

# Seven-dimensional forest fires

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## Abstract

We show that in high dimensional Bernoulli percolation, removing from a thin infinite cluster a much thinner infinite cluster leaves an infinite component. This observation has implications for the van den Berg-Brouwer forest fire process, also known as self-destructive percolation, for dimension high enough.

## 1 Introduction

Think about the open vertices of supercritical percolation as if they were trees, and about the infinite cluster as a forest. Suddenly a fire breaks out and the entire forest is cleared. New trees then start growing randomly. When can one expect a new infinite cluster to appear? The surprising conjecture in [vdBB04] is that in the two-dimensional case, even if the original forest were extremely thin, still a considerable amount of trees must be added to create a new infinite cluster. Heuristically, the conjecture claims that the infinite cluster might occupy a very low proportion of vertices but they sit in a way that separates the remaining finite clusters by gaps that cannot be easily bridged. This conjecture is still open. See [vdBB04, vdBBV08, vdBdL09] for connections to other models of forest fires and more.

Let us define the model formally, in three steps. The model was originally introduced as a site percolation model, but we will define it for bonds, as some of the auxiliary results we need have only been proved for bond percolation. We are given a graph  $G$ , a probability  $p \in [0, 1]$  (“the original density”) and a probability  $\varepsilon \in [0, 1]$  (“the recovered density”). Let  $\mathbb{P}_p$  be the Bernoulli bond percolation measure on  $G$  with parameter  $p$ .

1. Assign independent uniformly distributed values from  $[0, 1]$  to the edges of  $G$ . Let  $\omega_p \in \{0, 1\}^{E(G)}$  denote the set of edges with value at most  $p$ . The configuration  $\omega_p$  is distributed as  $\mathbb{P}_p$ , and a *cluster* refers

to a maximal connected component of edges. It will be of importance below that as  $p$  ranges over  $[0, 1]$ , we obtain a simultaneous coupling of Bernoulli configurations on  $G$  such that  $\omega_{p_1} \subset \omega_{p_2}$  when  $p_1 \leq p_2$ .

2. Let  $\tilde{\mathbb{P}}_p$  be the law of the configuration  $\tilde{\omega}_p$  constructed as follows: for any edge  $e$ ,

$$\tilde{\omega}_p(e) = \begin{cases} \omega_p(e) & \text{if } e \text{ is in a finite cluster of } \omega_p, \\ 0 & \text{otherwise.} \end{cases}$$

3. Let  $\tilde{\mathbb{P}}_{p,\varepsilon}$  be the law of  $\tilde{\omega}_{p,\varepsilon}$  where  $\tilde{\omega}_{p,\varepsilon}$  is defined as follows: for any edge  $e$ ,  $\tilde{\omega}_{p,\varepsilon}(e) = \max\{\tilde{\omega}_p(e), \omega'_\varepsilon(e)\}$ , where  $\omega'_\varepsilon$  is a percolation configuration with edge-weight  $\varepsilon$  which is independent of  $\omega_p$ .

We can now define our property of interest.

**Definition** We say that the graph  $G$  *recovers from fires* if for every  $\varepsilon > 0$ , there exists  $p > p_c(G)$  such that  $\tilde{\mathbb{P}}_{p,\varepsilon}$  has an infinite connected component, with probability 1. We say that  $G$  *site-recovers from fires* if the analogous definitions for site percolation hold.

In [vdBB04] the authors showed that a binary tree site-recovers from fires and conjectured that  $\mathbb{Z}^2$  lattice does *not* site-recover from fires. The binary tree is an example of a non-amenable graph, that is, a graph in which the boundary of a (finite) set of vertices is comparable in size to the set itself. Recovery from fires, both in edge and site sense, was proven in [AST13] for a large class of non-amenable transitive graphs. Our result concerns hyper-cubic lattices.

