

Higher genus modular graph functions

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Introduction

- **String theory naturally generalizes real-analytic Eisenstein series**
 - for genus one surfaces,
 - ★ multiple Kronecker-Eisenstein sums,
 - ★ multiple integrations of Green function on the torus,
 - ⇒ modular graph functions invariant under $SL(2, \mathbb{Z})$.

(Green, Russo, Vanhove 2008; ED, Green, Vanhove 2015; ED, Green, Gurdogan, Vanhove 2015)

- **String theory includes contributions from surfaces of all genera**
 - ⇒ **expect modular graph functions for higher genus surfaces.**
- **Focus of this talk is on genus two and higher**
 - simplest is Kawazumi-Zhang invariant (to be explained below);
 - genus-two string theory predicts an infinite number of higher invariants;
 - genus greater than two string theory offers no predictions,
 - but mathematical constructions produce higher invariants.
- **I will give an account of what we know and do not know to date.**

Bibliography

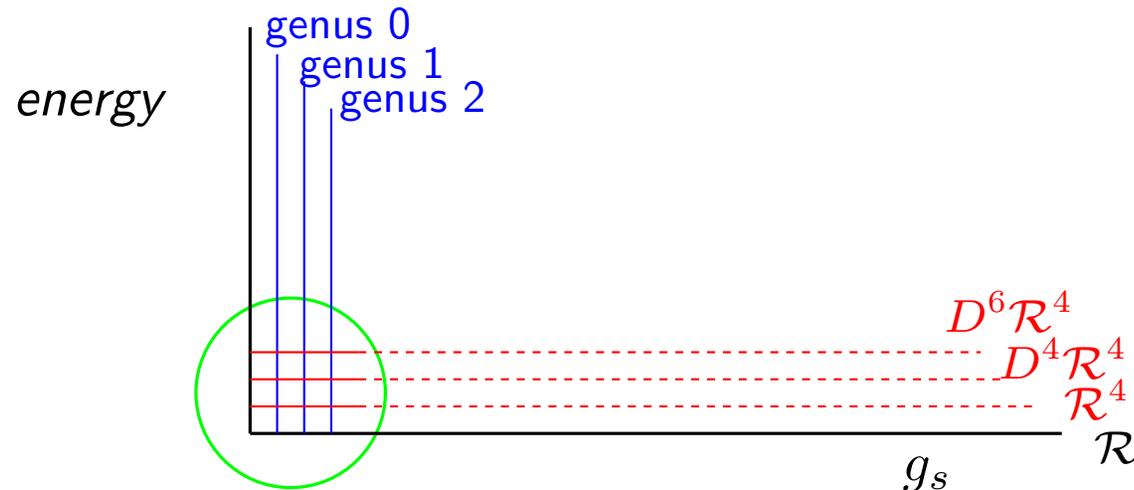
Based on

- ED, Michael Green and Boris Pioline, arXiv:1712.06135,
Higher genus modular graph functions, string invariants, and their exact asymptotics
- ED, Michael Green and Boris Pioline, (in preparation)
Asymptotics of the $D^8\mathcal{R}^4$ genus-two string invariant

and earlier work

- ED, Michael Green, arXiv:1308.4597, Journal of Number Theory, Vol 144 (2014) 111-150,
Zhang-Kawazumi invariants and Superstring Amplitudes
- ED, Michael Green, Boris Pioline, Rudolfo Russo, arXiv:1405.6226, JHEP 1501 (2015) 031,
Matching the $D^6\mathcal{R}^4$ interaction at two-loops
- Boris Pioline, arXiv:1504.04182, Journal of Number Theor. **163**, 520 (2016),
A Theta lift representation for the Kawazumi-Zhang and Faltings invariants of genus-two Riemann surfaces

Expansion in genus and low energy



- **Superstring Perturbation Theory** in $g_s^{(2-2h)}$ with $h \geq 0$ is the genus
 - holds for small string coupling $g_s \ll 1$
 - but for all energies
- **Supergravity** \mathcal{R}
 - leading low energy expansion of string theory
 - holds for all couplings g_s
- **String induced effective interactions** $\mathcal{R}^4, D^4\mathcal{R}^4, D^6\mathcal{R}^4$
 - Evaluated in superstring perturbation theory
 - Accessible via the four-graviton scattering amplitude

Effective Interactions

- Four-graviton amplitude in Type II at genus 0,

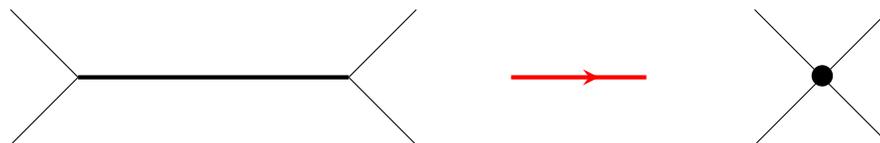
$$\mathcal{A}^{(0)}(s_{ij}) = \mathcal{R}^4 \frac{1}{stu} \frac{\Gamma(1-s)\Gamma(1-t)\Gamma(1-u)}{\Gamma(1+s)\Gamma(1+t)\Gamma(1+u)}$$

- \mathcal{R}^4 = unique maximally supersymmetric contraction of 4 Weyl tensors
 - External momenta k_i for $i = 1, 2, 3, 4$ with $k_i^2 = 0$ and $\sum_i k_i = 0$
 - Introduce dimensionless Lorentz-invariants $s_{ij} = -\alpha' k_i \cdot k_j / 2$
 - $s = s_{12} = s_{34}$, $t = s_{13} = s_{24}$, $u = s_{14} = s_{23}$ with $s + t + u = 0$
- Low energy expansion corresponds to $|s|, |t|, |u| \ll 1$

$$\frac{1}{stu} + 2\zeta(3) + \zeta(5)(s^2 + t^2 + u^2) + 2\zeta(3)^2 stu + \dots$$

massless
 \mathcal{R}^4
 $D^4\mathcal{R}^4$
 $D^6\mathcal{R}^4$

Exchange of massive string states produces local effective interactions.



