

Particle life time in

De Sitter Space

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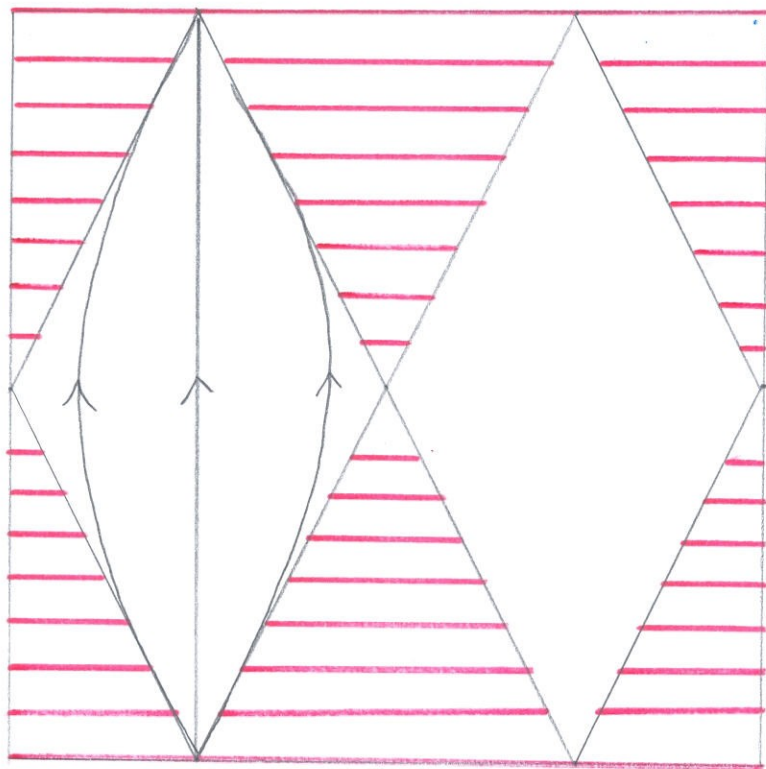
CMP and in preparation.