**Theorem 1.** *For  $d$  sufficiently large,  $\mathbb{Z}^d$  recovers from fires.*

Here and below,  $\mathbb{Z}^d$  refers to the  $\mathbb{Z}^d$  nearest neighbour lattice. The main property of  $\mathbb{Z}^d$  that we will use is that  $\mathbb{P}_{p_c}(0 \leftrightarrow \partial B(0,r)) \leq Cr^{-2}$  (see below for a discussion on this condition, and also for the notations). This was proved in [KN11] based on results of Hara, van der Hofstad & Slade [HvdHS03, Har08]. These establish the necessary estimate for  $d$  sufficiently large (19 seems to be enough, though this can be improved) and also for *stretched-out* lattices in  $d > 6$ . The number 6 is actually meaningful and is the limit of the technique involved, lace expansion. Our proof easily extends to stretched-out 7-dimensional lattice (hence the title of the article), but for simplicity we will prove the theorem only for nearest-neighbour percolation in  $d$  sufficiently high. In fact, our proof provides further information in the supercritical percolation regime. Recall the common notation  $\mathcal{C}_\infty(\omega_p)$  for the infinite cluster of edges present in  $\omega_p$ .

**Theorem 2.** *For every  $\varepsilon > 0$  and  $d$  sufficiently large, there exists  $p > p_c$  such that  $\omega_{p_c+\varepsilon} \setminus \mathcal{C}_\infty(\omega_p)$  contains an infinite cluster almost surely.*

Theorem 1 is clearly a corollary of Theorem 2. Another consequence is that for every  $\varepsilon > 0$ , the critical probability for percolation on the random graph obtained from  $\mathbb{Z}^d$  by removing a sufficiently ‘thin’ supercritical percolation cluster is almost surely at most  $p_c + \varepsilon$ . Theorem 2 and the last statement cannot possibly hold for (site) percolation on  $\mathbb{Z}^2$ , since an infinite cluster cuts space up into finite pieces.

**Proof sketch** We will show that for every  $\varepsilon > 0$ , there exists some  $p > p_c$  such that when removing the infinite cluster of  $p$ -percolation from  $(p_c + \varepsilon)$ -percolation, the remainder still percolates. The proof proceeds by a renormalization procedure.

1. We first choose  $\ell \in \mathbb{N}$  sufficiently large such that for any  $L \geq \ell$ , connectivity properties of boxes of size  $L^2 \times \ell^{d-2}$  in  $(p_c + \varepsilon)$ -percolation behave like  $(1 - \eta)$ -percolation on a coarse grain lattice for some small  $\eta$ . This is a standard application of Grimmett-Marstrand [GM90] and renormalization theory.
2. We then use the fact that the one-arm exponent in high dimensions is 2 to note that for any  $L$ , only a small number  $M$  of vertices in a box of size  $L^2 \times \ell^{d-2}$  can connect to distance  $L$  in *critical percolation*.
3. Picking  $L$  sufficiently large, one can argue that these  $M$  points do not alter the connectivity properties of boxes of size  $L^2 \times \ell^{d-2}$  for  $(p_c + \varepsilon)$ -percolation. In particular, the coarse grain percolation still behaves like  $(1 - \eta)$ -percolation even after removing that small number of vertices.
4. We now pick  $p$  sufficiently close to  $p_c$  that the behaviour (for  $\omega_p$ ) at scale  $L$  is not altered by moving from  $p_c$  to  $p$ . Since there are less sites in  $\mathcal{C}_\infty(\omega_p)$  than sites connected to distance  $L$  in  $\omega_p$ , this  $p$  gives the result.

Examining this a little shows that what the proof really needs is that the one-arm exponent is bigger than 1 i.e. that

$$\mathbb{P}_{p_c}(0 \longleftrightarrow \partial B(0, r)) \leq r^{-1-c} \quad c > 0.$$

(the number of points removed in the second renormalization step will no longer be bounded independently of  $L$ , but would still be too small to block the cluster of the boxes at scale  $\ell$ ). This is interesting as it is conjectured to hold also below 6 dimensions. While nothing is proved, simulations hint

that it might hold for  $\mathbb{Z}^5$  [AS94, §2.7]. On the other hand, let us note that in  $\mathbb{Z}^3$  this probability is larger than  $cr^{-1}$  (this is well-known but we are not aware of a precise reference – compare to [vdBK85, (3.15)] and [Kes82, Theorem 5.1]). Hence, the approach used here has no hope of working in  $\mathbb{Z}^3$  (though, of course, this does not preclude the possibility that  $\mathbb{Z}^3$  does recover from fires). We remark that a similar renormalization technique was recently used in [GHK13], also under the assumption that the one-arm exponent is bigger than 1.