Genus-one string amplitude

- **Effective \mathcal{R}^4 -type interactions in Type II**

- Generated by a multiple integral over a torus $\Sigma_1 = \mathbb{C}/(\mathbb{Z} + \tau\mathbb{Z})$, of modulus $\tau \in \mathbb{H}$, namely $\tau = \tau_1 + i\tau_2$ with $\tau_1, \tau_2 \in \mathbb{R}$ and $\tau_2 > 0$,

$$\mathcal{B}^{(1)}(s_{ij}|\tau) = \prod_{i=1}^N \int_{\Sigma_1} \frac{d^2 z_i}{\tau_2} \exp \left\{ \sum_{1 \leq i < j \leq N} s_{ij} g(z_i - z_j|\tau) \right\}$$

- **Mathematically, one may consider this integral for arbitrary N**

- $g(z|\tau)$ is the translation-invariant Green function on Σ_1 ,

$$\tau_2 \partial_{\bar{z}} \partial_z g(z|\tau) = -\pi \delta^{(2)}(z) + \pi \quad \int_{\Sigma_1} d^2 z g(z|\tau) = 0$$

- Integrals absolutely convergent for $|s_{ij}| < 1$; analytic near $s_{ij} = 0$;
- $\mathcal{B}^{(1)}(s_{ij}|\tau)$ is invariant under the modular group $SL(2, \mathbb{Z})$.

- **String amplitude obtained by integral over modulus of the torus,**

$$\mathcal{A}^{(1)}(s_{ij}) = \int_{\mathbb{H}/SL(2, \mathbb{Z})} \frac{d^2 \tau}{\tau_2^2} \mathcal{B}^{(1)}(s_{ij}|\tau)$$

- requires analytic continuation in s_{ij} (ED, Phong 1994).

Genus-one modular graph functions

- Taylor series expansion of $\mathcal{B}^{(1)}(s_{ij}|\tau)$ for fixed τ in powers of s_{ij}

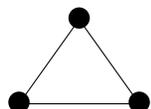
– An integration point z_i is represented by a vertex \bullet

– A Green function is represented by an edge $\bullet \text{---} \bullet = g(z_i - z_j|\tau)$

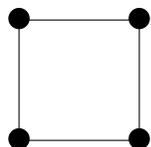
$D^4 \mathcal{R}^4$



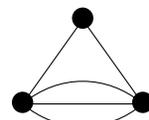
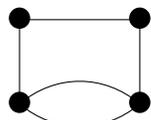
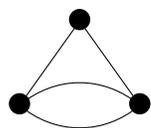
$D^6 \mathcal{R}^4$



$D^8 \mathcal{R}^4$



$D^{10} \mathcal{R}^4$



one-loop

two-loops

three-loops

four-loops

Properties of genus-one modular graph functions

- One-loop graphs with k vertices give real analytic Eisenstein series E_s

$$\prod_{i=1}^k \int_{\Sigma} \frac{d^2 z_i}{\tau_2} g(z_i - z_{i+1} | \tau) = \sum_{(m,n) \neq (0,0)} \frac{\tau_2^k}{\pi^k |m\tau + n|^{2k}} = E_k$$

- convergent sums for $\text{Re}(s) > 1$; modular $SL(2, \mathbb{Z})$ -invariant
- Laplace-eigenvalue equation, $(\Delta - s(s-1)) E_s = 0$ with $\Delta = 4\tau_2^2 \partial_\tau \partial_{\bar{\tau}}$

- Two-loop graphs evaluate to the series

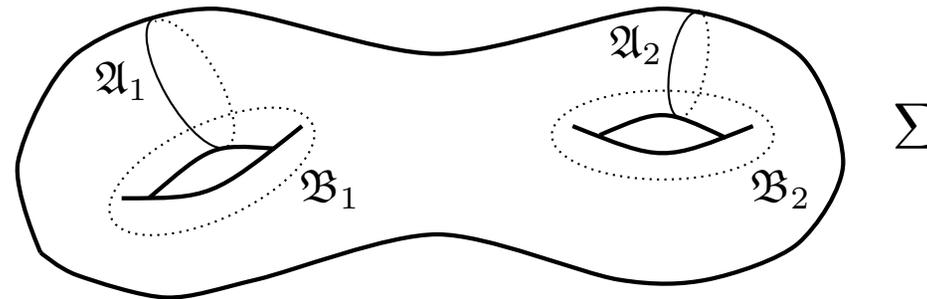
$$C_{s_1, s_2, s_3}(\tau) = \sum_{(m_r, n_r) \neq (0,0)} \prod_{r=1}^3 \left(\frac{\tau_2}{\pi |m_r \tau + n_r|^2} \right)^{s_r} \delta\left(\sum_r m_r\right) \delta\left(\sum_r n_r\right)$$

- convergent sums for $\text{Re}(s_r) \geq 1$; modular $SL(2, \mathbb{Z})$ -invariant;
- satisfy inhomogeneous Laplace-eigenvalue equations, e.g.