DE SITTER

Sphere $-T^2 + X^2 = R^2$

No global Killing vector

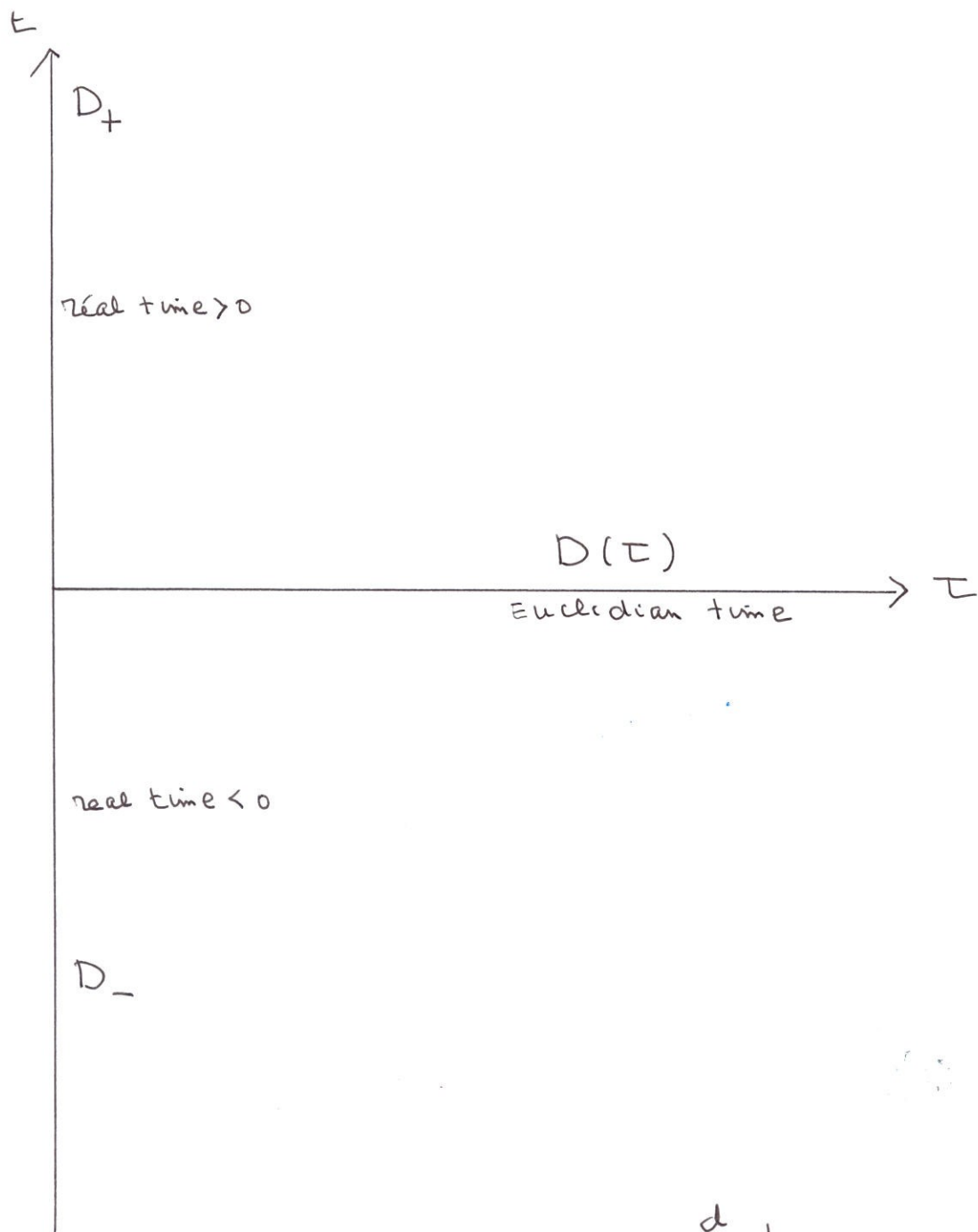
$$\begin{pmatrix} 0 \\ \cos \theta \\ \sin \theta \end{pmatrix} \longrightarrow \begin{pmatrix} \sinh t & \cos \theta \\ \cosh t & \cos \theta \\ \sin \theta \end{pmatrix}$$



Wightman Schwinger functions
Minkowski space

$$\square_x D(x, x') = 0$$

$$D(\sqrt{(x-x')^2})$$

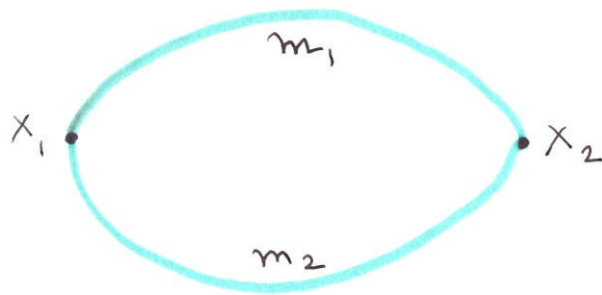


$$D_m(\tau) = (m\tau)^{-\frac{d}{2}+1} K_{\frac{d}{2}-1}(m\tau)$$

Källén - Lehmann

$$S^+ S = 1, \quad S = 1 + iT$$

$$i(T - T^+) = -T^+ T$$



$$D_{m_1} D_{m_2} = \int \rho(m) D_m dm$$

Fermi Golden rule:

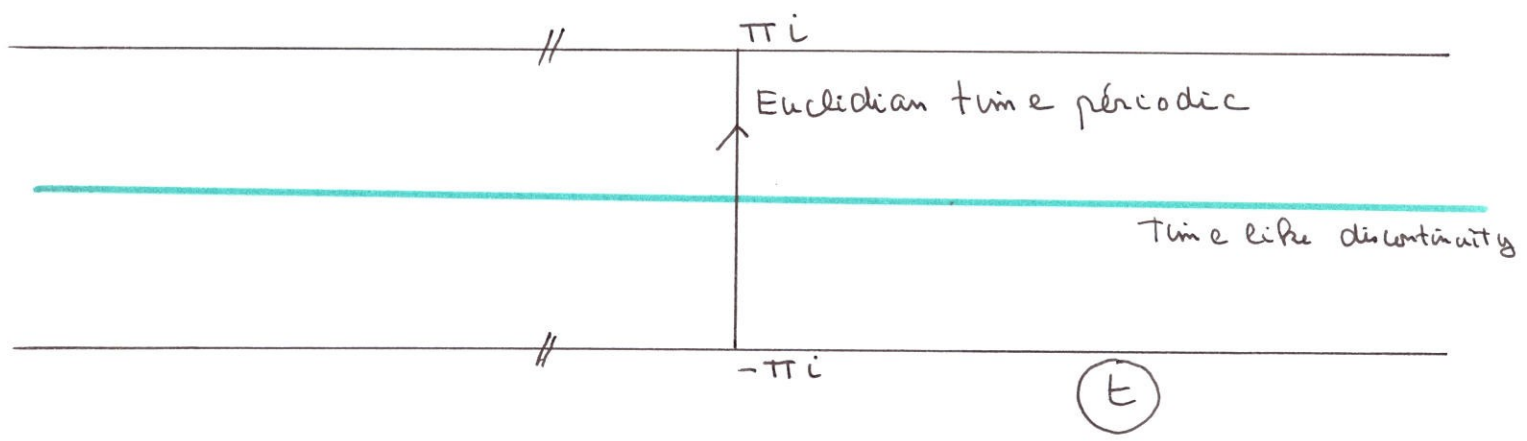
$$\rho = \int \delta_+(p^2 - m_1^2) \delta_+(q^2 - m_2^2) \delta^{(d)}(p+q-M)$$

$$= \left[\frac{\prod_{4\epsilon} (m \pm m_1 \pm m_2)}{m^2} \right]^{\frac{d-3}{2}} \Theta(m - m_1 - m_2)$$

De Sitter ?