**Notations** Identify  $\mathbb{Z}^2$  with the subgraph of  $\mathbb{Z}^d$  of points with the  $d-2$  last coordinates equal to 0. Let  $S_\ell = \{x \in \mathbb{Z}^d : |x_i| \leq \ell \ \forall i \geq 3\}$  be the two-dimensional slab of height  $2\ell + 1$ . We will also use the following standard notations: For a subgraph  $G$  of  $\mathbb{Z}^d$ , we say that  $x$  is connected to  $y$  in  $G$  if they are in the same connected component of  $G$ . We denote it by  $x \xleftrightarrow{G} y$  or simply  $x \longleftrightarrow y$  when the context is clear. We also use the notation  $x \longleftrightarrow \infty$  to denote the fact that  $x$  is contained in an infinite connected component.

Let  $\|\cdot\|_\infty$  be the infinity norm on  $\mathbb{R}^d$  defined by

$$\|x\|_\infty = \max\{|x_i| : i = 1, \dots, d\}.$$

We consider the hypercubic lattice  $\mathbb{Z}^d$  for some large but fixed  $d$ . For  $\ell, L > 0$ , define the ball  $B_x(L) = \{y \in \mathbb{Z}^d : \|y - x\|_\infty \leq L\}$  and let  $\partial B_x(L)$  be its inner vertex boundary.

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## 2 Proof

From now on,  $d$  is fixed and is large enough. For  $x \in \mathbb{Z}^2$ , let  $\mathcal{A}(x, \ell, L, M)$  be the event that there are less than  $M$  sites  $y$  in the  $(6L + 1) \times (6L + 1) \times (2\ell + 1)^{d-2}$  box  $S_\ell \cap B_x(3L)$  that are connected to a site at distance  $L$  from themselves. Note that we do not assume that this connection is contained in the slab  $S_\ell$ , the connection may be anywhere in  $B_y(L)$ .

**Lemma 3.** *Let  $\eta > 0$  and  $\ell > 0$ . There exists  $M > 0$  such that for any integer  $L$ , there exists  $p > p_c$  such that*

$$\mathbb{P}_p(\mathcal{A}(x, \ell, L, M)) \geq 1 - \eta.$$

*Proof.* By [KN11], there exists  $C > 0$  such that (for large enough  $d$ )

$$\mathbb{P}_{p_c}(0 \longleftrightarrow \partial B_0(n)) \leq \frac{C}{n^2}. \quad (1)$$

Choose  $M$  in such a way that  $\frac{49(2\ell+1)^{d-2}C}{M} < \eta$ . For any integer  $L$ , Markov's inequality implies

$$\begin{aligned} \mathbb{P}_{p_c} \left[ |\{y \in S_\ell \cap B_x(3L) : y \leftrightarrow \partial B_y(L)\}| \geq M \right] \\ \leq \frac{1}{M} \sum_{y \in S_\ell \cap B_x(3L)} \mathbb{P}_{p_c}(y \leftrightarrow \partial B_y(L)). \end{aligned}$$

By (1) and the choice of  $M$ , the right-hand side is thus strictly smaller than  $\eta$ . By choosing  $p$  close enough to  $p_c$ , we obtain that

$$\mathbb{P}_p \left[ |\{y \in S_\ell \cap B_x(3L) : y \leftrightarrow \partial B_y(L)\}| \geq M \right] \leq \eta. \quad \square$$

For a set  $S \subset \mathbb{Z}^d$ , let  $\omega^S$  be the configuration obtained from  $\omega$  by closing the edges adjacent to a site in  $S$ . Let  $\mathcal{B}(x, \ell, L, M)$  be the event that for any set  $S$  of  $M$  sites contained in  $B_x(3L)$ ,  $\omega^S$  contains

- a cluster crossing from  $\partial B_x(L)$  to  $\partial B_x(3L)$  contained in the slab  $S_\ell$ ,
- a unique cluster in the box  $S_\ell \cap B_x(3L)$  of radius larger than  $L$ .

