Effective-One-Body approach to the Two-Body Problem in General Relativity

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Renewed importance of 2-body problem

- Gravitational wave (GW) signal emitted by binary black hole coalescences : a prime target for LIGO/Virgo/GEO
- GW signal emitted by binary neutron stars : target for advanced LIGO….

BUT

- Breakdown of analytical approach in such strong-field situations ? expansion parameter during coalescence ! ? $x \sim \frac{v^2}{c^2} \sim \mathcal{O}(1)$
- Give up analytical approach, and use only Numerical Relativity ?

Binary black hole coalescence

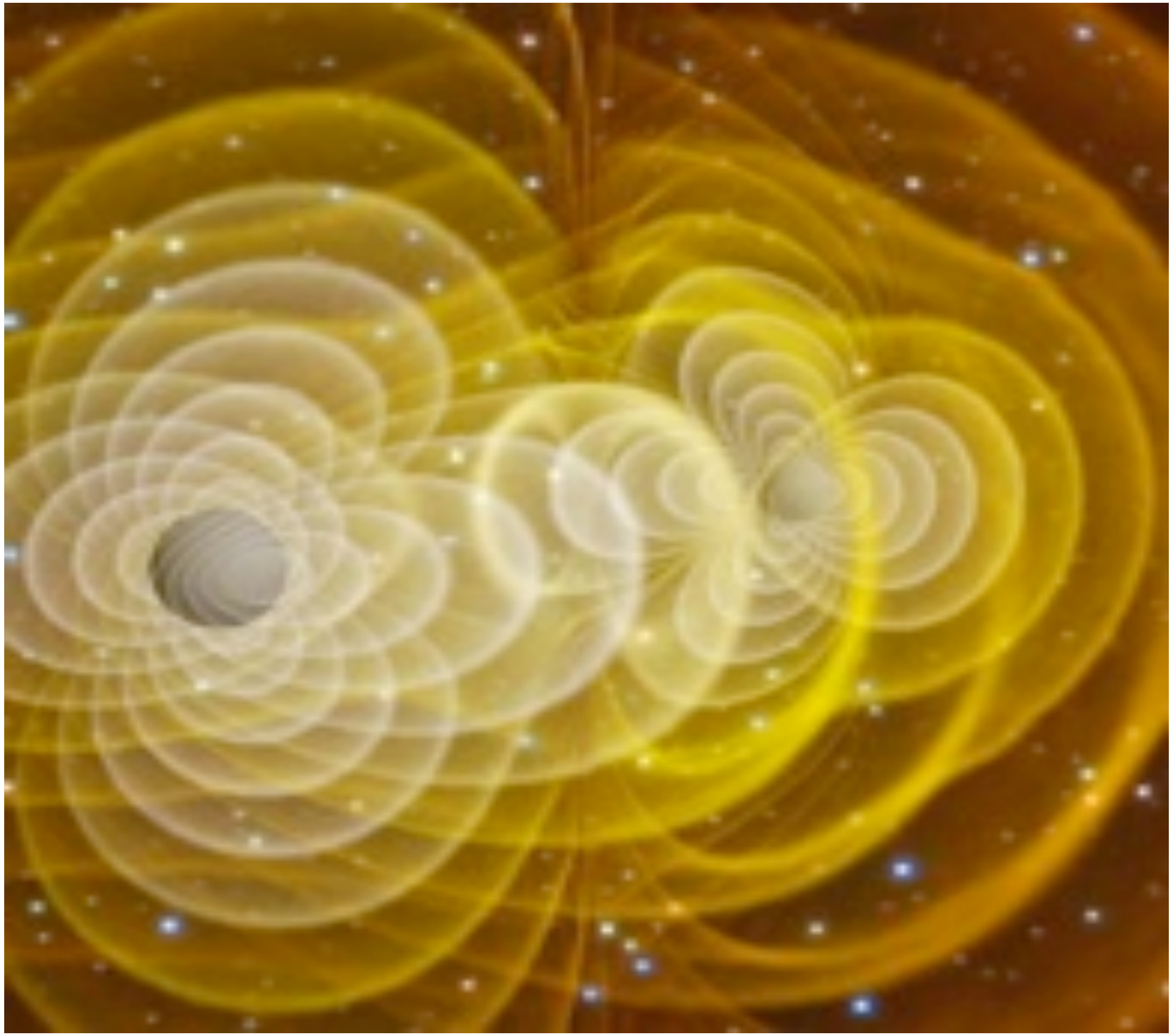
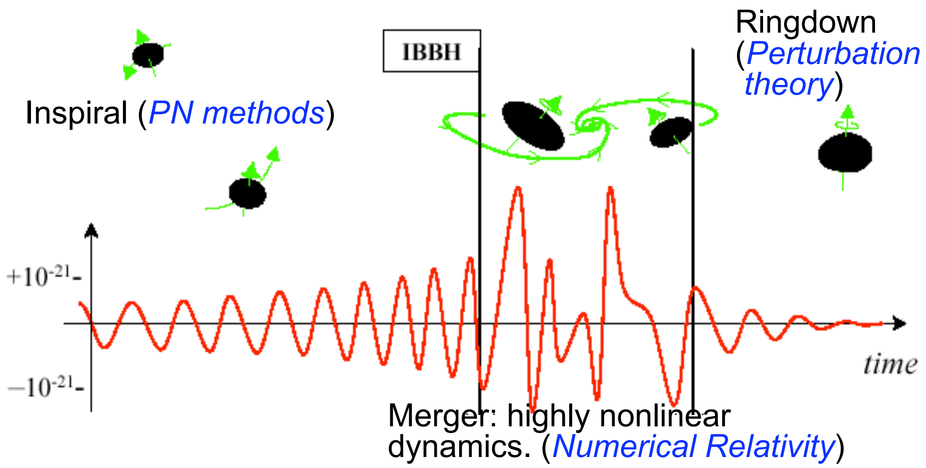


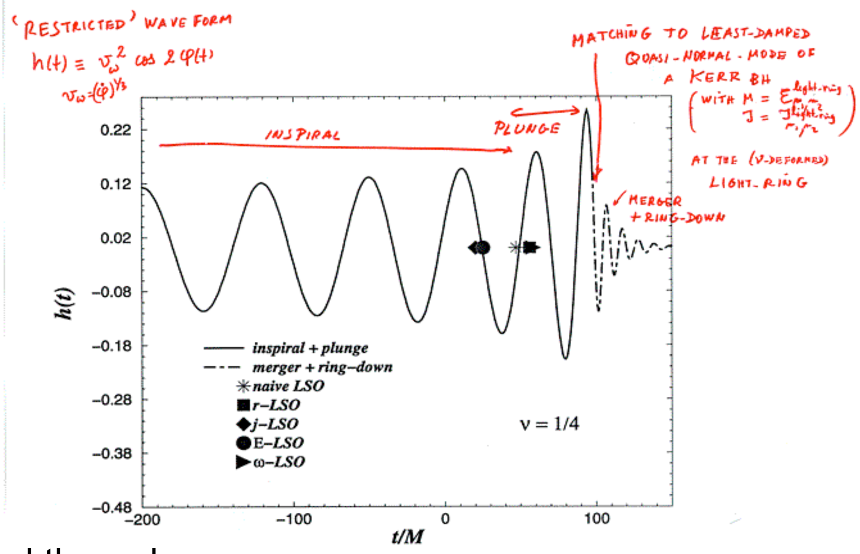
Image: NASA/GSFC

Templates for GWs from BBH coalescence

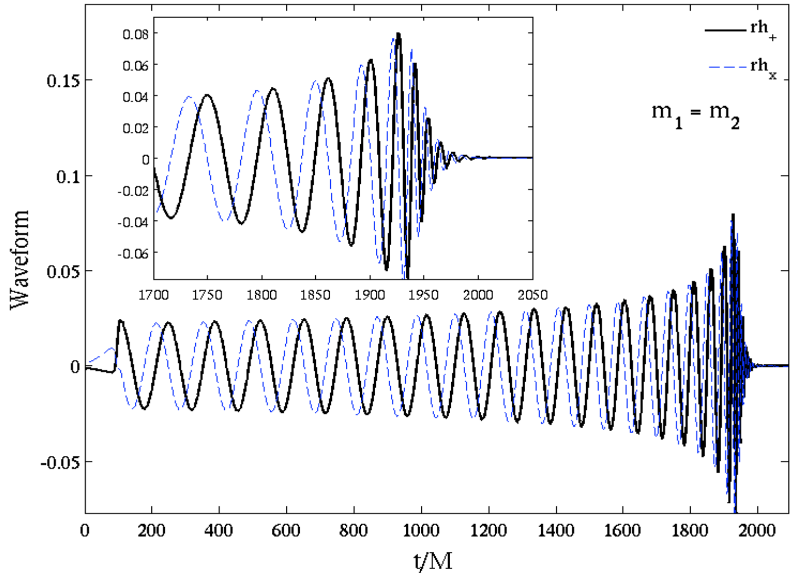
(Brady, Craighton, Thorne 1998)



(Buonanno & Damour 2000)



Numerical Relativity, the 2005 breakthrough:
Pretorius, Campanelli et al., Baker et al. …



An improved analytical approach

EFFECTIVE ONE BODY (EOB)

approach to the two-body problem

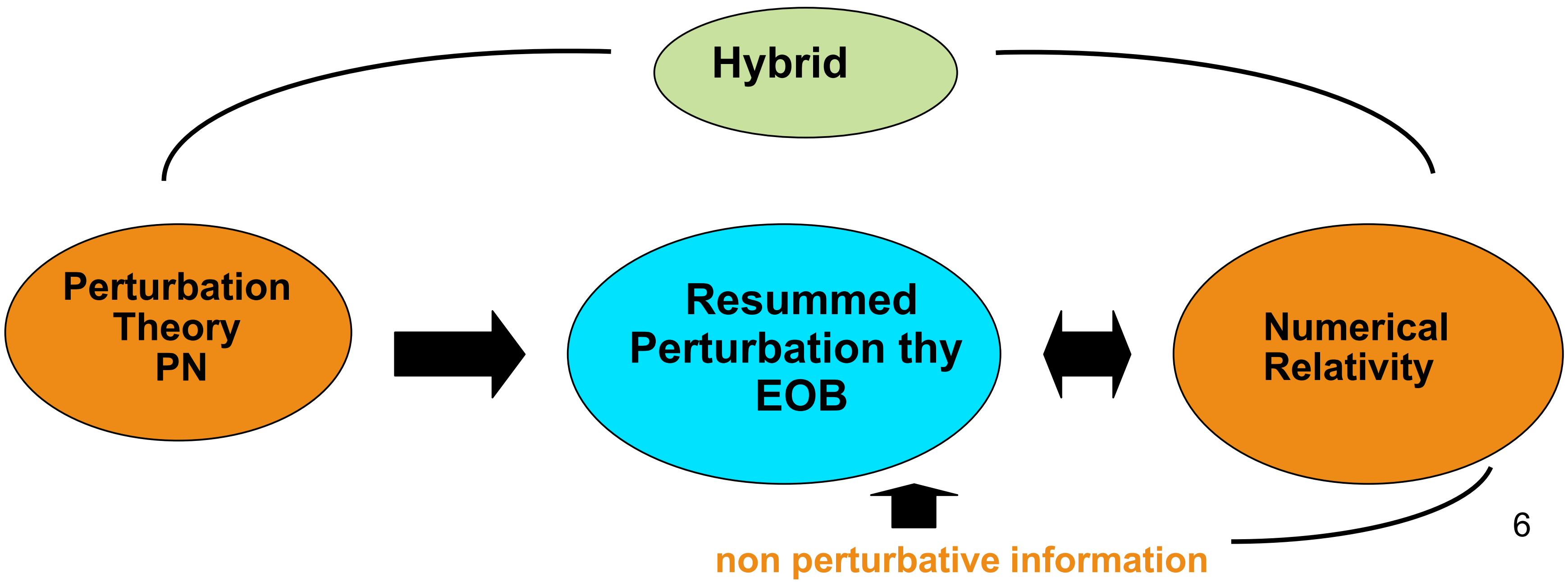
Buonanno,Damour 99	(2 PN Hamiltonian)
Buonanno,Damour 00	(Rad.Reac. full waveform)
Damour, Jaranowski,Schäfer 00	(3 PN Hamiltonian)
Damour, 01	(spin)
Damour, Nagar 07, Damour, Iyer, Nagar 08	(factorized waveform)