$$\begin{aligned} \Delta C_{1,1,1} &= 6E_3 \\ (\Delta - 2)C_{2,1,1} &= 9E_4 - E_2^2 \\ (\Delta - 6)C_{3,1,1} &= 3C_{2,2,1} + 16E_5 - 4E_2E_3 \end{aligned} \tag{1}$$

Genus-two surfaces

- Σ is a compact Riemann surface of genus two
 - Key difference with genus-one: no translation symmetry



- Homology and cohomology

- One-cycles $H_1(\Sigma, \mathbb{Z}) \approx \mathbb{Z}^4$ with intersection pairing $\mathfrak{I}(\cdot, \cdot) \rightarrow \mathbb{Z}$
- Canonical basis $\mathfrak{I}(\mathfrak{A}_I, \mathfrak{A}_J) = 0$, $\mathfrak{I}(\mathfrak{B}_I, \mathfrak{B}_J) = 0$ with $I, J = 1, 2$
 $\mathfrak{I}(\mathfrak{A}_I, \mathfrak{B}_J) = \delta_{IJ}$, $\mathfrak{I}(\mathfrak{B}_I, \mathfrak{A}_J) = -\delta_{IJ}$
- Canonical dual basis of holomorphic one-forms ω_I in $H^{(1,0)}(\Sigma)$

$$\oint_{\mathfrak{A}_I} \omega_J = \delta_{IJ} \qquad \oint_{\mathfrak{B}_I} \omega_J = \Omega_{IJ}$$

- Period matrix Ω obeys Riemann relations $\Omega^t = -\Omega$, $\text{Im}(\Omega) > 0$

Modular transformations and geometry

- Transformation $Sp(4, \mathbb{Z}) : H_1(\Sigma, \mathbb{Z}) \rightarrow H_1(\Sigma, \mathbb{Z})$ leaves $\mathfrak{J}(\cdot, \cdot)$ invariant
 - action on basis cycles given by

$$\begin{pmatrix} \mathfrak{B}_I \\ \mathfrak{A}_I \end{pmatrix} \rightarrow \sum_J M_{IJ} \begin{pmatrix} \mathfrak{B}_J \\ \mathfrak{A}_J \end{pmatrix} \quad M^t \mathfrak{J} M = \mathfrak{J}$$

- action on 1-forms ω_I and periods Ω_{IJ} given by

$$\begin{aligned} \omega &\rightarrow \omega (C\Omega + D)^{-1} \\ \Omega &\rightarrow (A\Omega + B) (C\Omega + D)^{-1} \end{aligned} \quad M = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$$

- Siegel upper half space \mathcal{S}_2

$$\mathcal{S}_2 = \left\{ \Omega_{IJ} \in \mathbb{C} \text{ with } \Omega^t = \Omega \text{ and } Y = \text{Im}(\Omega) > 0 \right\}$$

- $\mathcal{S}_2 = \frac{Sp(4, \mathbb{R})}{SU(2) \times U(1)} = \frac{SO(3, 2)}{SO(3) \times SO(2)}$ is Kähler with invariant metric

$$ds_2^2 = \sum_{I, J, K, L=1, 2} Y_{IJ}^{-1} d\bar{\Omega}_{JK} Y_{KL}^{-1} d\Omega_{LI}$$

- Moduli space of genus-two surfaces is $\mathcal{S}_2 / Sp(4, \mathbb{Z})$ (minus diagonal Ω)

Green function and volume form

- How to generalize the genus-one formula to a genus-two formula ?
 - recall the genus-one formula

$$\mathcal{B}^{(1)}(s_{ij}|\tau) = \prod_{i=1}^N \int_{\Sigma_1} \frac{d^2 z_i}{\tau_2} \exp \left\{ \sum_{1 \leq i < j \leq N} s_{ij} g(z_i - z_j|\tau) \right\}$$

- Canonical metric and Kähler form for genus-two Σ
 - modular invariant and smooth

$$\kappa = \frac{i}{4} \sum_{I,J} Y_{IJ}^{-1} \omega_I \wedge \bar{\omega}_J \quad \int_{\Sigma} \kappa = 1$$

- Natural “Arakelov Green function” $\mathcal{G}(w, z|\Omega) = \mathcal{G}(z, w|\Omega)$
 - Inverse of scalar Laplace operator for canonical metric

$$\begin{aligned} \partial_{\bar{w}} \partial_w \mathcal{G}(w, z|\Omega) &= -\pi \delta(w, z) + \pi \kappa(w) \\ \int_{\Sigma} \kappa(w) \mathcal{G}(w, z|\Omega) &= 0 \end{aligned}$$

A natural genus-two candidate

- A natural candidate formula for a string amplitude would be

$$\mathcal{C}^{(2)}(s_{ij}|\Omega) = \prod_{i=1}^N \int_{\Sigma} \kappa(z_i) \exp \left\{ \sum_{1 \leq i < j \leq N} s_{ij} \mathcal{G}(z_i, z_j|\Omega) \right\}$$