**Lemma 4.** *Let  $\eta > 0$  and  $\varepsilon > 0$ . There exists  $\ell > 0$  such that for any  $M > 0$ , there is  $L > 0$  so that*

$$\mathbb{P}_{p_c+\varepsilon}(\mathcal{B}(x, \ell, L, M)) \geq 1 - \eta.$$

*Proof.* For a given  $\ell$  and  $L$  denote by  $E = E(x, \ell, L)$  the event that:

1. There is a crossing from  $\partial B_x(L)$  to  $\partial B_x(3L)$  in  $S_\ell$ .
2. There is exactly one cluster in  $S_\ell \cap B_x(3L)$  of radius larger than  $L$ .

Shortly, the event  $E$  is just  $\mathcal{B}$  without the set  $S$ , or if you want  $\mathcal{B}$  is the event that  $E$  occurred in  $\omega^S$  for all  $S$  with  $|S| \leq M$ .

We claim that for  $\ell$  sufficiently large,  $\mathbb{P}_{p_c+\varepsilon}(\neg E) \leq \exp(-cL)$  for some  $c = c(\varepsilon, \ell) > 0$  independent of  $L$ . Finding such an  $\ell$  is a standard exercise in renormalization theory, but let us give a few details nonetheless. Call a box of side-length  $2\ell+1$  *good* if it contains crossings between opposite faces in all directions, and if all clusters of diameter at least  $\frac{1}{4}\ell$  connect inside the box. By choosing  $\ell$  large, we can require that a box is good with arbitrarily high probability (see e.g. the appendix of [BBHK08]). Considering such boxes centered around the sites in  $\ell\mathbb{Z}^2$ . The events that these boxes are good are 2-dependent (in the sense of [LSS97] i.e. any box is independent

of all boxes not neighbouring it), and hence by [LSS97], if the probability that a box is good is sufficiently large, then the good boxes stochastically dominate two-dimensional percolation at density, say,  $\frac{9}{10}$ . Now, a cluster of good boxes contains a cluster in the underlying percolation, since the crossings of adjacent boxes must intersect. This means that if either of the conditions in the definition of  $E$  fail, then there is a cluster of bad boxes with at least  $L/\ell$  boxes. But the probability for that, from Peierls' argument, is at most  $(4/10)^{L/\ell} \cdot (6L/\ell)^2$ . This shows the claim.

Fix  $M > 0$ . Let  $F_M$  be the set of configurations in  $B_x(3L)$  for which there exists  $S \subset B_x(3L)$  with  $|S| = M$  and  $\omega^S \notin E$ . We have

$$\begin{aligned} \mathbb{P}_{p_c+\varepsilon}(F_M) &\leq \sum_{S \subset B_x(3L): |S|=M} \mathbb{P}_{p_c+\varepsilon}(\omega^S \notin E) \\ &\leq \sum_{S \subset B_x(3L): |S|=M} (1-p_c-\varepsilon)^{-2dM} \mathbb{P}_{p_c+\varepsilon}(-E) \\ &\leq (1-p_c-\varepsilon)^{-2dM} (6L+1)^{dM} \mathbb{P}_{p_c+\varepsilon}(-E) \\ &\leq (1-p_c-\varepsilon)^{-2dM} (6L+1)^{dM} \exp(-cL). \end{aligned}$$

For  $L$  large enough, this quantity is smaller than  $\eta$ . The lemma follows from the fact that if  $\omega \notin \mathcal{B}(x, \ell, L, M)$ , then there exists  $S \subset B_x(3L)$  with  $|S| = M$  and  $\omega^S \notin E$ , i.e.  $\omega \in F_M$ .  $\square$

In order to prove Theorem 1 and 2, we will use Lemma 4 to construct an infinite cluster at density  $p_c + \varepsilon$ , and Lemma 3 to make sure that the infinite cluster present at the lower density  $p$  does not interfere too much with this construction.