Importance of an analytical formalism

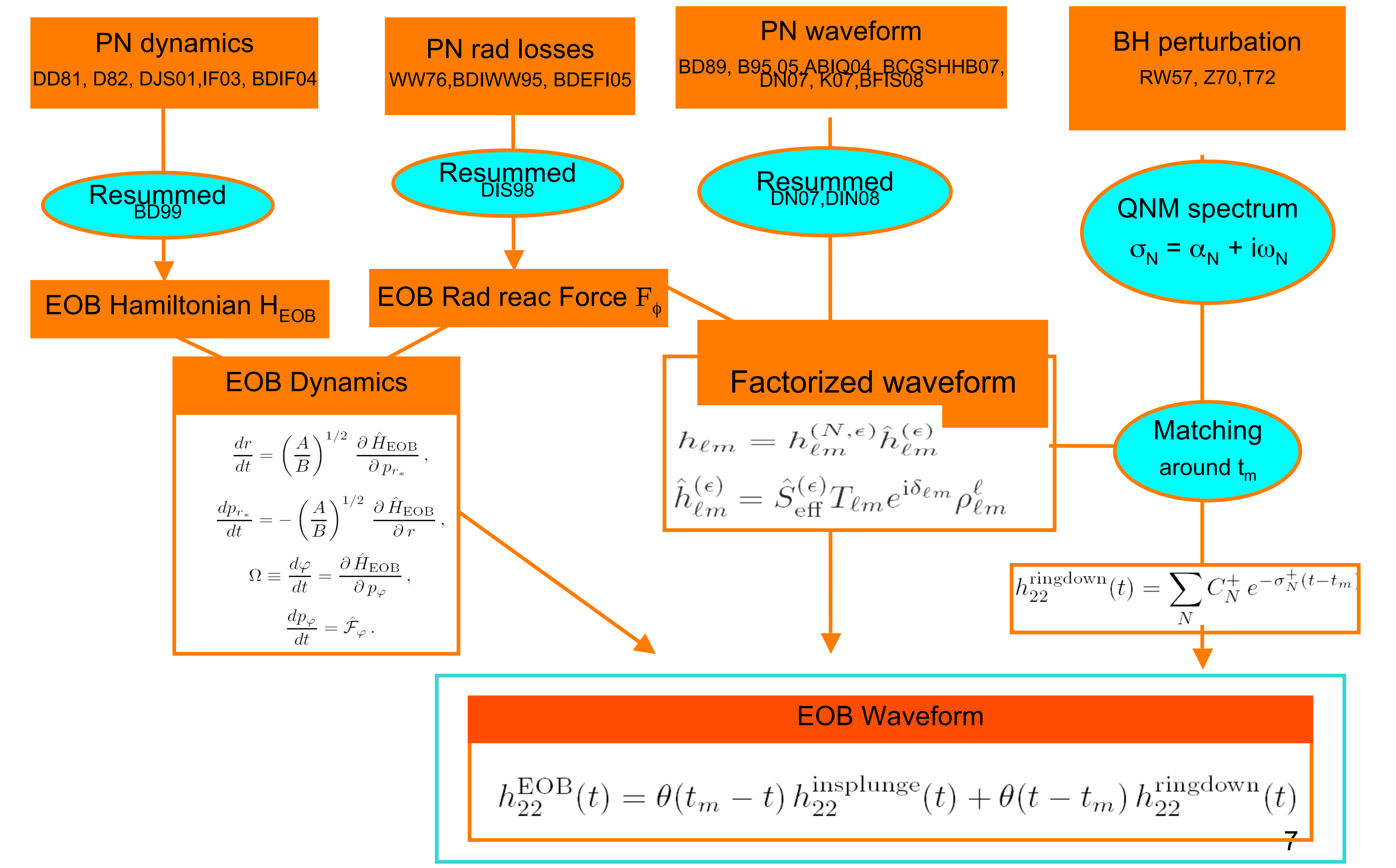
- **Theoretical:** physical understanding of the coalescence process, especially in complicated situations (arbitrary spins)

- **Practical:** need many thousands of accurate GW templates for detection & data analysis;
  need some "analytical" representation of waveform templates as $f(m_1,m_2,\mathbf{S}_1,\mathbf{S}_2)$

- Solution: **synergy between analytical & numerical relativity**



Structure of EOB formalism



Historical roots of EOB

- H_{EOB} : QED positronium states [Brezin, Itzykson, Zinn-Justin 1970]
  "Quantum" Hamiltonian $H(I_a)$ [Damour-Schäfer 1988]

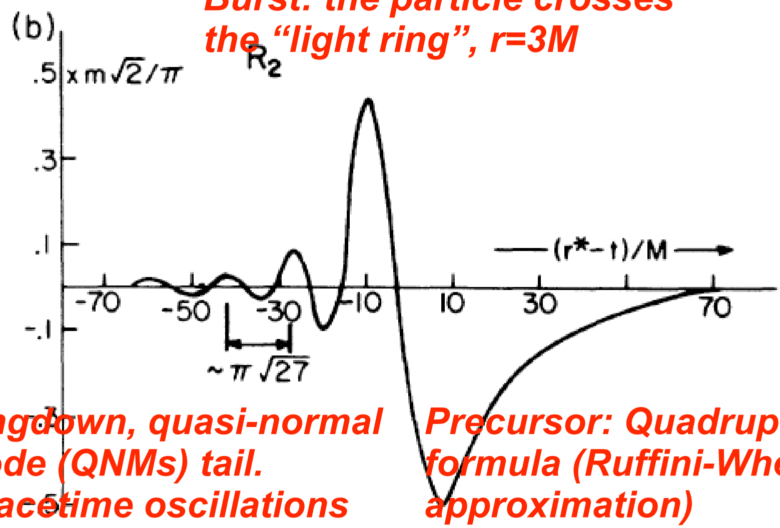
- Padé resummation [Padé1892]

- h(t) : [Davis, Ruffini, Tiomno 1972]
  CLAP [Price-Pullin 1994]

F_ϕ [DIS1998]
A(r) [DJS00]
Factorized waveform [DN07]

Discovery of the structure:
Precursor (plunge)-Burst (merger)-Ringdown

Burst: the particle crosses the "light ring", r=3M



Ringdown, quasi-normal mode (QNMs) tail. Spacetime oscillations

Precursor: Quadrupole formula (Ruffini-Wheeler approximation)

Some key references

PN

Wagoner & Will 76
Damour & Deruelle 81,82;
Blanchet & Damour 86
Damour & Schafer 88
Blanchet & Damour 89;
Blanchet, Damour Iyer, Will, Wiseman 95
Blanchet 95
Jaranowski & Schafer 98
Damour, Jaranowski, Schafer 01
Blanchet, Damour, Esposito-Farese & Iyer 05
Kidder 07
Blanchet, Faye, Iyer & Sinha, 08

NR

Brandt & Brugmann 97
Baker, Brugmann, Campanelli, Lousto & Takahashi 01
Baker, Campanelli, Lousto & Takahashi 02
Pretorius 05
Baker et al. 05
Campanelli et al. 05
Gonzalez et al. 06
Koppitz et al. 07
Pollney et al. 07
Boyle et al. 07
Scheel et al. 08

EOB

Buonanno & Damour 99, 00
Damour 01
Damour Jaranowski & Schafer 00
Buonanno et al. 06-09
Damour & Nagar 07-09
Damour, Iyer & Nagar 08

2-body Taylor-expanded 3PN Hamiltonian [JS98, DJS00,01]

$$H_{\rm N}(\mathbf{x}_a,\mathbf{p}_a)=\sum_a \frac{\mathbf{p}_a^2}{2\,m_a}-\frac{1}{2}\sum_a\sum_{b\neq a}\frac{G\,m_a\,m_b}{r_{ab}}\,.$$

$$H_{\rm 1PN}(\mathbf{x}_a,\mathbf{p}_a)=-\frac{1}{8}\frac{(\mathbf{p}_1^2)^2}{m_1^3}+\frac{1}{8}\frac{Gm_1m_2}{r_{12}}\left[-12\frac{\mathbf{p}_1^2}{m_1^2}+14\frac{(\mathbf{p}_1\cdot\mathbf{p}_2)}{m_1m_2}+2\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1m_2}\right]+\frac{1}{4}\frac{Gm_1m_2}{r_{12}}\frac{G(m_1+m_2)}{r_{12}}+(1\longleftrightarrow 2),$$

1PN

$$\begin{aligned}
H_{\rm 2PN}(\mathbf{x}_a,\mathbf{p}_a)=&\frac{1}{16}\frac{(\mathbf{p}_1^2)^3}{m_1^5}+\frac{1}{8}\frac{Gm_1m_2}{r_{12}}\left[5\frac{(\mathbf{p}_1^2)^2}{m_1^4}-\frac{11}{2}\frac{\mathbf{p}_1^2\,\mathbf{p}_2^2}{m_1^2m_2^2}-\frac{(\mathbf{p}_1\cdot\mathbf{p}_2)^2}{m_1^2m_2^2}+5\frac{\mathbf{p}_1^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^2m_2^2}\right.\\
&\left.-6\frac{(\mathbf{p}_1\cdot\mathbf{p}_2)\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1^2m_2^2}-\frac{3}{2}\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^2m_2^2}\right]\\
&+\frac{1}{4}\frac{G^2m_1m_2}{r_{12}^2}\left[m_2\left(10\frac{\mathbf{p}_1^2}{m_1^2}+19\frac{\mathbf{p}_2^2}{m_2^2}\right)-\frac{1}{2}(m_1+m_2)\frac{27\,(\mathbf{p}_1\cdot\mathbf{p}_2)+6\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1m_2}\right]\\
&-\frac{1}{8}\frac{Gm_1m_2}{r_{12}}\frac{G^2(m_1^2+5m_1m_2+m_2^2)}{r_{12}^2}+(1\longleftrightarrow 2).
\end{aligned}$$