- Integrals absolutely convergent for $|s_{ij}| < 1$; analytic near $s_{ij} = 0$,
- Expanding in powers of s_{ij} gives *genus-two modular graph functions*.
- **But ... Genus-two string amplitudes are NOT given by $\mathcal{C}^{(2)}(s_{ij}|\Omega)$**
- For integration over a single copy of Σ
 - κ is the only natural modular invariant volume form.
- For integration over multiple copies of Σ
 - $Sp(4, \mathbb{Z})$ modular invariants other than $\prod_i \kappa(z_i)$ allowed.
 - For example, when $N = 2$ we can have $\kappa(z_1)\kappa(z_2)$ as well as

$$\sum_{I,J,K,L} Y_{IL}^{-1} Y_{JK}^{-1} \omega_I(z_1) \overline{\omega_J(z_1)} \omega_K(z_2) \overline{\omega_L(z_2)}$$

Genus-two string amplitude

- Instead the $N = 4$ graviton amplitude was calculated (ED, Phong 2005)

$$\mathcal{B}^{(2)}(s_{ij}|\Omega) = \int_{\Sigma^4} \frac{\mathcal{Y} \wedge \bar{\mathcal{Y}}}{(\det Y)^2} \exp \left\{ \sum_{1 \leq i < j \leq 4} s_{ij} \mathcal{G}(z_i, z_j) \right\}$$

- The key difference with the candidate $\mathcal{C}^{(2)}$ is the structure of \mathcal{Y}

$$\begin{aligned} 3\mathcal{Y} &= (t - u) \Delta(z_1, z_2) \wedge \Delta(z_3, z_4) & s &= s_{12} = s_{34} \\ &+ (s - t) \Delta(z_1, z_3) \wedge \Delta(z_4, z_2) & t &= s_{13} = s_{24} \\ &+ (u - s) \Delta(z_1, z_4) \wedge \Delta(z_2, z_3) & u &= s_{14} = s_{23} \end{aligned}$$

- where Δ is a holomorphic $(1, 0)_i \times (1, 0)_j$ form on $\Sigma \times \Sigma$

$$\Delta(z_i, z_j) = \varepsilon^{IJ} \omega_I(z_i) \wedge \omega_J(z_j)$$

- The combination $\mathcal{Y} \wedge \bar{\mathcal{Y}} / (\det Y)^2$ is $Sp(4, \mathbb{Z})$ -invariant,
- and produces a modular invariant $\mathcal{B}^{(2)}(s_{ij}|\Omega)$.

Low energy expansion

- **Contributions to local effective interactions**

- Expand $\mathcal{B}^{(2)}(s_{ij}|\Omega)$ in powers of s_{ij} and integrate over $\mathcal{M}_2 = \mathcal{S}_2/Sp(4, \mathbb{Z})$

$\mathcal{R}^4, D^2\mathcal{R}^4$ zero, since \mathcal{Y} vanishes for $s = t = u = 0$
 $D^4\mathcal{R}^4$ Siegel volume form on \mathcal{M}_2
 $D^6\mathcal{R}^4$ one factor of \mathcal{G} in expansion in powers of s_{ij}

$$\mathcal{B}^{(2)}(s_{ij}|\Omega) = 32(s^2 + t^2 + u^2) + 192 stu \varphi(\Omega) + \mathcal{O}(s_{ij}^4)$$

$$\varphi(\Omega) = -\frac{1}{4} \sum_{I,J,K,L} Y_{IL}^{-1} Y_{JK}^{-1} \int_{\Sigma^2} \mathcal{G}(x, y) \omega_I(x) \overline{\omega_J(x)} \omega_K(y) \overline{\omega_L(y)}$$

- $\varphi(\Omega)$ coincides with the **Kawazumi-Zhang invariant** (ED, Green 2013)

- introduced as a spectral invariant (Kawazumi 0801.4218 and Zhang 0812.0371)
- related to the genus-two Faltings invariant (De Jong 2010)
- formulated in terms of modular tensors (Kawazumi OIST lecture notes 2016)

$$\mathcal{A}_{IJ;KL} = \int_{\Sigma^2} \mathcal{G}(x, y) \omega_I(x) \overline{\omega_J(x)} \omega_K(y) \overline{\omega_L(y)}$$

Higher string-invariants

- **The KZ-invariant exists for all genera ≥ 2** (Kawazumi 2008; Zhang 2008)
 - unknown if correct object for string theory at genus ≥ 3 .
- **But the Taylor expansion coefficients of $\mathcal{B}^{(2)}(s_{ij}|\Omega)$**
 - are *modular graph functions at genus-two*;
 - do naturally emerge from string theory at genus-two;
 - provide a string-motivated generalization of KZ-invariants at genus-two.
- **Higher string-invariants** (contribute to $D^8\mathcal{R}^4$ and $D^{10}\mathcal{R}^4$) (ED, Green 2013)

$$\mathcal{B}_{(2,0)}^{(2)} = \int_{\Sigma^4} \frac{|\Delta(1,2)\Delta(3,4)|^2}{(\det Y)^2} \left(\mathcal{G}(1,4) + \mathcal{G}(2,3) - \mathcal{G}(1,3) - \mathcal{G}(2,4) \right)^2$$

$$\mathcal{B}_{(1,1)}^{(2)} = \int_{\Sigma^4} \frac{|\Delta(1,2)\Delta(3,4) - \Delta(1,4)\Delta(2,3)|^2}{(\det Y)^2} \left(\mathcal{G}(1,2) + \mathcal{G}(3,4) \right. \\ \left. + \mathcal{G}(1,4) + \mathcal{G}(2,3) - 2\mathcal{G}(1,3) - 2\mathcal{G}(2,4) \right)^3$$

