

# Resurgence

formal divergent power series  $\hat{\Phi} := \sum_{n=0}^{\infty} a_n \bar{z}^{n-1} \in \bar{z}^{-1} \mathbb{C}[[\bar{z}]]_1$

↳ 0-radius of conv.  
 "an grow faster than  $A^n$ "

$a_n \sim \underline{n!} A^n$       1-Genrey

- perturbation theory
- asymptotics  $\Gamma(z) = \int_0^{\infty} t^{z-1} e^{-t} dt \sim e^{-z} z^{-1/2} \sqrt{\pi} \left( 1 + \frac{1}{12z} + \frac{1}{288z^2} + \dots \right)$
- Solving im. sing. ODEs  $\hat{\Phi}' = -\hat{\Phi} + \frac{1}{z}$

GOAL: find non-perturbative contributions  
 "Resum" the formal solution to get an analytic one

① Borel transform "remove divergence"

formal  $\mathcal{B}_\xi [z^{n-1}] \rightsquigarrow \frac{\xi^n}{n!}$  + linearity

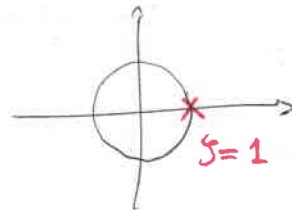
$$\mathcal{B}_\xi : \sum_{n=0}^{\infty} a_n \bar{z}^{n-1} \xrightarrow{\mathcal{B}_\xi} \sum_{n=0}^{\infty} a_n \frac{\xi^n}{n!} =: \hat{\phi}$$

$\bar{z}^{-1} \mathbb{C}[[\bar{z}]]_1$        $\mathbb{C}\{\xi\}$       finite radius of conv.

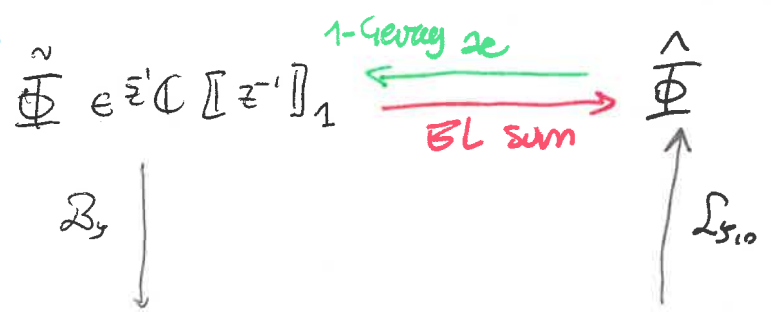
e.g.  $\sum_{n=0}^{\infty} n! \bar{z}^{n-1} \xrightarrow{\mathcal{B}_\xi} \sum_{n=0}^{\infty} n! \frac{\xi^n}{n!} = \sum_{n=0}^{\infty} \xi^n \quad |\xi| < 1$

- study analytic continuation of  $\hat{\phi}$

e.g.  $\hat{\phi} = \frac{1}{1-\xi}$  sum of  $\hat{\phi}$



$z$ -frequency



$s$ -position  
(Borel plane)



to go back to  $z$ -plane we want inverse of  $B_s$  but ANALYTIC

② Laplace

$$L_{s,0} f := \int_0^{+\infty} e^{-zs} f ds \quad \begin{array}{l} f \text{ integrable} \\ |f(s)| \leq A e^{c_0|s|} \quad s \gg 1 \end{array}$$

$$L_{s,0} \left[ \frac{s^n}{n!} \right] = z^{-n-1} \quad (\text{Borel is formal inverse of Laplace})$$