2PN

$$\begin{aligned}
H^{\rm reg}_{\rm 3PN}(\mathbf{x}_a,\mathbf{p}_a)=&-\frac{5}{128}\frac{(\mathbf{p}_1^2)^4}{m_1^7}+\frac{1}{32}\frac{Gm_1m_2}{r_{12}}\left[-14\frac{(\mathbf{p}_1^2)^3}{m_1^6}+4\frac{\big((\mathbf{p}_1\cdot\mathbf{p}_2)^2+4\,\mathbf{p}_1^2\,\mathbf{p}_2^2\big)\mathbf{p}_1^2}{m_1^4m_2^2}+\frac{\big(\mathbf{p}_1^2\,\mathbf{p}_2^2-2\,(\mathbf{p}_1\cdot\mathbf{p}_2)^2\big)(\mathbf{p}_1\cdot\mathbf{p}_2)}{m_1^3m_2^3}\right.\\
&-10\frac{\big(\mathbf{p}_1^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2+\mathbf{p}_2^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2\big)\mathbf{p}_1^2}{m_1^4m_2^2}+24\frac{\mathbf{p}_1^2\,(\mathbf{p}_1\cdot\mathbf{p}_2)(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1^4m_2^2}+2\frac{\mathbf{p}_1^2\,(\mathbf{p}_1\cdot\mathbf{p}_2)(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^3m_2^3}\\
&+\frac{\big(7\,\mathbf{p}_1^2\,\mathbf{p}_2^2-10\,(\mathbf{p}_1\cdot\mathbf{p}_2)^2\big)(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1^3m_2^3}+6\frac{\mathbf{p}_1^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^4m_2^2}\\
&\left.+15\frac{(\mathbf{p}_1\cdot\mathbf{p}_2)(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^3m_2^3}-18\frac{\mathbf{p}_1^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)^3}{m_1^3m_2^3}+5\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^3(\mathbf{n}_{12}\cdot\mathbf{p}_2)^3}{m_1^3m_2^3}\right]\\
&+\frac{G^2m_1m_2}{r_{12}^2}\left[\frac{1}{16}(m_1-27m_2)\frac{(\mathbf{p}_1^2)^2}{m_1^4}-\frac{115}{16}m_1\frac{\mathbf{p}_1^2\,(\mathbf{p}_1\cdot\mathbf{p}_2)}{m_1^3m_2}+\frac{1}{48}m_2\frac{25\,(\mathbf{p}_1\cdot\mathbf{p}_2)^2+371\,\mathbf{p}_1^2\,\mathbf{p}_2^2}{m_1^2m_2^2}\right.\\
&+\frac{17}{16}\frac{\mathbf{p}_1^2(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2}{m_1^3}-\frac{1}{8}m_1\frac{\big(15\,\mathbf{p}_1^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_2)+11\,(\mathbf{p}_1\cdot\mathbf{p}_2)\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)\big)(\mathbf{n}_{12}\cdot\mathbf{p}_1)}{m_1^3m_2}+\frac{5}{12}\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^4}{m_1^3}\\
&-\frac{3}{2}m_1\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^3(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1^3m_2}+\frac{125}{12}m_2\frac{(\mathbf{p}_1\cdot\mathbf{p}_2)\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1^2m_2^2}+\frac{10}{3}m_2\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^2m_2^2}\\
&\left.-\frac{1}{48}(220m_1+193m_2)\frac{\mathbf{p}_1^2(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^2m_2^2}\right]+\frac{G^3m_1m_2}{r_{12}^3}\left[-\frac{1}{48}\left(466\,m_1^2+\left(473-\frac{3}{4}\pi^2\right)m_1m_2+150\,m_2^2\right)\frac{\mathbf{p}_1^2}{m_1^2}\right.\\
&+\frac{1}{16}\left(77(m_1^2+m_2^2)+\left(143-\frac{1}{4}\pi^2\right)m_1m_2\right)\frac{(\mathbf{p}_1\cdot\mathbf{p}_2)}{m_1m_2}+\frac{1}{16}\left(61\,m_1^2-\left(43+\frac{3}{4}\pi^2\right)m_1m_2\right)\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2}{m_1^2}\\
&\left.+\frac{1}{16}\left(21(m_1^2+m_2^2)+\left(119+\frac{3}{4}\pi^2\right)m_1m_2\right)\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1m_2}\right]\\
&+\frac{1}{8}\frac{G^4m_1m_2^3}{r_{12}^4}\left[\left(\frac{227}{3}-\frac{21}{4}\pi^2\right)m_1+m_2\right]+(1\longleftrightarrow 2).
\end{aligned}\tag{12}$$

3PN

Taylor-expanded 3PN waveform

Blanchet,Iyer, Joguet 02, Blanchet, Damour, Esposito-Farese, Iyer 04, Kidder 07, Blanchet et al. 08

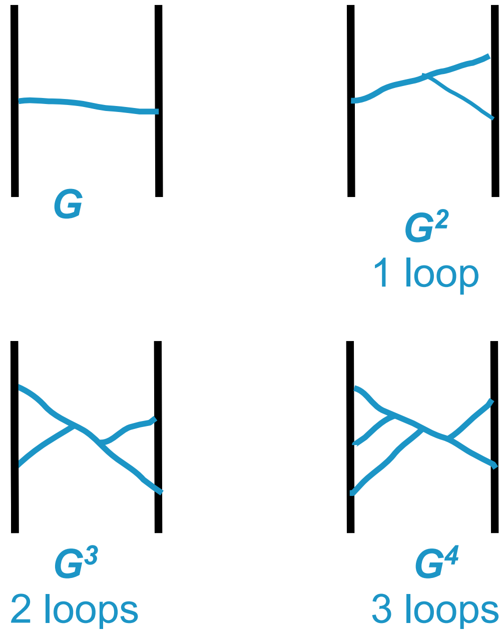
$$
\begin{aligned}
h^{22} = & -8\sqrt{\frac{\pi}{5}}\frac{G\nu m}{c^2 R} e^{-2i\phi} x\Bigg\{1 - x\left(\frac{107}{42} - \frac{55}{42}\nu\right) + x^{3/2}\left[2\pi + 6i\ln\left(\frac{x}{x_0}\right)\right] - x^2\left(\frac{2173}{1512} + \frac{1069}{216}\nu - \frac{2047}{1512}\nu^2\right) \\
& - x^{5/2}\left[\left(\frac{107}{21} - \frac{34}{21}\nu\right)\pi + 24i\nu + \left(\frac{107i}{7} - \frac{34i}{7}\nu\right)\ln\left(\frac{x}{x_0}\right)\right] \\
& + x^3\left[\frac{27\,027\,409}{646\,800} - \frac{856}{105}\gamma_E + \frac{2}{3}\pi^2 - \frac{1712}{105}\ln 2 - \frac{428}{105}\ln x\right. \\
& \left. - 18\left[\ln\left(\frac{x}{x_0}\right)\right]^2 - \left(\frac{278\,185}{33\,264} - \frac{41}{96}\pi^2\right)\nu - \frac{20\,261}{2772}\nu^2 + \frac{114\,635}{99\,792}\nu^3 + \frac{428i}{105}\pi + 12i\pi\ln\left(\frac{x}{x_0}\right)\right] + O(\epsilon^{7/2})\Bigg\},
\end{aligned}
$$

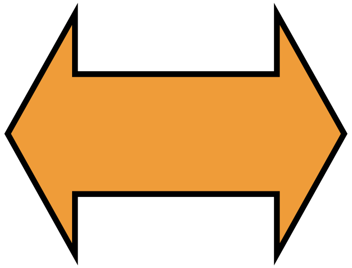
$$x = (M\Omega)^{2/3} \sim v^2/c^2$$

$$
\boxed{
\begin{aligned}
M &= m_1 + m_2 \\
\nu &= m_1 m_2/(m_1 + m_2)^2
\end{aligned}
}
$$

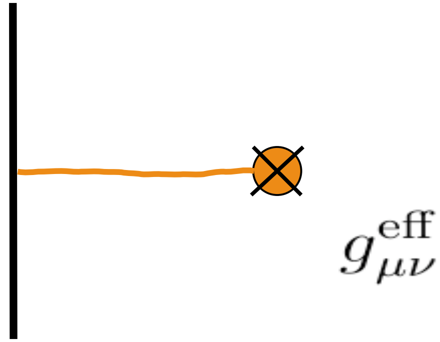
Real dynamics versus Effective dynamics

Real dynamics





Effective dynamics



$$S = -\int \mu ds + \ldots$$

$$H = H_0 + \left(G H_1 + \frac{G^2}{c^2} H_2 + \frac{G^3}{c^4} H_3 + \frac{G^4}{c^6} H_4 \right) \left(1 + \frac{1}{c^2} + \ldots \right)$$

Effective metric

$$ds^2_{\mathrm{eff}} = -A(r)dt^2 + B(r)dr^2 + r^2 \left(d\theta^2 + \sin^2\theta d\varphi^2 \right)$$

Two-body/EOB "correspondence": think quantum-mechanically (Wheeler)

Real 2-body system (m_1, m_2) (in the c.o.m. frame)

an effective particle of mass μ in some effective metric $g_{\mu\nu}^{eff}(M)$

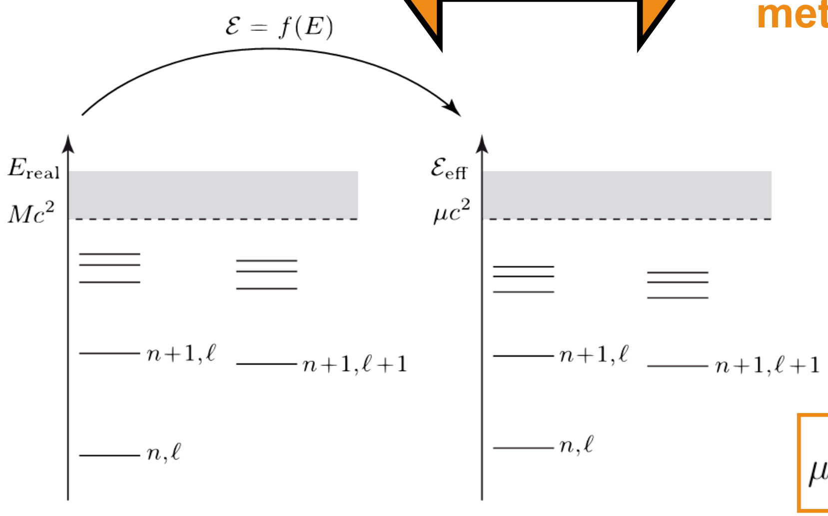


Figure 1: Sketch of the correspondence between the quantized energy levels of the real and effective conservative dynamics. n denotes the 'principal quantum

$$\mu^2 + g_{\rm eff}^{\mu\nu} \frac{\partial S_{\rm eff}}{\partial x^\mu} \frac{\partial S_{\rm eff}}{\partial x^\nu} + \mathcal{O}(p^4) = 0$$

Sommerfeld "Old Quantum Mechanics":

$$J = \ell\hbar = \frac{1}{2\pi} \oint p_\varphi \, d\varphi$$
$$N = n\hbar = I_r + J$$
$$I_r = \frac{1}{2\pi} \oint p_r \, dr$$

$H^{\rm classical}(q,p)$ → $H^{classical}(I_a)$ →

$$E^{\rm quantum}(I_a = n_a h) = f^{-1}\left[\mathcal{E}_{\rm eff}^{\rm quantum}\left(I_a^{\rm eff} = n_a h\right)\right]$$

The 2-body Hamiltonian [at 2PN for clarity]

The 2-body Hamiltonian in the c.o.m frame at 2PN:

$$H_{2\mathrm{PN}}^{\mathrm{relative}}(\boldsymbol{q},\boldsymbol{p}) = H_0(\boldsymbol{q},\boldsymbol{p}) + \frac{1}{c^2} H_2(\boldsymbol{q},\boldsymbol{p}) + \frac{1}{c^4} H_4(\boldsymbol{q},\boldsymbol{p})$$