$$\dots = \dots$$

Differential equations and Asymptotics

- **Genus-one modular graph functions satisfy**
 - System of inhomogeneous Laplace-eigenvalue equations
 - Laurent polynomial behavior near the cusp of degree bounded by weight
- **Genus-two KZ-invariant φ satisfies**
 - Characteristic class relations and role in Johnson homomorphism
(Kawazumi 2008; Zhang 2008; DeJong 2010)
 - Eigenvalue equations for $Sp(4, \mathbb{R})$ -invariant differential operators
(ED, Green, Pioline, Russo 2014; Pioline 2015)
 - Theta-lift analogous to Borchers for Igusa cusp form (Pioline 2015)
 - Laurent polynomial of degree $(1, 1)$ near non-separating divisor
(De Jong 2010; Pioline 2015; ED, Green, Pioline 2017)
- **Genus-two higher string-invariants satisfy**
 - Laurent polynomial of bounded degree near non-separating degeneration
(ED, Green, Pioline 2017)
 - System of differential equations for invariant differential operators
(in progress ED, Green, Pioline 201?)

Differential equations

- **Theorem 1** *Laplace eigenvalue equation* (ED, Green, Pioline, Russo 2014)

$$(\Delta - 5)\varphi = -2\pi\delta_{SN}$$

- δ_{SN} has support on separating node (into two genus-one surfaces)
- Δ is the Laplace-Beltrami operator on \mathcal{S}_2 with Siegel metric

$$\Delta = 4 \sum_{I,J,K,L} Y_{IK} Y_{JL} \bar{\partial}^{IJ} \partial^{KL} \quad \partial^{IJ} = \frac{1}{2}(1 + \delta^{IJ}) \frac{\partial}{\partial \Omega_{IJ}}$$

- proven by theory and methods of deformations of complex structures

- **Theorem 2** *Quartic differential operator eigenvalue equation* (Pioline 2015)

$$(32 \square_2^* \square_0 - 15)\varphi = 0$$

- \square_0, \square_2 are Maass-Siegel operators

$$\square_0 = \varepsilon^{IJ} \varepsilon^{KL} \nabla_{IK} \nabla_{JL} \quad \nabla_{IJ} = Y_{IK} Y_{JL} \partial^{KL}$$

- **System of differential equations satisfied by higher string invariants ?**

Asymptotics

For genus one, there is only one type of degeneration, as modulus $\tau \rightarrow i\infty$

- Behavior of holomorphic Eisenstein series

$$\mathbb{G}_{2k}(\tau) = -\frac{B_{2k}}{2k} + \sum_{n=1}^{\infty} \sigma_{2k-1}(n)q^n \quad q = e^{2\pi i\tau}$$

- Behavior of non-holomorphic Eisenstein series, $y = \pi\text{Im}(\tau)$

$$E_k(\tau) = \frac{2\zeta(2k)}{\pi^{2k}}y^k + 4 \binom{2k-3}{k-1} \frac{\zeta(2k-1)}{(4y)^{k-1}} + \mathcal{O}(|q|)$$

- Behavior of general genus-one modular graph functions.
 - Finite degree Laurent polynomial in y plus exponentials,
 - eg two-loop modular graph functions (ED, Bill Duke 2017)

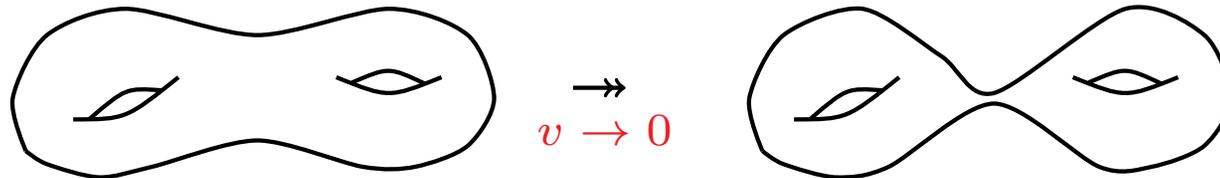
$$C_{a,b,c}(\tau) = c_w y^w + \frac{c_{2-w}}{y^{w-2}} + \sum_{k=1}^{w-1} \frac{c_{w-2k-1} \zeta(2k+1)}{y^{2k+1-w}} + \mathcal{O}(|q|)$$

$$c_{2-w} = \sum_{m=1}^{w-2} \gamma_m \zeta(2m+1) \zeta(2w-2m-3)$$

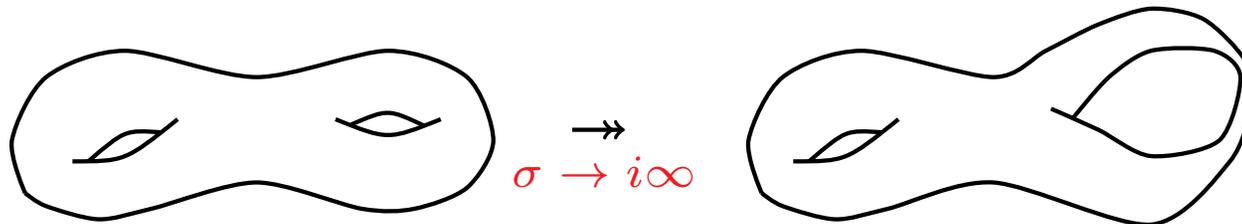
- with $w = a + b + c$ and $c_w, \gamma_m, c_{w-2k-1} \in \mathbb{Q}$.