*Proof of Theorems 1 and 2.* Recall the notations  $\omega_p$ ,  $\tilde{\omega}_p$  and  $\omega'_\varepsilon$  from page 1. We need to show that for any  $\varepsilon > 0$ , there exists  $p > p_c$  such that  $\tilde{\omega}_{p,\varepsilon}$  has an infinite component. Note that  $(\omega_{p_c} \cup \omega'_\varepsilon) \setminus \mathcal{C}_\infty(\omega_p)$  is stochastically dominated by  $\tilde{\omega}_{p,\varepsilon}$ . Thus, it suffices to show that for every  $\varepsilon > 0$ , there is  $p > p_c$  such that  $\omega_{p_c+\varepsilon} \setminus \mathcal{C}_\infty(\omega_p)$  contains an infinite component. That is, Theorem 1 follows from Theorem 2, and it suffices to prove the latter.

Let therefore  $\varepsilon > 0$ . Fix  $\eta > 0$  such that  $1-2\eta$  exceeds the critical parameter for any 8-dependent percolation on vertices of  $\mathbb{Z}^2$ . Define successively  $\ell, M, L$  and  $p_0$  as follows. Fix  $\ell = \ell(\varepsilon, \eta) > 0$  as defined in Lemma 4. Pick  $M = M(\eta, \ell) > 0$  as defined in Lemma 3. This defines  $L = L(\eta, \varepsilon, \ell, M) > 0$  by Lemma 4, and then  $p = p(\eta, \ell, M, L) > p_c$  by Lemma 3.

Let  $\mathbf{P}$  denote the joint law of  $(\omega_p, \omega_{p_c+\varepsilon})$  under the increasing coupling described above. A site  $x \in L\mathbb{Z}^2$  is said to be *good* if  $\omega_p \in \mathcal{A}(x, \ell, L, M)$  and  $\omega_{p_c+\varepsilon} \in \mathcal{B}(x, \ell, L, M)$ . By definition,

$$\mathbf{P}[\mathcal{A}(x, \ell, L, M) \cap \mathcal{B}(x, \ell, L, M)] \geq 1 - 2\eta.$$

Since these events depend on edges in  $B_x(4L)$  only, the site percolation (on  $L\mathbb{Z}^2$ ) thus obtained is 8-dependent. As a consequence, there exists an infinite cluster of good sites on the coarse grained lattice  $L\mathbb{Z}^2$ .

On the event that there exists an infinite cluster of good sites on the coarse grained lattice, there exists an infinite path in  $\omega_{p_c+\varepsilon} \setminus \mathcal{C}_\infty(\omega_p)$ . Indeed, by induction, consider a path of adjacent good sites  $x_1, \dots, x_n$ . Consider  $C_i$  to be a cluster in

$$[\omega_{p_c+\varepsilon} \setminus \mathcal{C}_\infty(\omega_p)] \cap [B_{x_i}(3L) \setminus B_{x_i}(L)]$$

of radius larger than  $L$ . By the definition of  $\mathcal{A}$  there are at most  $M$  sites in the box  $S_l \cap B_{x_i}(3L)$  connected to distance  $L$  in  $\omega_p$ . Hence the same box also contains no more than  $M$  sites in  $\mathcal{C}_\infty(\omega_p)$  since any site connected to infinity must be connected to distance  $L$ . Using the definition of  $\mathcal{B}$  with  $S$  being exactly  $\mathcal{C}_\infty(\omega_p) \cap S_l \cap B_{x_i}(3L)$  we see that  $\omega_{p_c+\varepsilon} \setminus \mathcal{C}_\infty(\omega_p)$  contains a crossing cluster for the box  $S_l \cap B_{x_i}(3L)$  with all the properties listed before Lemma 4. In particular, the uniqueness property ensures two such crossing clusters in two neighbouring boxes must intersect. The result follows readily.  $\square$

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