The Newtonian limit :

$$H_0(\boldsymbol{q},\boldsymbol{p}) = \frac{1}{2\mu}\boldsymbol{p}^2 + \frac{GM\mu}{|\boldsymbol{q}|}$$

4 additional terms at 1PN

7 additional terms at 2PN

11 additional terms at 3PN

Rewrite the c.o.m. energy using *action variables* (à la Sommerfeld):
obtain the "quantum" energy levels [from Damour&Schaefer 1988]

$$E_{2\mathrm{PN}}^{\mathrm{relative}}(n,\ell) = -\frac{1}{2}\mu\frac{\alpha^2}{n^2}\left[1 + \frac{\alpha^2}{c^2}\left(\frac{c_{11}}{n\ell} + \frac{c_{20}}{n^2}\right) + \frac{\alpha^4}{c^4}\left(\frac{c_{13}}{n\ell^3} + \frac{c_{22}}{n^2\ell^2} + \frac{c_{31}}{n^3\ell} + \frac{c_{40}}{n^4}\right)\right]$$

"Balmer" formula

"Delaunay Hamiltonian"

$$\nu \equiv \mu/M$$

$$\left.\begin{aligned} J &= \ell\hbar = \frac{1}{2\pi}\oint p_\varphi\, d\varphi \\ N &= n\hbar = I_r + J \\ I_r &= \frac{1}{2\pi}\oint p_r\, dr \end{aligned}\right\} \quad \alpha \equiv GM\mu/\hbar = G m_1 m_2/\hbar$$

The 3PN EOB Hamiltonian

Real 2-body system *(m_1, m_2)* (in the c.o.m. frame)

1:1 map

an effective particle of mass $\mu = m_1 m_2/(m_1+m_2)$ in some effective metric $g_{\mu\nu}^{eff}(M)$

Simple energy map

$$\boxed{\mathcal{E}_{eff} = \frac{s - m_1^2 - m_2^2}{2M}}$$

$$s = E_{\rm real}^2$$

$$\boxed{H_{\rm EOB} = M\sqrt{1 + 2\nu\left(\hat{H}_{\rm eff} - 1\right)}}$$

$$\boxed{\begin{aligned} M &= m_1 + m_2 \\ \nu &= m_1 m_2/(m_1+m_2)^2 \end{aligned}}$$

Simple effective Hamiltonian

$$\boxed{\hat{H}_{\rm eff} \equiv \sqrt{p_{r_*}^2 + A\left(1 + \frac{p_\varphi^2}{r^2} + z_3 \frac{p_{r_*}^4}{r^2}\right)}.}$$

crucial EOB "radial potential" A(r)

$$p_{r_*} = \left(\frac{A}{B}\right)^{1/2} p_r$$

Explicit form of the effective metric

The effective metric $g_{\mu\nu}^{eff}(M)$ at 3PN

$$ds^2 = -A(r)dt^2 + B(r)dr^2 + r^2(d\theta^2 + \sin^2\theta d\varphi^2).$$

where the coefficients are a ν-dependent "deformation" of the Schwarzschild ones:

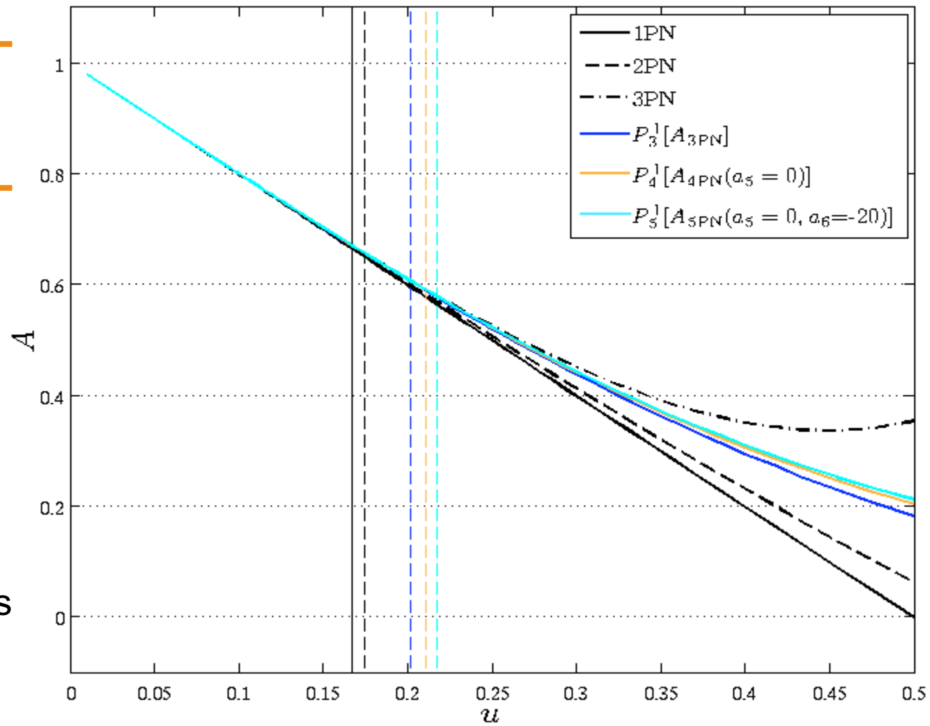
$$\boxed{A_{3PN}(R) = 1 - 2u + 2\nu\, u^3 + a_4 \nu\, u^4} \qquad a_4 = \frac{94}{3} - \frac{41}{32}\pi^2 \simeq 18.6879027$$

$$(A(R)B(R))_{3PN} = 1 - 6\nu u^2 + 2(3\nu - 26)\nu u^3$$

u = 1/r

$$A(u; a5, a6, \nu) = P_5^1\left[A_{3PN}(u) + \nu a_5 u^5 + \nu a_6 u^6\right]$$

- Compact representation of PN dynamics
- Bad behaviour at 3PN. Use Padé resummation of A(r) to have an effective horizon.
- Impose [by continuity with the ν=0 case] that A(r) has a simple zero [at r≈2].
- The a_5 and a_6 constants parametrize (yet) uncalculated 4PN corrections and 5PN corrections



2-body Taylor-expanded 3PN Hamiltonian [JS98, DJS00,01]

$$H_{\rm N}(\mathbf{x}_a, \mathbf{p}_a) = \sum_a \frac{\mathbf{p}_a^2}{2\, m_a} - \frac{1}{2} \sum_a \sum_{b\neq a} \frac{G\, m_a\, m_b}{r_{ab}}\,.$$

$$H_{\rm 1PN}(\mathbf{x}_a, \mathbf{p}_a) = -\frac{1}{8}\frac{(\mathbf{p}_1^2)^2}{m_1^3} + \frac{1}{8}\frac{G m_1 m_2}{r_{12}}\left[-12\frac{\mathbf{p}_1^2}{m_1^2} + 14\frac{(\mathbf{p}_1\cdot\mathbf{p}_2)}{m_1 m_2} + 2\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1 m_2}\right] + \frac{1}{4}\frac{G m_1 m_2}{r_{12}}\frac{G(m_1+m_2)}{r_{12}} + (1\longleftrightarrow 2),$$

1PN

$$\begin{aligned}
H_{\rm 2PN}(\mathbf{x}_a, \mathbf{p}_a) = {} & \frac{1}{16}\frac{(\mathbf{p}_1^2)^3}{m_1^5} + \frac{1}{8}\frac{G m_1 m_2}{r_{12}}\Bigg[5\frac{(\mathbf{p}_1^2)^2}{m_1^4} - \frac{11}{2}\frac{\mathbf{p}_1^2\,\mathbf{p}_2^2}{m_1^2 m_2^2} - \frac{(\mathbf{p}_1\cdot\mathbf{p}_2)^2}{m_1^2 m_2^2} + 5\frac{\mathbf{p}_1^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^2 m_2^2} \\
& -6\frac{(\mathbf{p}_1\cdot\mathbf{p}_2)\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1^2 m_2^2} - \frac{3}{2}\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^2 m_2^2}\Bigg] \\
& +\frac{1}{4}\frac{G^2 m_1 m_2}{r_{12}^2}\left[m_2\left(10\frac{\mathbf{p}_1^2}{m_1^2} + 19\frac{\mathbf{p}_2^2}{m_2^2}\right) - \frac{1}{2}(m_1+m_2)\frac{27\,(\mathbf{p}_1\cdot\mathbf{p}_2) + 6\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1 m_2}\right] \\
& -\frac{1}{8}\frac{G m_1 m_2}{r_{12}}\frac{G^2(m_1^2 + 5 m_1 m_2 + m_2^2)}{r_{12}^2} + (1\longleftrightarrow 2).
\end{aligned}$$

2PN

$$\begin{aligned}
H^{\rm reg}_{\rm 3PN}(\mathbf{x}_a, \mathbf{p}_a) = {} & -\frac{5}{128}\frac{(\mathbf{p}_1^2)^4}{m_1^7} + \frac{1}{32}\frac{G m_1 m_2}{r_{12}}\Bigg[-14\frac{(\mathbf{p}_1^2)^3}{m_1^6} + 4\frac{\big((\mathbf{p}_1\cdot\mathbf{p}_2)^2 + 4\,\mathbf{p}_1^2\,\mathbf{p}_2^2\big)\mathbf{p}_1^2}{m_1^4 m_2^2} + \frac{\big(\mathbf{p}_1^2\,\mathbf{p}_2^2 - 2\,(\mathbf{p}_1\cdot\mathbf{p}_2)^2\big)(\mathbf{p}_1\cdot\mathbf{p}_2)}{m_1^3 m_2^3} \\
& -10\frac{\big(\mathbf{p}_1^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2 + \mathbf{p}_2^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2\big)\mathbf{p}_1^2}{m_1^4 m_2^2} + 24\frac{\mathbf{p}_1^2\,(\mathbf{p}_1\cdot\mathbf{p}_2)(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1^4 m_2^2} + 2\frac{\mathbf{p}_1^2\,(\mathbf{p}_1\cdot\mathbf{p}_2)(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^3 m_2^3} \\
& +\frac{\big(7\,\mathbf{p}_1^2\,\mathbf{p}_2^2 - 10\,(\mathbf{p}_1\cdot\mathbf{p}_2)^2\big)(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1^3 m_2^3} + 6\frac{\mathbf{p}_1^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^4 m_2^2} \\
& +15\frac{(\mathbf{p}_1\cdot\mathbf{p}_2)(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^3 m_2^3} - 18\frac{\mathbf{p}_1^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)^3}{m_1^3 m_2^3} + 5\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^3(\mathbf{n}_{12}\cdot\mathbf{p}_2)^3}{m_1^3 m_2^3}\Bigg] \\
& +\frac{G^2 m_1 m_2}{r_{12}^2}\Bigg[\frac{1}{16}(m_1 - 27 m_2)\frac{(\mathbf{p}_1^2)^2}{m_1^4} - \frac{115}{16}m_1\frac{\mathbf{p}_1^2\,(\mathbf{p}_1\cdot\mathbf{p}_2)}{m_1^3 m_2} + \frac{1}{48}m_2\frac{25\,(\mathbf{p}_1\cdot\mathbf{p}_2)^2 + 371\,\mathbf{p}_1^2\,\mathbf{p}_2^2}{m_1^2 m_2^2} \\
& +\frac{17}{16}\frac{\mathbf{p}_1^2(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2}{m_1^3} - \frac{1}{8}m_1\frac{\big(15\,\mathbf{p}_1^2\,(\mathbf{n}_{12}\cdot\mathbf{p}_2) + 11\,(\mathbf{p}_1\cdot\mathbf{p}_2)\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)\big)(\mathbf{n}_{12}\cdot\mathbf{p}_1)}{m_1^3 m_2} + \frac{5}{12}\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^4}{m_1^3} \\
& -\frac{3}{2}m_1\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^3(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1^3 m_2} + \frac{125}{12}m_2\frac{(\mathbf{p}_1\cdot\mathbf{p}_2)\,(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1^2 m_2^2} + \frac{10}{3}m_2\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^2 m_2^2} \\
& -\frac{1}{48}(220 m_1 + 193 m_2)\frac{\mathbf{p}_1^2(\mathbf{n}_{12}\cdot\mathbf{p}_2)^2}{m_1^2 m_2^2}\Bigg] + \frac{G^3 m_1 m_2}{r_{12}^3}\Bigg[-\frac{1}{48}\left(466\,m_1^2 + \left(473 - \frac{3}{4}\pi^2\right)m_1 m_2 + 150\,m_2^2\right)\frac{\mathbf{p}_1^2}{m_1^2} \\
& +\frac{1}{16}\left(77(m_1^2 + m_2^2) + \left(143 - \frac{1}{4}\pi^2\right)m_1 m_2\right)\frac{(\mathbf{p}_1\cdot\mathbf{p}_2)}{m_1 m_2} + \frac{1}{16}\left(61\,m_1^2 - \left(43 + \frac{3}{4}\pi^2\right)m_1 m_2\right)\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)^2}{m_1^2} \\
& +\frac{1}{16}\left(21(m_1^2 + m_2^2) + \left(119 + \frac{3}{4}\pi^2\right)m_1 m_2\right)\frac{(\mathbf{n}_{12}\cdot\mathbf{p}_1)(\mathbf{n}_{12}\cdot\mathbf{p}_2)}{m_1 m_2}\Bigg] \\
& +\frac{1}{8}\frac{G^4 m_1 m_2^3}{r_{12}^4}\left[\left(\frac{227}{3} - \frac{21}{4}\pi^2\right)m_1 + m_2\right] + (1\longleftrightarrow 2).
\end{aligned} \tag{12}$$