$$D_n(t)$$

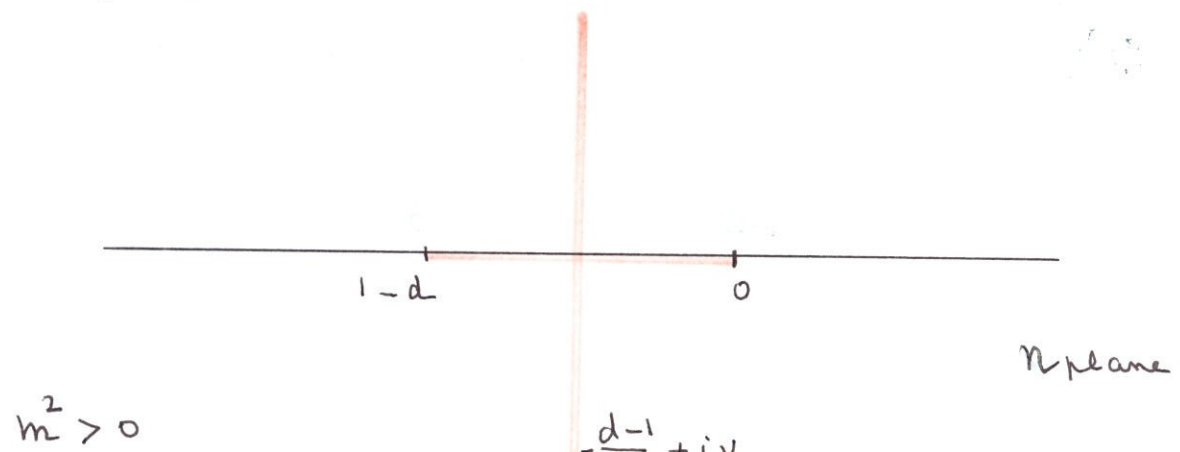
$$(x - x')^2 = \text{ch } t - 1$$



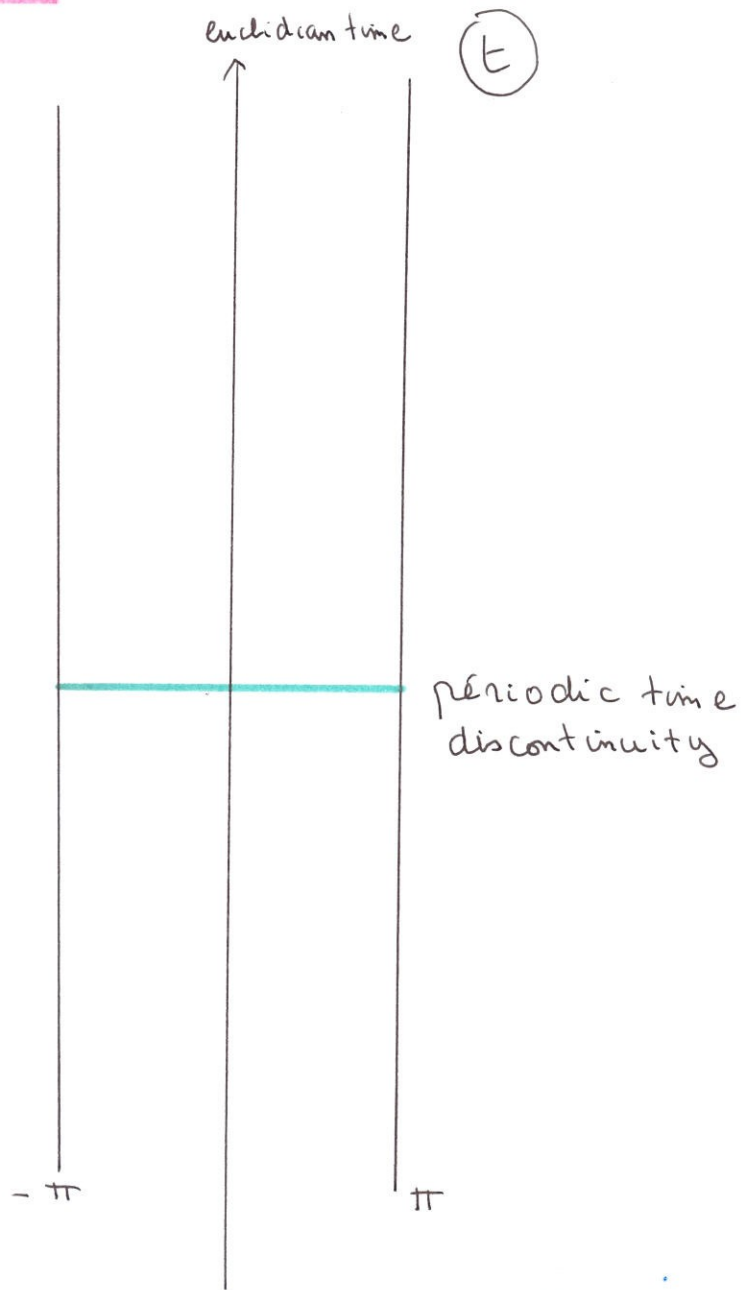
$$D = P_n^{(\frac{d-1}{2})}(-\text{ch } t)$$

Periodicity in Euclidian time has thermic interpretation:

KMS condition.