Degenerations of genus-two Riemann surfaces $\Omega = \begin{pmatrix} \tau & v \\ v & \sigma \end{pmatrix}$

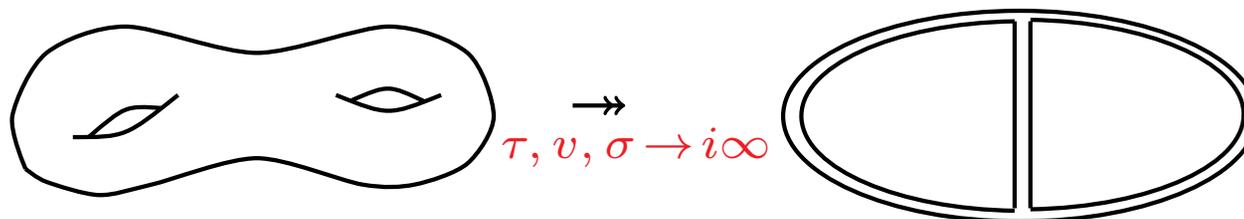
– Separating degeneration



– Non-separating degeneration



– Maximal degeneration (or “tropical limit”)



Non-separating degeneration

- A genus-two surface Σ degenerates to a torus Σ_1 with two punctures p_a, p_b
 - keep the cycles $\mathfrak{A}_1, \mathfrak{B}_1, \mathfrak{A}_2$ fixed, and let $\mathfrak{B}_2 \rightarrow \infty$

$$\Omega = \begin{pmatrix} \tau & v \\ v & \sigma \end{pmatrix} \quad \text{Im}(\sigma) \rightarrow \infty$$

- τ is the modulus of Σ_1 and $v = \int_{p_a}^{p_b} \omega_1 = p_b - p_a$

- The genus-two $Sp(4, \mathbb{Z})$ restricts to $SL(2, \mathbb{Z}) \times \mathbb{Z}^3$ (Fourier-Jacobi group)
 - $SL(2, \mathbb{Z})$ subgroup of $M \in Sp(4, \mathbb{Z})$ such that $M\mathfrak{B}_2 = \mathfrak{B}_2$ is

$$M = \begin{pmatrix} a & 0 & b & 0 \\ 0 & 1 & 0 & 0 \\ c & 0 & d & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \quad \begin{cases} \tau' & = & (a\tau + b)/(c\tau + d) \\ v' & = & v/(c\tau + d) \\ \sigma' & = & \sigma - cv^2/(c\tau + d) \end{cases}$$

- \mathbb{Z}^3 subgroup shifts v by $\mathbb{Z} + \tau\mathbb{Z}$ and $\text{Re}(\sigma)$ by \mathbb{Z} ;
- The degeneration parameter σ is not invariant under $SL(2, \mathbb{Z})$

Degeneration of Siegel modular forms

- A genus-two (“rank two”) Siegel modular form S of modular weight k
 - is holomorphic and transforms under $Sp(4, \mathbb{Z})$ by,

$$S(\Omega') = \det(C\Omega + D)^k S(\Omega) \quad \Omega' = (A\Omega + B)(C\Omega + D)^{-1}$$

- Fourier expansion in powers of $e^{2\pi i\sigma}$ near non-separating node

$$S(\Omega) = \sum_{m=0}^{\infty} e^{2\pi im\sigma} \phi_m(v|\tau)$$

- $\phi_m(v|\tau)$ is a *Jacobi form* of modular weight k and index m
- transforms under the residual modular group $SL(2, \mathbb{Z})$ by

$$\phi_m \left(\frac{v}{c\tau + d} \middle| \frac{a\tau + b}{c\tau + d} \right) = (c\tau + d)^k e^{2\pi imcv^2/(c\tau + d)} \phi_m(v|\tau)$$

(Eichler and Zagier, “The theory of Jacobi forms”, 1985)

Degeneration of the KZ-invariant

- **Non-separating degeneration of non-holomorphic modular functions**
 - is governed by a *real* $SL(2, \mathbb{Z})$ -invariant parameter $t > 0$

$$t \equiv \frac{\det(\operatorname{Im} \Omega)}{\operatorname{Im} \tau} \quad \Omega = \begin{pmatrix} \tau & v \\ v & \sigma \end{pmatrix}$$