3PN

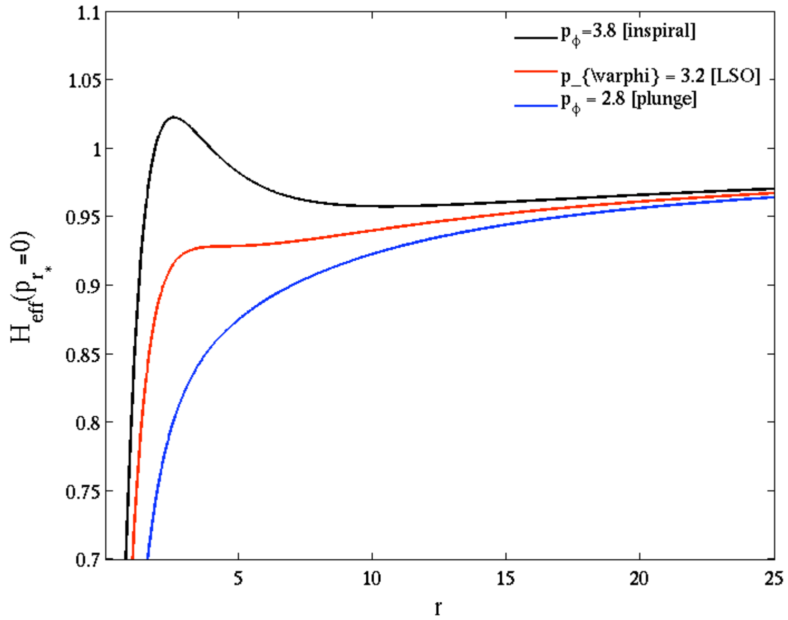
Hamilton's equation + radiation reaction

$$\frac{dr}{dt} = \left(\frac{A}{B}\right)^{1/2} \frac{\partial \hat{H}_{\rm EOB}}{\partial p_{r_*}} ,$$

$$\frac{dp_{r_*}}{dt} = -\left(\frac{A}{B}\right)^{1/2} \frac{\partial \hat{H}_{\rm EOB}}{\partial r} ,$$

$$\Omega \equiv \frac{d\varphi}{dt} = \frac{\partial \hat{H}_{\rm EOB}}{\partial p_\varphi} ,$$

$$\frac{dp_\varphi}{dt} = \hat{\mathcal{F}}_\varphi \, .$$

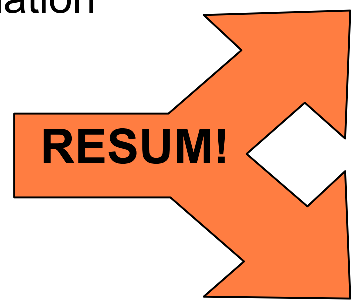


The system must lose mechanical angular momentum

Use PN-expanded result for ***GW angular momentum flux*** as a starting point.
Needs resummation to have a better behavior during late-inspiral and plunge.

PN calculations are done in the circular approximation

$$\hat{\mathcal{F}}_\varphi^{\rm Taylor} = -\frac{32}{5}\,\nu\,\Omega^5\,r_\omega^4\,\hat{F}^{\rm Taylor}(v_\varphi)$$



Parameter-dependent

EOB 1.* **[DIS 1998, DN07]**

Parameter -free:

EOB 2.0 **[DIN 2008, DN09]**

Taylor-expanded 3PN waveform

Blanchet,Iyer, Joguet 02, Blanchet, Damour, Esposito-Farese, Iyer 04, Kidder 07, Blanchet et al. 08

$$
\begin{aligned}
h^{22} = & -8\sqrt{\frac{\pi}{5}}\frac{G\nu m}{c^2 R}e^{-2i\phi}x\Bigg\{1 - x\left(\frac{107}{42} - \frac{55}{42}\nu\right) + x^{3/2}\left[2\pi + 6i\ln\left(\frac{x}{x_0}\right)\right] - x^2\left(\frac{2173}{1512} + \frac{1069}{216}\nu - \frac{2047}{1512}\nu^2\right) \\
& - x^{5/2}\left[\left(\frac{107}{21} - \frac{34}{21}\nu\right)\pi + 24i\nu + \left(\frac{107i}{7} - \frac{34i}{7}\nu\right)\ln\left(\frac{x}{x_0}\right)\right] \\
& + x^3\Bigg[\frac{27\,027\,409}{646\,800} - \frac{856}{105}\gamma_E + \frac{2}{3}\pi^2 - \frac{1712}{105}\ln 2 - \frac{428}{105}\ln x \\
& - 18\left[\ln\left(\frac{x}{x_0}\right)\right]^2 - \left(\frac{278\,185}{33\,264} - \frac{41}{96}\pi^2\right)\nu - \frac{20\,261}{2772}\nu^2 + \frac{114\,635}{99\,792}\nu^3 + \frac{428i}{105}\pi + 12i\pi\ln\left(\frac{x}{x_0}\right)\Bigg] + O(\epsilon^{7/2})\Bigg\},
\end{aligned}
$$

$$x = (M\Omega)^{2/3} \sim v^2/c^2$$

$$\boxed{\begin{aligned} M &= m_1 + m_2 \\ \nu &= m_1 m_2/(m_1+m_2)^2 \end{aligned}}$$

EOB 2.0: new resummation procedures (DN07, DIN 2008)

- Resummation of the waveform **multipole by multipole**
- **Factorized** waveform for any (l,m) at the highest available PN order (start from PN results of Blanchet et al.)

Next-to-Quasi-Circular correction

$$h_{\ell m} = h_{\ell m}^{(N)} \hat{h}_{\ell m}^{(\epsilon)} f_{\ell m}^{\rm NQC}$$

Newtonian x PN-correction

remnant phase correction

$$\hat{h}_{\ell m}^{(\epsilon)} = \hat{S}_{\rm eff}^{(\epsilon)} T_{\ell m} e^{{\rm i}\delta_{\ell m}} \rho_{\ell m}^{\ell}$$

remnant modulus correction:

- l-th power of the (expanded) l-th root of f_{lm}
- improves the behavior of PN corrections

The "Tail factor"

$$\boxed{T_{\ell m} = \frac{\Gamma(\ell+1-2{\rm i}\hat{\hat{k}})}{\Gamma(\ell+1)} e^{\pi \hat{\hat{k}}} e^{2{\rm i}\hat{\hat{k}}\log(2kr_0)}}$$

Effective source:

EOB (effective) energy (even-parity)

Angular momentum (odd-parity)

resums an infinite number of leading logarithms in tail effects

Radiation reaction: parameter-free resummation

$$\boxed{\mathcal{F}_\varphi \equiv -\frac{1}{8\pi\,\Omega}\sum_{\ell=2}^{\ell_{\max}}\sum_{m=1}^{\ell}(m\,\Omega)^2\,|R\,h_{\ell m}^{(\epsilon)}|^2}$$

$$h_{\ell m} = h_{\ell m}^{(N)}\hat{h}_{\ell m}^{(\epsilon)} f_{\ell m}^{\rm NQC}$$

$$\hat{h}_{\ell m}^{(\epsilon)} = \hat{S}_{\rm eff}^{(\epsilon)} T_{\ell m} e^{{\rm i}\delta_{\ell m}} \rho_{\ell m}^{\ell}$$

$$\begin{aligned}\rho_{22}(x;\nu) = 1 &+ \left(\frac{55\nu}{84} - \frac{43}{42}\right)x + \left(\frac{19583\nu^2}{42336} - \frac{33025\nu}{21168} - \frac{20555}{10584}\right)x^2 \\ &+ \left(\frac{10620745\nu^3}{39118464} - \frac{6292061\nu^2}{3259872} + \frac{41\pi^2\nu}{192} - \frac{48993925\nu}{9779616} - \frac{428}{105}\text{eulerlog}_2(x) + \frac{1556919113}{122245200}\right)x^3 \\ &+ \left(\frac{9202}{2205}\text{eulerlog}_2(x) - \frac{387216563023}{160190110080}\right)x^4 + \left(\frac{439877}{55566}\text{eulerlog}_2(x) - \frac{16094530514677}{533967033600}\right)x^5 + \mathcal{O}(x^6),\end{aligned}$$

- Different possible representations of the residual amplitude correction [Padé]
- The "adiabatic" EOB parameters (***a***$_5$, ***a***$_6$) propagate in radiation reaction via the effective source.