AdS



$$(x - x')^2 = \cos t - 1$$

$$D_n = Q_n^{(\frac{d-1}{2})} (\cos t)$$

Results

$$\lambda = \frac{d-1}{2}$$

$$\nu = n + \lambda, \text{ Imaginary}$$

$$W_\nu = \Gamma(\lambda + \nu) \Gamma(\lambda - \nu) F\left(\lambda + \nu, \lambda - \nu, \lambda + \frac{1}{2}, \frac{1-z}{2}\right)$$

$$W_{\nu_1} W_{\nu_2} = \int_{i\mathbb{R}} d\nu \nu \rho \sin(\pi\nu) \phi(\nu) \frac{W_\nu}{\Gamma(\lambda + \nu) \Gamma(\lambda - \nu)}$$

$$\phi(\nu) = \prod_{\text{8 factors}} \Gamma(\lambda \pm \nu_1 \pm \nu_2 \pm \nu)$$

Formula giving the
Källén Lehmann weight

Relation with Clebsch Gordan Coefficients

Gegenbauer Polynomials are bi-invariant functions on the group $SO(d+1)$
 $(SO(d) \backslash SO(d+1) / SO(d))$ with respect to $SO(d)$

$$\binom{\lambda}{n} = \frac{\Gamma(n+2\lambda)}{\Gamma(n+1)\Gamma(2\lambda)} F_{\nu}$$

$$\nu = n + \lambda \quad \text{and} \quad n \in \mathbb{N}$$

$$\int_{-1}^1 d\mu(x) \binom{\lambda}{n_1}(x) \binom{\lambda}{n_2}(x) \binom{\lambda}{n_3}(x) = C_{n_1 n_2 n_3}^{\lambda}$$

$C_{n_1 n_2 n_3}^{\lambda} \neq 0$ only if

$n_1 + n_2 + n_3$ is even and n_1, n_2, n_3 form a triangle. Szegő Orthogonal Polynomials.

$C_{n_1 n_2 n_3}^{\lambda}$ is a Clebsch Gordan coefficient
and K.L. is an extension of the results
to the $SO(d,1)$ Unitary series

Intriguing Observation

(G. E. Volovick)

De Sitter violates the Lorentz invariance mass inequality:

$$M \longmapsto m_1 + m_2$$

$$M \geq m_1 + m_2$$

In the limit where the violation is small:

$$\Delta M \ll M$$

$$w \propto \exp\left(-\frac{\Delta M}{T}\right)$$

Thermal activation process with

$$T = \frac{\hbar H}{\pi} = 2 T^{dS}$$

Idea for derivation

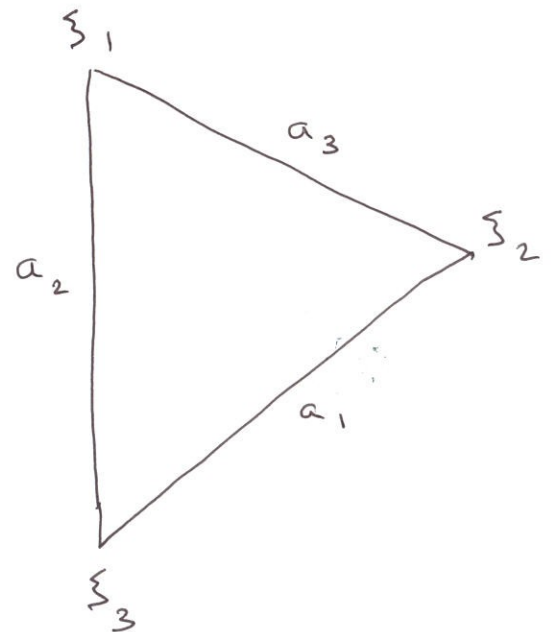
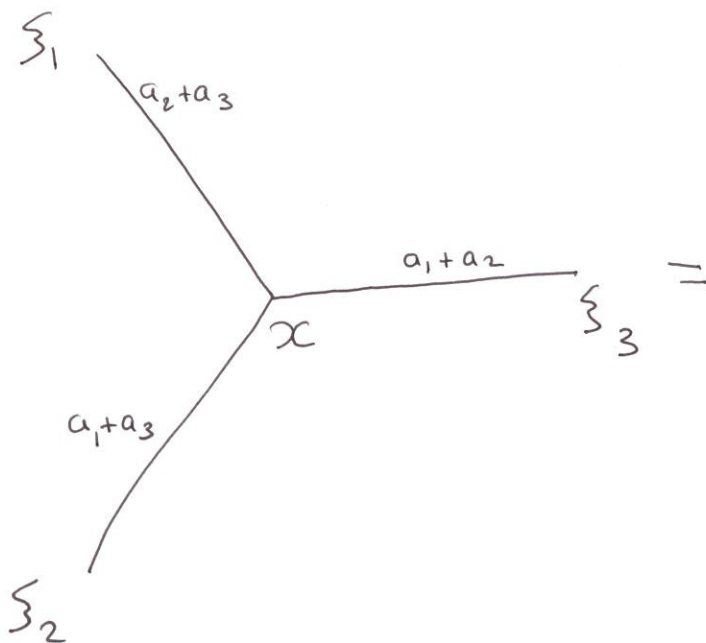
$$C_n^{(\lambda)}(x) = \int d\xi, (x \cdot \xi)^{-n}$$