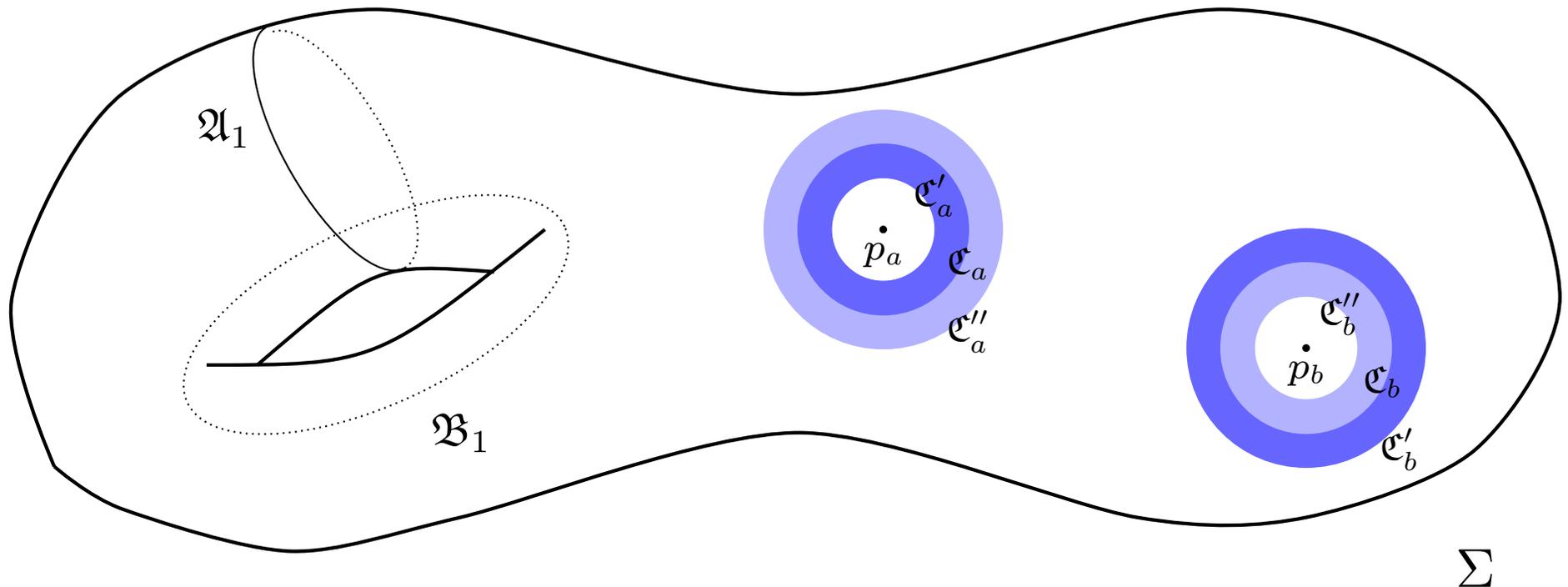
- with the non-separating node characterized by $t \rightarrow \infty$
- **Theorem 3** Non-separating degeneration of KZ-invariant (Pioline 2015)

$$\varphi(\Omega) = \frac{\pi t}{6} + \frac{1}{2}g(v|\tau) + \frac{1}{4\pi t} \left(E_2(\tau) - g_2(v|\tau) \right) + \mathcal{O}(e^{-2\pi t})$$

- $g(v|\tau)$ is the torus Green function;
- E_2 is the non-holo Eisenstein series;
- $g_2(v|\tau) = \int_{\Sigma_1} d^2z/\tau_2 g(v-z|\tau)g(z|\tau)$;
- Derived using the Laplace-eigenvalue equation for φ .
- The Laurent polynomial is of **finite degree** in the variable t
 - but it is not of finite degree in, say, $t + 1$

Non-separating node from a punctured surface

- **Standard construction of a surface near a non-separating node** (Fay 1973)
 - start from a genus-one surface with two punctures p_a, p_b
 - local coordinates z_a, z_b which vanish respectively at p_a, p_b
 - identify points $z_a z_b = t$ between two annuli $\mathcal{C}'_a = \mathcal{C}'_b$ and $\mathcal{C}''_a = \mathcal{C}''_b$



- in practice not easy to implement, unless one can define cycles \mathcal{C} naturally

Non-separating node from a punctured surface (cont'd)

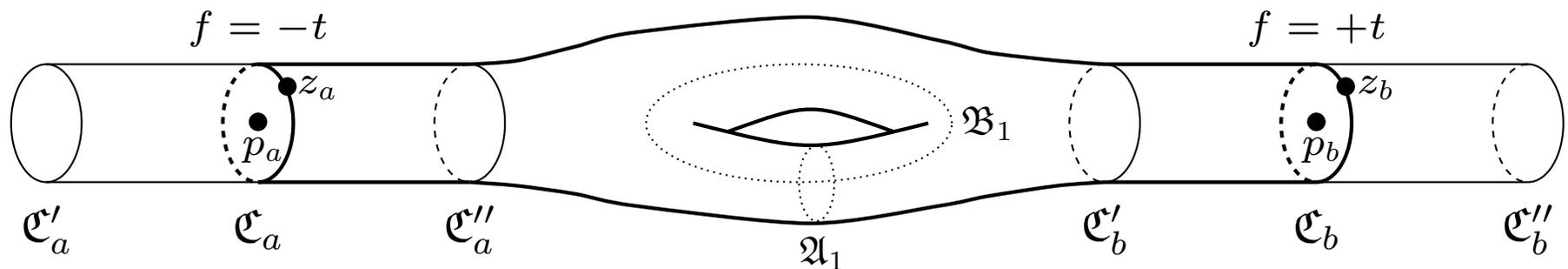
- Key is the existence of a real single-valued harmonic function $f(z)$ on Σ
 - such that in the degeneration limit

$$f(z) \rightarrow -\infty \text{ as } z \rightarrow p_a \text{ and } f(z) \rightarrow +\infty \text{ as } z \rightarrow p_b$$

- for large t cycles prescribed by $f(\mathcal{C}_a) \rightarrow -2\pi t$ and $f(\mathcal{C}_b) \rightarrow +2\pi t$

$$f(z) = -2\pi t + 4\pi \operatorname{Im} \int_{z_a}^z \left(\omega_2 - \operatorname{Im}(v) \omega_1 / \operatorname{Im}(\tau) \right)$$

- Funnel construction of the non-separating degeneration



- The cycle \mathcal{A}_2 is homologous to the cycles $\mathcal{C}_a, \mathcal{C}'_a, \mathcal{C}''_a$ and $\mathcal{C}_b, \mathcal{C}'_b, \mathcal{C}''_b$;
- The cycles are pairwise identified by $\mathcal{C}_a \approx \mathcal{C}_b, \mathcal{C}'_a \approx \mathcal{C}'_b$ and $\mathcal{C}''_a \approx \mathcal{C}''_b$;
- The points are pairwise identified by $z_a \approx z_b$;
- The cycle \mathcal{B}_2 may be chosen to be a simple curve connecting z_a to z_b .