The **factorized** waveform: getting into details

The Newtonian factor

$$h_{\ell m}^{(N,\epsilon)} = \frac{GM\nu}{c^2 R} n_{\ell m}^{(\epsilon)} c_{\ell+\epsilon}(\nu) x^{(\ell+\epsilon)/2} Y^{\ell-\epsilon,-m}\left(\frac{\pi}{2}, \Phi\right)$$

$$n_{\ell m}^{(0)} = (\mathrm{i}m)^{\ell} \frac{8\pi}{(2\ell+1)!!} \sqrt{\frac{(\ell+1)(\ell+2)}{\ell(\ell-1)}},$$

$$n_{\ell m}^{(1)} = -(\mathrm{i}m)^{\ell} \frac{16\pi \mathrm{i}}{(2\ell+1)!!} \sqrt{\frac{(2\ell+1)(\ell+2)(\ell^2-m^2)}{(2\ell-1)(\ell+1)\ell(\ell-1)}};$$

The effective source

$$\hat{S}_{\rm eff}^{(0)}(x) = \hat{H}_{\rm eff}(x) \qquad \ell+m \quad \text{even}$$

$$\hat{S}_{\rm eff}^{(1,J)} = \hat{j}(x) \equiv x^{1/2} j(x) \qquad \ell+m \quad \text{odd}$$

$$\hat{H}_{\rm eff} = \frac{H_{\rm eff}}{\mu} = \sqrt{A(u)(1+j^2u^2)}$$

Residual phase

$$\delta_{22} = \frac{7}{3} y^{3/2} + \frac{428\pi}{105} y^3 - 24\nu \bar{y}^{5/2}$$

$$y \equiv (H_{\rm real}\Omega)^{2/3}$$

Tail factor (resums an infinite number of leading logarithms)

$$T_{\ell m} = \frac{\Gamma(\ell+1-2\mathrm{i}\hat{\hat{k}})}{\Gamma(\ell+1)} e^{\pi \hat{\hat{k}}} e^{2\mathrm{i}\hat{\hat{k}}\log(2kr_0)}$$

The factorized waveform: resumming the quadrupole amplitude

Residual amplitude factor

$$f_{22}(x;\nu) = 1 + \frac{1}{42}(55\nu - 86)x + \frac{(2047\nu^2 - 6745\nu - 4288)}{1512}x^2 + \left(\frac{114635\nu^3}{99792} - \frac{227875\nu^2}{33264} + \frac{41}{96}\pi^2\nu - \frac{34625\nu}{3696} - \frac{856}{105}\text{eulerlog}_2(x) + \frac{21428357}{727650}\right)x^3 + \left(\frac{36808}{2205}\text{eulerlog}_2(x) - \frac{5391582359}{198648450}\right)x^4 + \left(\frac{458816}{19845}\text{eulerlog}_2(x) - \frac{93684531406}{893918025}\right)x^5 + \mathcal{O}(x^6), \quad (31)$$

Resumming f_{22}:

1. Go straight Padé: $$P^3_2\{f_{22}(x;\nu)\}$$

2. Replace f_{22} by its square-root

$$\rho_{22}(x;\nu) = 1 + \left(\frac{55\nu}{84} - \frac{43}{42}\right)x + \left(\frac{19583\nu^2}{42336} - \frac{33025\nu}{21168} - \frac{20555}{10584}\right)x^2 + \left(\frac{10620745\nu^3}{39118464} - \frac{6292061\nu^2}{3259872} + \frac{41\pi^2\nu}{192} - \frac{48993925\nu}{9779616} - \frac{428}{105}\text{eulerlog}_2(x) + \frac{1556919113}{122245200}\right)x^3 + \left(\frac{9202}{2205}\text{eulerlog}_2(x) - \frac{387216563023}{160190110080}\right)x^4 + \left(\frac{439877}{55566}\text{eulerlog}_2(x) - \frac{16094530514677}{533967033600}\right)x^5 + \mathcal{O}(x^6),$$

The factorized waveform: resumming higher multipoles

Understanding the PN residual amplitude correction

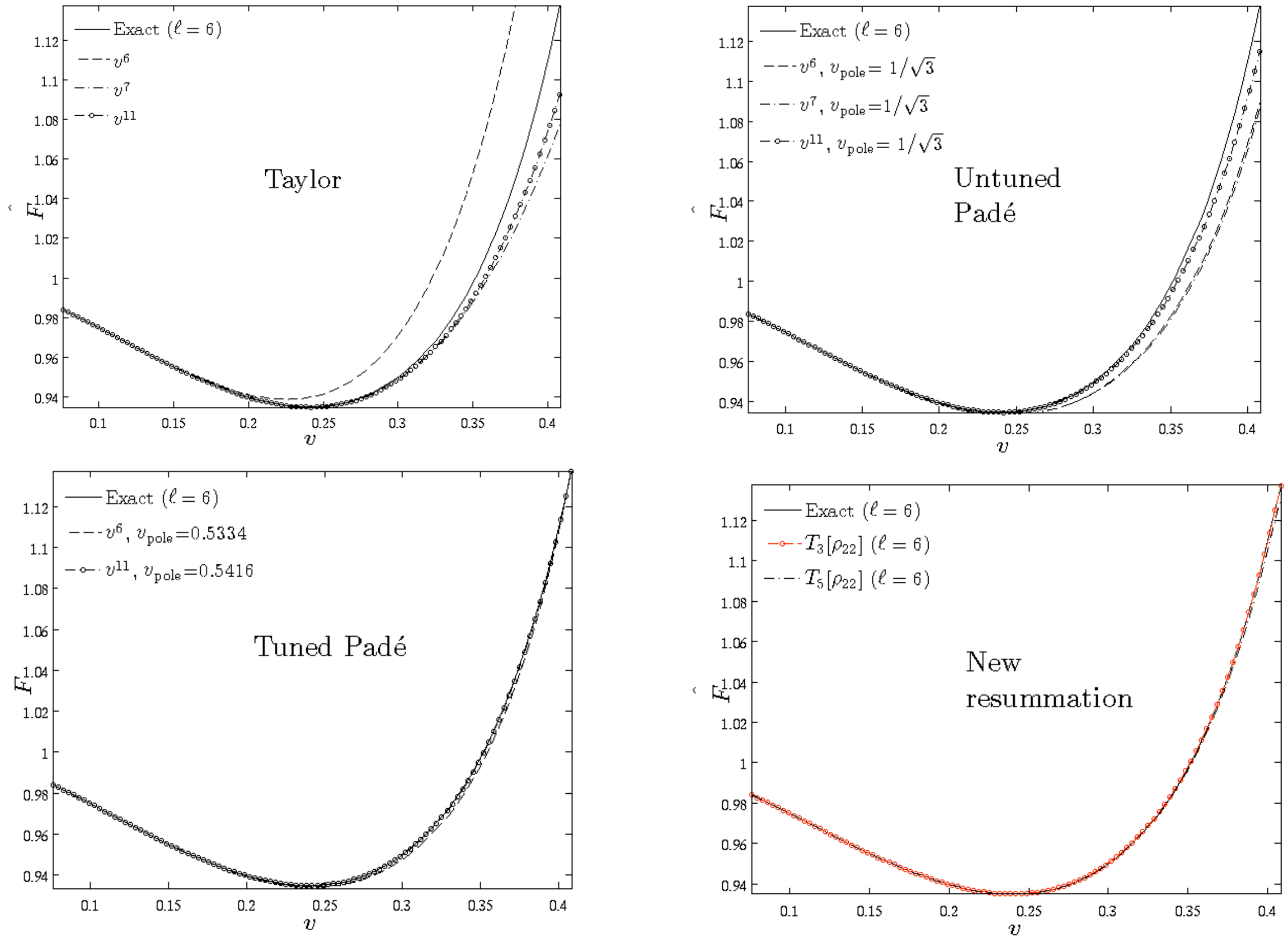
$$f_{\ell m}^{\rm even}(x;0) = 1-\ell x\left(1-\frac{1}{\ell}+\frac{m^2(\ell+9)}{2\ell(\ell+1)(2\ell+3)}\right)+\mathcal{O}(x^2).$$

For higher multipolar orders, use l-th root as default resummation

$$\rho_{\ell m}(x;\nu) = [f_{\ell m}(x;\nu)]^{1/\ell}$$

$$\rho_{\ell m}^{\rm even}(x;0) = 1-x\left(1-\frac{1}{\ell}+\frac{m^2(\ell+9)}{2\ell(\ell+1)(2\ell+3)}\right)+\mathcal{O}(x^2).$$

Test-mass limit (ν=0): circular orbits



Parameter free resummation technique!

EOB 2.0: Next-to-Quasi-Circular correction: EOB ∪ NR

Next-to quasi-circular correction to the *l=m=2 amplitude*

$$\boxed{f_{22}^{\mathrm{NQC}}(a_1, a_2) = 1 + a_1\, p_{r_*}^2/(r\Omega)^2 + a_2\, \ddot{r}/r\,\Omega^2}$$

a_1 & a_2 are determined by requiring:

- The maximum of the (Zerilli-normalized) EOB metric waveform is equal to the maximum of the NR waveform
- That this maximum occurs at the EOB "light-ring" [i.e., maximum of EOB orbital frequency].
- Using **two** NR data: maximum $\boxed{\varphi(\nu) \simeq 0.3215\nu(1 - 0.131(1 - 4\nu))}$
- NQC correction is added consistently in RR. **Iteration until a_1 & a_2 stabilize**

Remaining EOB 2.0 flexibility:

$$A(u; a_5, a_6, \nu) \equiv P_5^1[A^{3\mathrm{PN}}(u) + \nu a_5 u^5 + \nu a_6 u^6]$$

Use Caltech-Cornell [inspiral-plunge] data to constrain ***(a_5,a_6)***

A wide region of correlated values ***(a_5,a_6)*** exists where the phase difference can be reduced at the level of the numerical error (<0.02 radians) during the inspiral

EOB *metric* gravitational waveform: merger and ringdown

EOB approximate representation of the merger (DRT1972 inspired) :

- sudden change of description around the "EOB light-ring" $t=t_m$ (maximum of orbital frequency)
- "match" the insplunge waveform to a superposition of QNMs of the final Kerr black hole
- matching on a 5-teeth comb (*found efficient in the test-mass limit, DN07a*)
- comb of width around 7M centered on the "EOB light-ring"
- use 5 positive frequency QNMs (found to be near-optimal in the test-mass limit)
- Final BH mass and angular momentum are computed from a fit to NR ringdown (*5 eqs for 5 unknowns*)

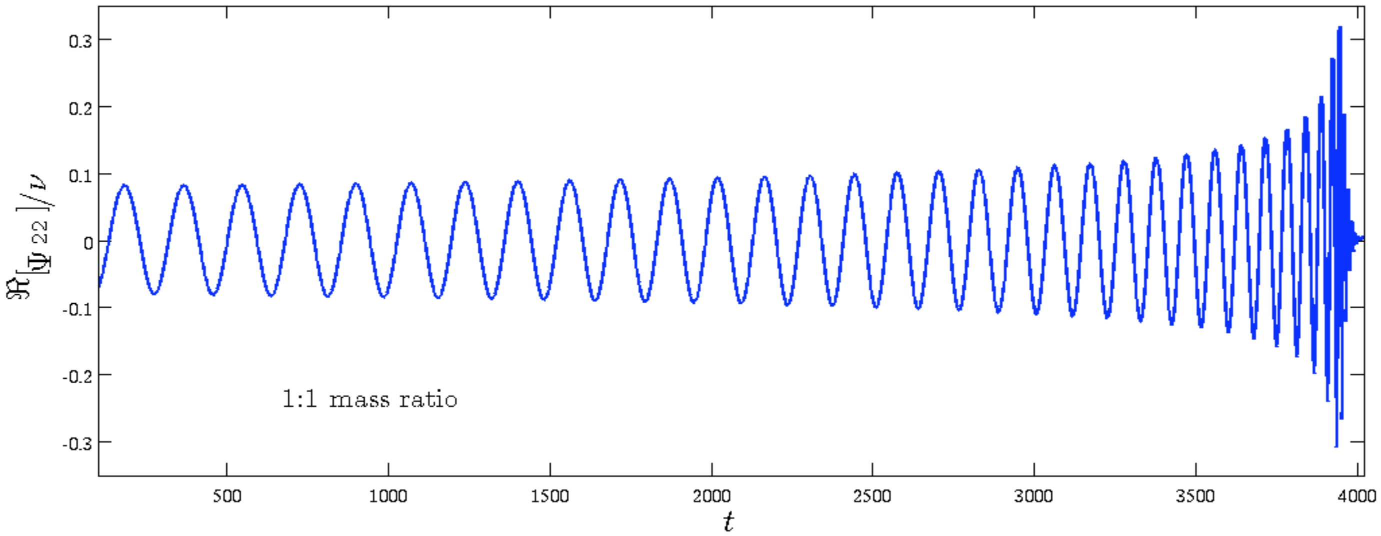
$$\Psi_{22}^{\text{ringdown}}(t) = \sum_N C_N^+ e^{-\sigma_N^+ t} .$$