ξ variable on the cone. $\xi^2 = 0$

$$\int_H C_{n_1} C_{n_2} C_{n_3} dx$$

$$= \int_{S^{d-1}} d\xi_1 d\xi_2 d\xi_3 \int_H dx (\xi_1 x)^{-n_1} (\xi_2 x)^{-n_2} (\xi_3 x)^{-n_3}$$

Integral over x

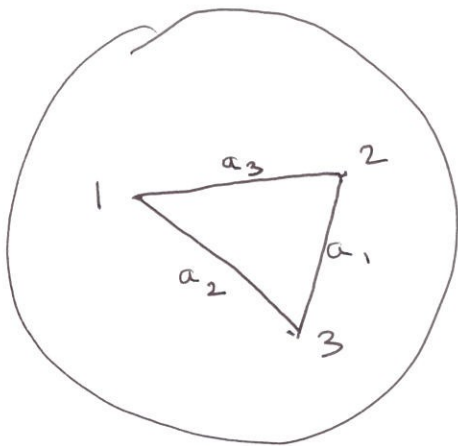


$$C(x) = \Gamma(-a-\lambda) \frac{\Gamma(-a_1) \Gamma(-a_2) \Gamma(-a_3)}{\Gamma(-a_2-a_3) \Gamma(-a_1-a_3) \Gamma(-a_1-a_2)}$$

$$\int (\xi_1, \xi_2)^{a_3} (\xi_2, \xi_3)^{a_1} (\xi_1, \xi_3)^{a_2}$$

$$= \int P(\tau_{12}, \tau_{23}, \tau_{13}) \tau_{12}^{a_3} \tau_{23}^{a_1} \tau_{13}^{a_2}$$

Moments of the probability that
 3 random points on the sphere
 S_{d-1} constitute a triangle
 with sides $\tau_{12}, \tau_{23}, \tau_{13}$.



$$= c(\lambda) \frac{\Gamma(a+2\lambda) \Gamma(\lambda+a_1) \Gamma(\lambda+a_2) \Gamma(\lambda+a_3)}{\Gamma(2\lambda+a_2+a_3) \Gamma(2\lambda+a_1+a_3) \Gamma(2\lambda+a_1+a_2)}$$

Conclusions

- De Sitter is very interesting to study for numerous reasons: Inflation, acceleration, ...
- No good notion of wave packets
acceleration \rightarrow + washes away all the structure, reason why it is interesting to study inflation.
- No killing vector makes Euclidean continuation difficult. How to use complex sphere to make predictions?
- Many connections with group theory
Källén-Lehmann weight = Clebsch-Gordan coefficients, particles organize into representations of $SO(d, 1)$ group.
- Is there any analogous of dS-gauge correspondence?