_1| + |\ell(\gamma_1)| + |\gamma_2| - 4$  or  $|\gamma_1| + |\ell(\gamma_1)| + |\gamma_2| - 6$ , and in the third the fact that f is at most  $100m^2$ -to-one. Using the claim, the previous inequality implies

$$\mathbb{P}_{\delta,x}(\operatorname{boxdist}(\gamma_{\delta}, \mathscr{V}_{F}) \leq 1) = \frac{Z_{\Theta_{F}}(x)}{Z_{(\Omega_{\delta}, a_{\delta}, b_{\delta})}(x)} \leq \frac{C(x, m)}{Z_{F}(x)} \leq \frac{C(x, m)}{Z_{m}(x)^{|F|}}.$$

**Proof of Theorem 6 in dimension 2.** Let  $x > 1/\mu$  and let  $(\Omega, a, b)$  be a domain with two points on the boundary. Let  $A_n$  be the number of connected subsets of  $\mathbb{Z}^2$  containing 0. It is well known that  $\overline{\lim} \sqrt[n]{A_n}$  is finite (see e.g. Theorem 4.20 in [Gri99]). Let therefore  $\lambda = \lambda(2)$  satisfy  $A_n \leq \lambda^n$  for all n. We now apply Proposition 3 and get some m = m(x, 2) such that  $Z_m(x) > 2\lambda$ .

Let  $\delta > 0$  and consider the event  $\mathscr{A}(s)$  that there exists a connected set S of cardinality s at distance larger than 6m of  $\gamma_{\delta}$ . Every box intersecting S must be disjoint from  $\gamma$ , so there must exist a connected family of at least  $s/(2m+1)^2$  boxes of size 2m+1 covering S. We may assume this family is maximal among families covering S and not intersecting  $\gamma_{\delta}$ . Since the family of boxes is maximal, and because of the condition of the theorem that the family of all boxes is connected, the box-distance between the union of boxes and  $\gamma_{\delta}$  is 1. Proposition 7 implies

$$\mathbb{P}_{(\Omega_{\delta}, a_{\delta}, b_{\delta}, x)}[\mathscr{A}(s)] \leq \sum_{F \in \mathscr{F}(\Omega_{\delta}, \xi) : |F| \geq s/(2m+1)^2} C(x, m) \left[ Z_m(x) \right]^{-|F|}.$$

By the definition of  $\lambda$ , the number of families of connected boxes of size K in  $\mathscr{F}(\Omega_{\delta}, \xi)$  is bounded by  $(C(\Omega)/\delta^2)\lambda^K$  (since up to translation they are connected subsets of a normalized square lattice), where  $C(\Omega) = C(\Omega, x, m)$  depends on the area of  $\Omega$ , x and m. Therefore, for c > 0,

$$\mathbb{P}_{(\Omega_{\delta}, a_{\delta}, b_{\delta}, x)} \left[ \mathscr{A}(s) \right] \leq C(x, m) \frac{C(\Omega)}{\delta^{2}} \sum_{i \geq s/(2m+1)^{2}} \left( \frac{\lambda}{Z_{m}(x)} \right)^{i}$$

$$\leq \frac{C(x, m, \Omega)}{\delta^{2}} 2^{-s/(2m+1)^{2}}$$

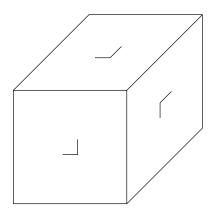


Figure 7: Cardinal edges in three dimensions.

and the theorem follows.

Remark 8. Let us briefly describe what needs to be changed in higher dimensions. The notion of cardinal edge must be extended: in the box  $[0, 2m+1]^d$ , cardinal edges for the face  $[0, 2m+1]^{d-1} \times \{0\}$  are all the edges joining vertices in  $\{m, m+1\}^{d-1} \times \{0\}$  of the form

$$\left[\left(\underbrace{m+1,\ldots,m+1}_{i-1 \text{ terms}},\underbrace{m,\ldots,m}_{d-i \text{ terms}},0\right), \left(\underbrace{m+1,\ldots,m+1}_{i \text{ terms}},\underbrace{m,\ldots,m}_{d-i-1 \text{ terms}},0\right)\right]$$

for  $1 \le i \le d-1$ . See figure 7 for an example in 3 dimensions. We only consider part of the edges joining vertices in  $\{m, m+1\}^{d-1} \times \{0\}$  because we want to contain all these edges, without forming a loop. Similarly, cardinal edges can be defined for every face. It can be shown that the number of polygons included in some box  $[0, 2m+1]^d$  and visiting all the cardinal edges grows exponentially at the same rate as the number of self-avoiding walks. The proofs then apply mutatis mutandis.

## 4 Questions

The supercritical phase exhibits an interesting behavior. We know that the curve becomes space-filling, yet we have very little additional information. For instance, a natural question is to study the length of the curve. It is not difficult to show that the length is of order  $1/\delta^2$ , yet a sharper result would be interesting:

Problem 9. For  $x > 1/\mu$ , show that there exists  $\theta(x) > 0$  such that for every  $\varepsilon > 0$  and every  $(\Omega, a, b)$ ,

$$\mathbb{P}_{(\Omega_{\delta}, a_{\delta}, b_{\delta}, x)}(||\gamma_{\delta}| - \theta(x) \cdot |\Omega_{\delta}||) > \varepsilon |\Omega_{\delta}|) \longrightarrow 0 \text{ when } \delta \to 0.$$

The quantity  $\theta(x)$  would thus be an averaged density of the walk. Note that the existence of  $\theta(x)$  seems natural since the space-filling curve should look fairly similar in different portions of the space.

Another challenge is to try to say something nontrivial about the critical phase. Recently, the uniformly chosen self-avoiding walk on  $\mathbb{Z}^d$  was proved to be sub-ballistic [DCH12]. A natural question would be to prove that it is *not* space-filling.

Problem 10. When  $x = 1/\mu$ , show that the sequence  $(\gamma_{\delta})$  does not become space-filling.

Finally, we recall the conjecture made in [Smi06] concerning the two-dimensional limit in the supercritical phase.

Conjecture 11 (Smirnov). Let  $(\Omega, a, b)$  be a simply connected domain of  $\mathbb{C}$  and consider approximations by the hexagonal lattice. The law of  $(\gamma_{\delta})$  converges to the chordal Schramm-Löwner Evolution in  $(\Omega, a, b)$ 

- with parameter 8/3 if  $x = 1/\mu$ ,
- with parameter 8 if  $x > 1/\mu$ .

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