Total EOB waveform covering inspiral-merger and ringdown

$$\boxed{h_{22}^{\text{EOB}}(t) = \theta(t_m - t)\, h_{22}^{\text{insplunge}}(t) + \theta(t - t_m)\, h_{22}^{\text{ringdown}}(t)}$$

Binary BH coalescence: Numerical Relativity waveform

1:1 (no spin) Caltech-Cornell simulation. Inspiral: $\Delta\phi<0.02$ rad; Ringdown: $\Delta\phi\sim0.05$ rad
Boyle et al 07, Scheel et al 09



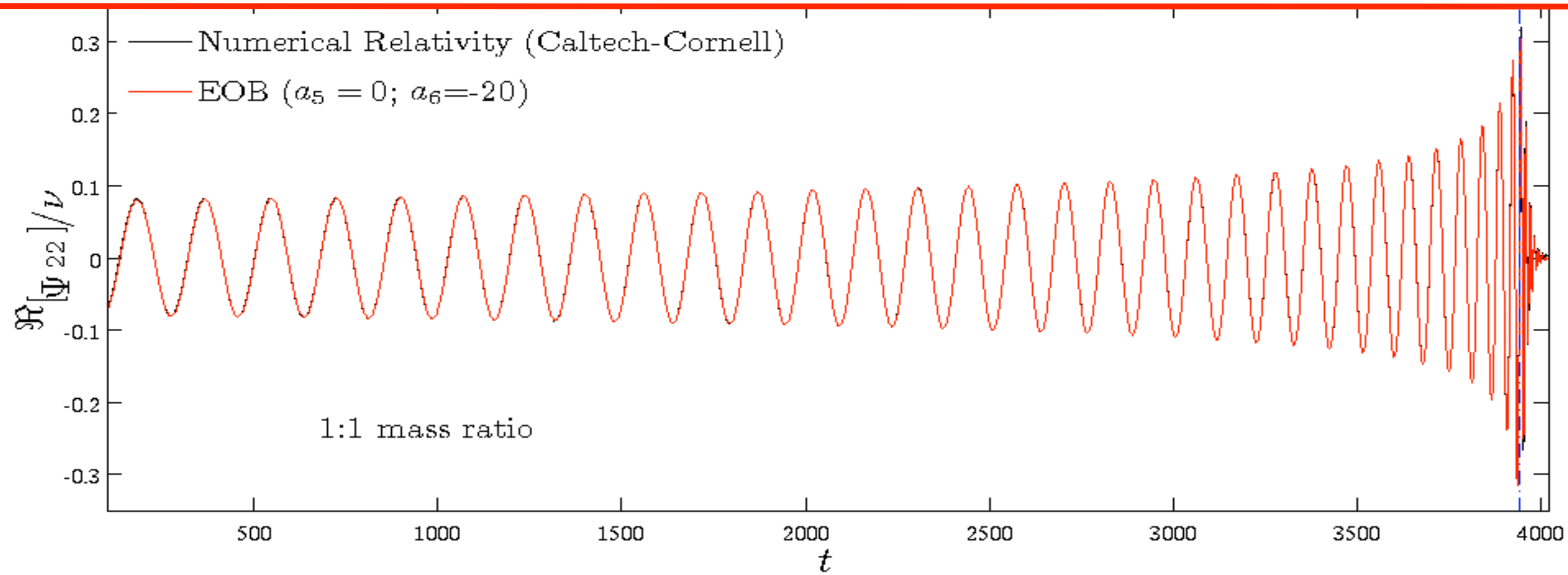
Ringdown

Early inspiral

Late inspiral & Merger

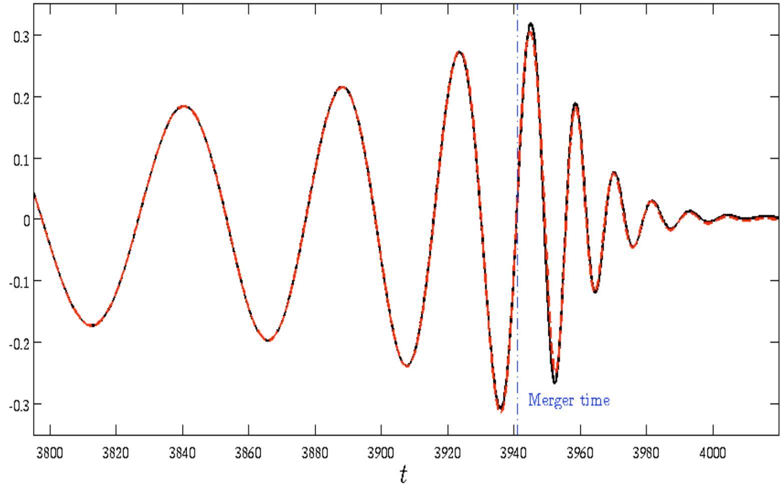
- Late inspiral and merger is non perturbative
- Only describable by NR ?

Comparison Effective-One-Body (EOB) vs NR waveforms



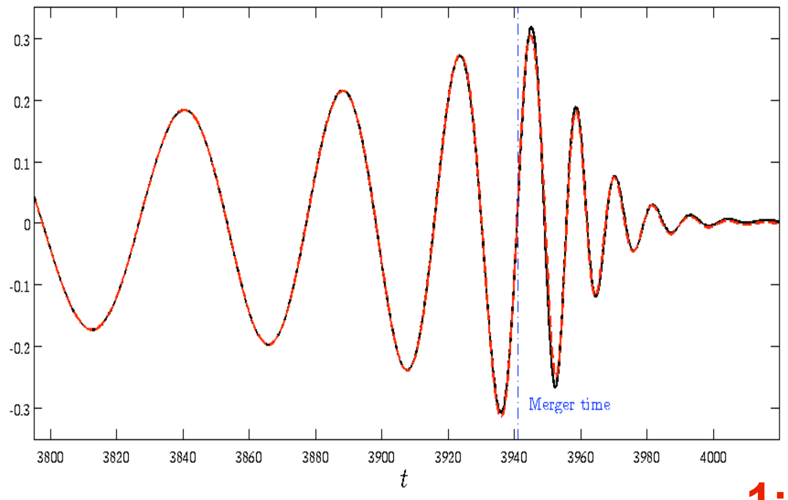
"New" EOB formalism: EOB 2.0$_{NR}$

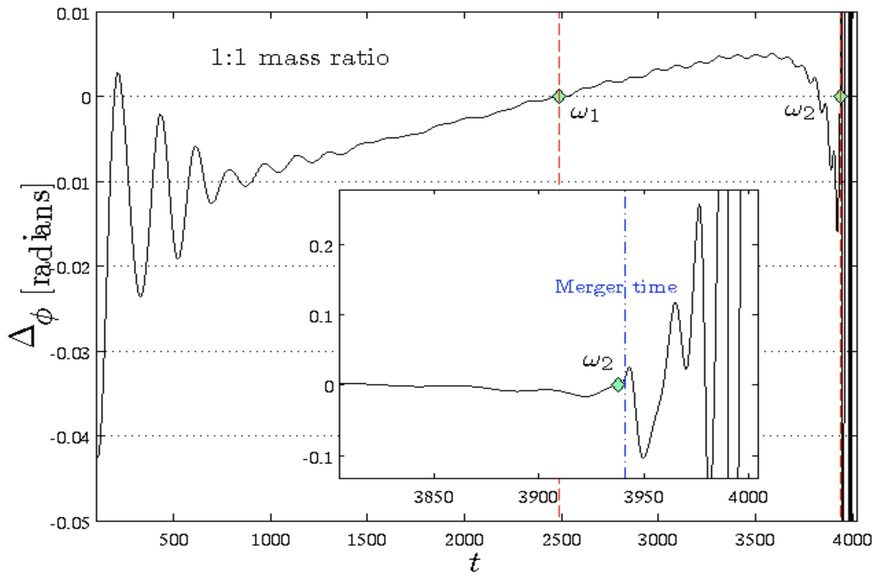
- Two unknown EOB parameters: 4PN and 5PN effective corrections in 2-body Hamiltonian, ***(a_5,a_6)***
- NR calibration of the maximum GW amplitude
- Need to "tune" only one parameter
- **Banana-like "best region"** in the (a_5,a_6) plane extending from (0,-20) to (-36, 520) (where $\Delta\phi \leq 0.02$)



Damour & Nagar, Phys. Rev. D **79**, 081503(R), (2009)
Damour, Iyer & Nagar, Phys. Rev. D **79**, 064004 (2009)