Degeneration of higher string invariants

- Consider the full genus-two string amplitude

$$\mathcal{B}(s_{ij}|\Omega) = \int_{\Sigma^4} \frac{\mathcal{Y} \wedge \bar{\mathcal{Y}}}{(\det Y)^2} \exp \left\{ \sum_{i < j} s_{ij} \mathcal{G}(z_i, z_j|\Omega) \right\}$$

- Expanding in powers of s_{ij} and collecting all terms homogeneous of degree w
- gives modular graph functions of weight w for genus-two surfaces Σ

$$\mathcal{B}_w(s_{ij}|\Omega) = \int_{\Sigma^4} \frac{\mathcal{Y} \wedge \bar{\mathcal{Y}}}{(\det Y)^2} \left(\sum_{i < j} s_{ij} \mathcal{G}(z_i, z_j|\Omega) \right)^w$$

- **Theorem 4** Under non-separating degeneration of Σ (ED, Green, Pioline 2017)

- $\mathcal{B}_w(s_{ij}|\Omega)$ has a Laurent polynomial of degree (w, w) in t

$$\mathcal{B}_w(s_{ij}|\Omega) = \sum_{k=-w}^w \mathcal{B}_w^{(k)}(s_{ij}|v, \tau) t^k + \mathcal{O}(e^{-2\pi t})$$

- $\mathcal{B}_w^{(k)}(s_{ij}|v, \tau)$ are invariant under $SL(2, \mathbb{Z}) \subset Sp(4, \mathbb{Z})$

$$\mathcal{B}_w^{(k)} \left(s_{ij} \left| \frac{v}{c\tau + d}, \frac{a\tau + b}{c\tau + d} \right. \right) = \mathcal{B}_w^{(k)}(s_{ij}|v, \rho)$$

- Note
 - similarity with the Laurent polynomial expansion for genus one;
 - $\mathcal{B}_w^{(k)}(s_{ij}|v, \tau)$ generalize Jacobi forms to a non-holomorphic setting;
 - combine genus-one modular graph functions and *elliptic polylogarithms*.

Ingredients in proof of Theorem 4

- **Uniform asymptotics of the Green function**

$$\mathcal{G}(x, y|\Omega) = \frac{\pi t}{12} + g(x - y|\tau) - \frac{f(x)f(y)}{4\pi t} + \cdots + \mathcal{O}(e^{-2\pi t})$$

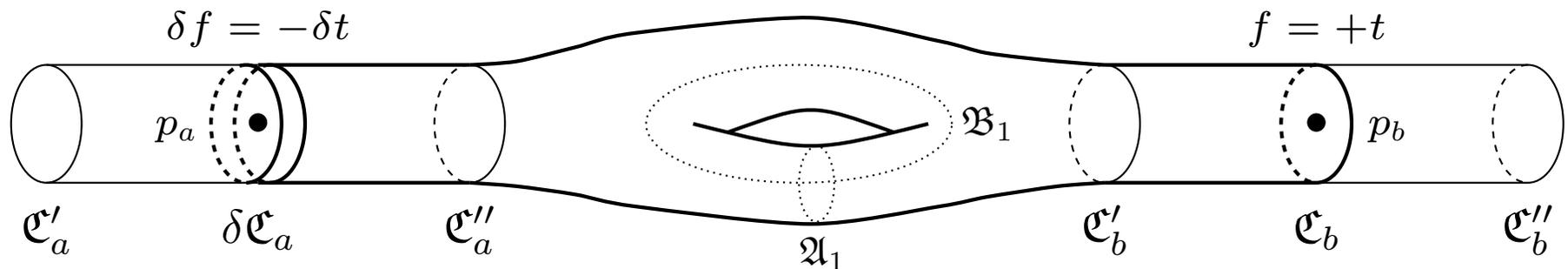
– see DGP 2017 for the complete expression (improving on Wentworth 1991)

- **Measure**

$$4\pi^2 \mathcal{Y} = dz_1 \wedge dz_2 \wedge dz_3 \wedge dz_4 \sum_{i < j} s_{ij} \partial_{z_i} f(z_i) \partial_{z_j} f(z_j)$$

- **The cycles $\mathcal{E}_a, \mathcal{E}_b$ have exponentially vanishing coordinate radius**

- $z \in \mathcal{E}_a$ satisfies $|z - p_a|^2 \approx e^{-2\pi t}$
- Power dependence in t arise from poles at p_a, p_b in the integrand
- Extract variation in t -dependence (cfr RG with cut-off t)



Summary and outlook

- **Low energy expansion of string theory has revealed a rich structure of**
 - Modular graph functions for genus-one Riemann surfaces;
 - Kawazumi-Zhang and higher string invariants for genus-two surfaces.
- **Asymptotics for higher string invariants**
 - Remarkable structure in Laurent polynomial for non-sep degeneration
 - Tested versus maximal degeneration limit (ED, Green, Pioline 2018)
 - Similar result for separating degeneration (ED, Green, Pioline 2018)
 - Crucial for the discovery of relations between modular graph functions
- **Differential equations for higher string invariants**
 - In progress (ED, Green, Pioline)
 - Asymptotics is expected to be a key guide