EOB 2.0 & NR comparison: 1:1 & 2:1 mass ratios

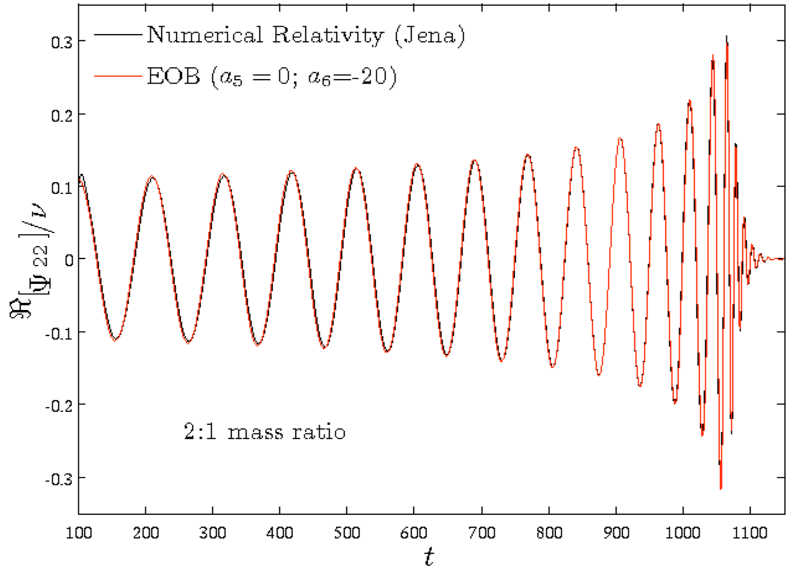




1:1

$a_5 = 0$, $a_6 = -20$

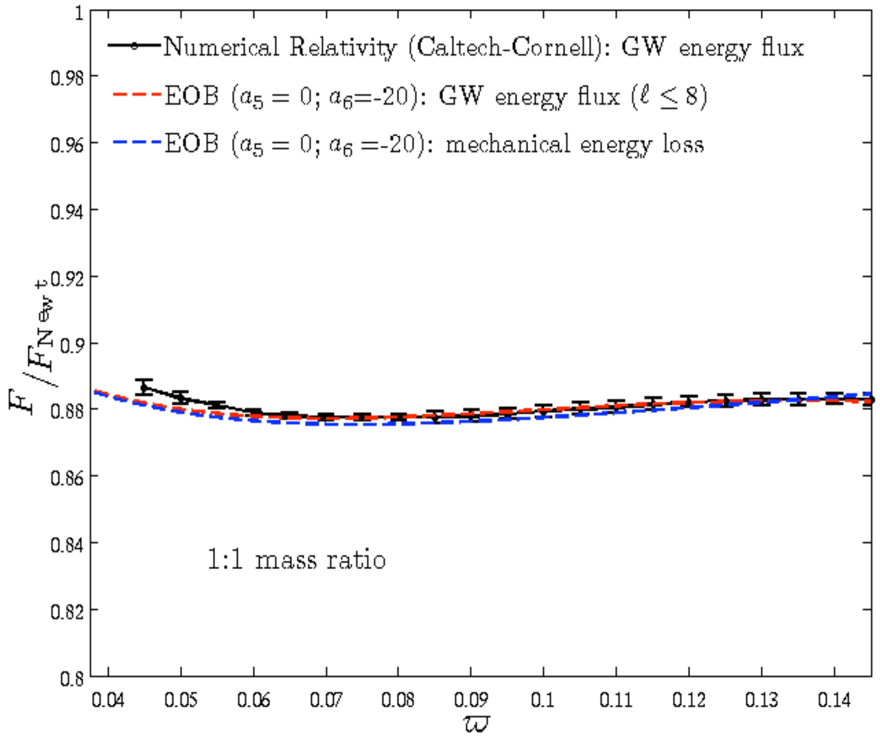
2:1



D, N, Hannam, Husa, Brügmann 08

Agreement: EOB Mechanical loss and NR energy flux

Highly accurate data from Boyle et al, Phys. Rev. D **78**, 104020 (2008) [inspiral only]



New, self-consistent EOB 2.0_{NR}

Damour & Nagar, Phys. Rev. D **79**, 081503(R) (2009)

EOB 1.3, Padé [DIS] & Taylor T4

Boyle et al, Phys. Rev. D **78**, 104020(2008)

EOB 1.5: Buonanno, Pan, Pfeiffer, Scheel, Buchman & Kidder, Phys Rev.D79, 124028 (2009)

- **EOB formalism: EOB 1.5 U NR**

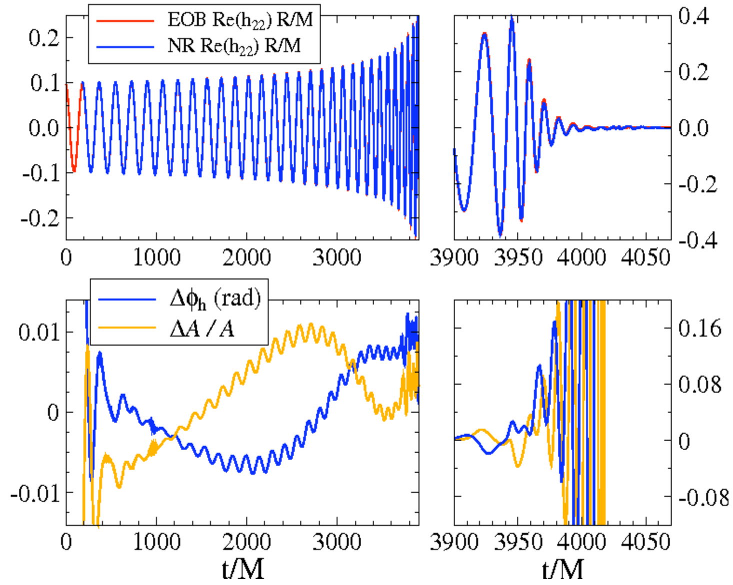
h_{lm} [RWZ] NR 1:1. **EOB resummed waveform** (à la DIN)

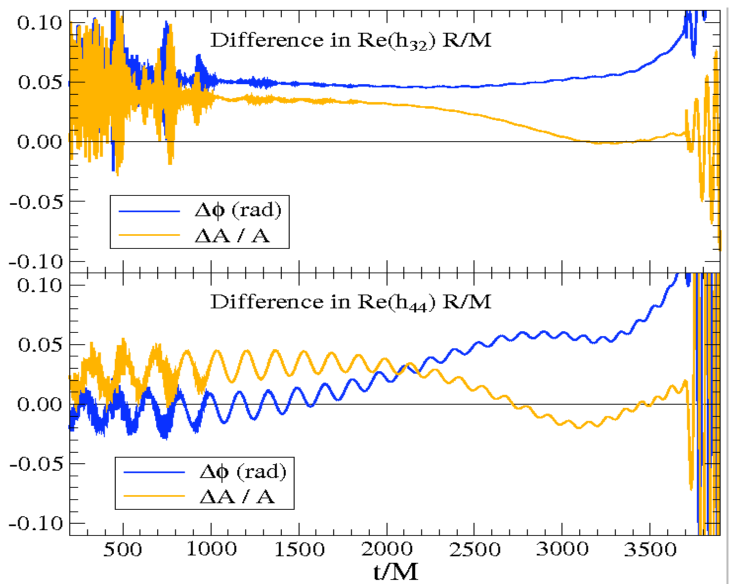
$$\left.\begin{aligned} a_5 &= 25.375 \\ v_{pole}(\nu=1/4) &= 0.85 \end{aligned}\right\} \quad \textbf{reference values}$$

$\Delta t^{22}_{match} = 3.0M$

$a_1 = -2.23$
$a_2 = 31.93$
$a_3 = 3.66$
$a_4 = -10.85$

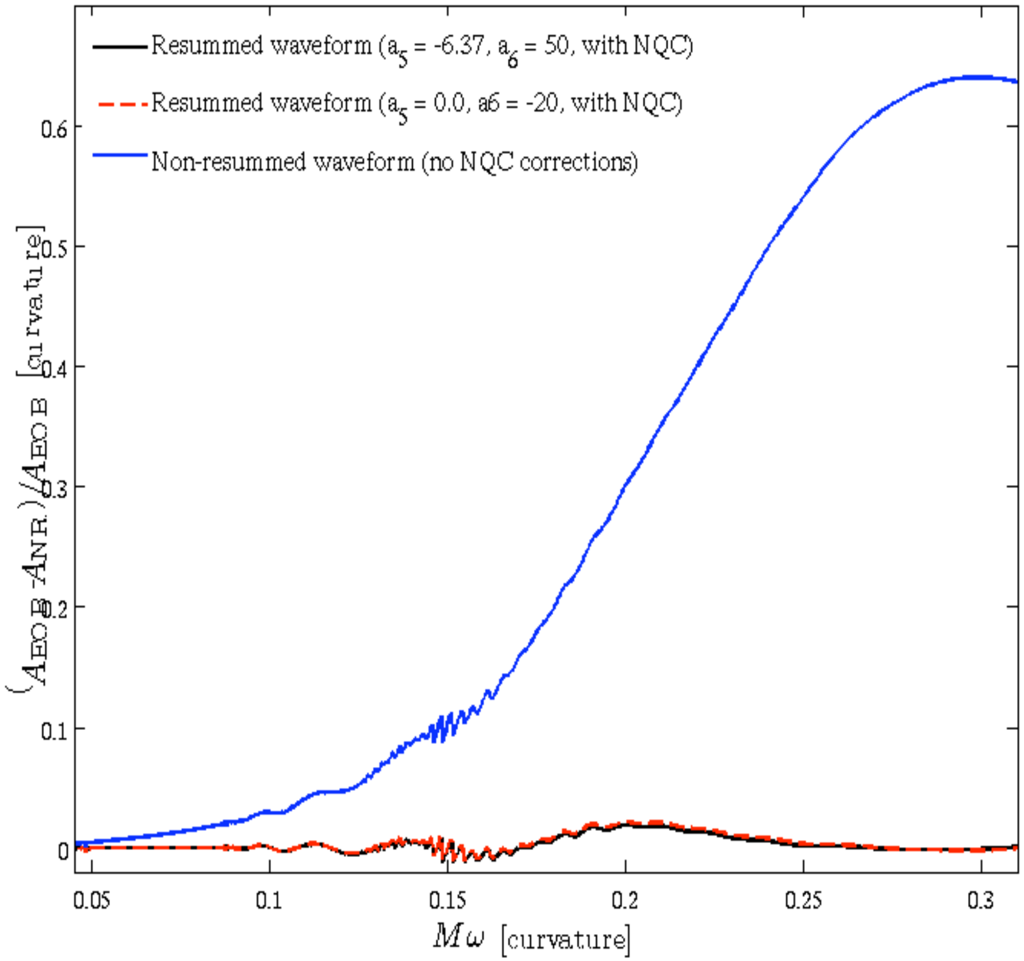
$-0.02 \leq \Delta\phi \leq +0.02 \quad -0.02 \leq DA/A \leq +0.02$ [l=m=2]





- Here, 1:1 mass ratio (**with higher multipoles**)
- **Plus** 2:1 & 3:1 [inspiral only] mass ratios

(Fractional) curvature amplitude difference EOB-NR



- Nonresummed: fractional differences start at the 0.5% level and build up to more than 60%! (just before merger)
- New resummed EOB amplitude+NQC corrections: fractional differences start at the 0.04% level and build up to only 2% (just before merger)
- *Resum+NQC: factor ~30 improvement!*

Shows the effectiveness of resummation techniques, even during (early) inspiral.

Tidal effects and EOB formalism

- tidal effects are important in late inspiral of binary neutron stars

Flanagan, Hinderer 08, Hinderer et al 09, Damour, Nagar 09, Binnington, Poisson 09

$\longrightarrow$ a possible handle on the nuclear equation of state

- tidal extension of EOB formalism : non minimal worldline couplings

$$\Delta S_{\text{nonminimal}} = \sum_A \frac{1}{4}\,\mu_2^A \int ds_A (u^\mu u^\nu R_{\mu\alpha\nu\beta})^2 + \ldots$$

Damour, Esposito-Farèse 96, Goldberger, Rothstein 06, Damour, Nagar 09

$\longrightarrow$ modification of EOB effective metric + … :

$$\begin{aligned} A(r) &= A^0(r) + A^{\text{tidal}}(r) \\ A^{\text{tidal}}(r) &= -\kappa_2\, u^6 (1 + \bar{\alpha}_1\, u + \bar{\alpha}_2\, u^2 + \ldots) + \ldots \end{aligned}$$

- need accurate NR simulation to “calibrate” the higher-order PN contributions that are quite important during late inspiral

Uryu et al 06, 09, Rezzolla et al 09

Conclusions

- Any strategy of building GW templates cannot avoid **synergy between analytical and numerical relativity.**

- **Complementarity** between resummed perturbation theory (EOB), and nonperturbative Numerical Relativity results.

- The EOB formalism made several (qualitative and semi-quantitative) predictions that have been broadly confirmed by NR (e.g. J/M^2 (final) within 10%)

- The EOB formalism (in all its various avatars) can provide high-accuracy ***parameter free*** templates ***$h(m_1,m_2)$*** for GWs from BBH coalescence, with unprecedented agreement with NR data (**and for any mass ratio**).

- Tidal effects have been recently included (Neutron Stars)

- **Next challenges: - SPIN**

  - eccentric orbits (LISA)