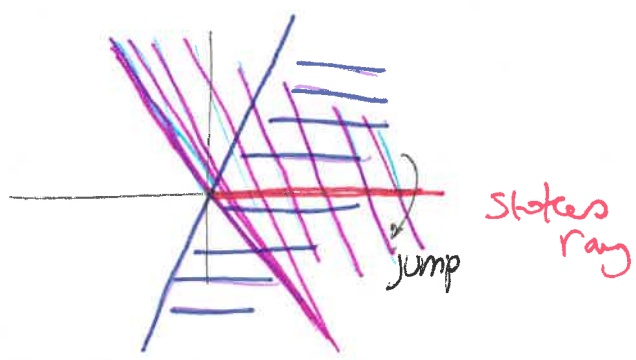
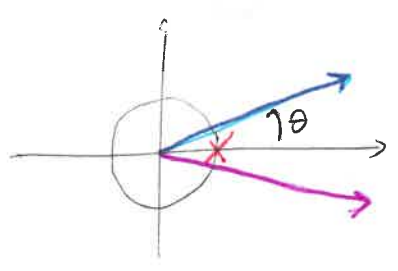
Thm:  $L_{s,0} f$  is holom. in  $\{ \text{Re } z > c_0 \}$  [draw diagram]

Def: Borel-Laplace sum

• change direction

$$L_{s,0}^\theta f := \int_0^{+\infty} e^{-ze^{i\theta}s} f ds \quad \text{holom. in } \{ \text{Re } ze^{-i\theta} > c_0 \}$$

e.g.



$$\hat{\Phi}_+ = L_{s,0}^\theta \frac{1}{1-s}$$

$$\hat{\Phi}_- = L_{s,0}^{-\theta} \frac{1}{1-s} \quad \underline{\underline{\text{two analytic functions}}}$$

Having singularity at  $s=1 \Rightarrow \hat{\Phi}_+ \neq \hat{\Phi}_-$  where both defined

$$\hat{\Phi}_+ - \hat{\Phi}_- = \int_{\text{loop}} e^{-zs} \frac{1}{1-s} ds = e^{-z} \cdot 2\pi i \quad \underline{\underline{\text{jump!}}} \quad \rightarrow \text{STOKES PHENOMENA}$$

Remark: asymptotically  $\hat{\Phi}_+, \hat{\Phi}_-$  have same behavior  $\text{Re } z \rightarrow +\infty$ .

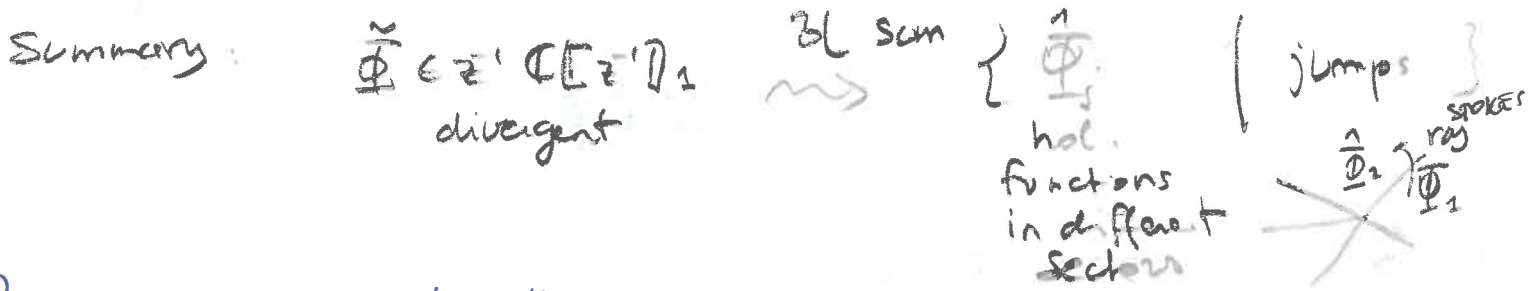
"Jump" is non-perturbative, subleading contribution

Thm (inverse of B-L sum)

$\hat{\Phi}$  is uniform Gevrey asympt. to  $\tilde{\Phi}$  in  $\text{Re } z > c \quad c > c_0$ .

$$|\hat{\Phi}(z) - \sum_{n=0}^N a_n z^{-n-1}| \leq C A^N N! |z|^{-N-2}$$

Remark: if you are familiar with ODEs, STOKES PHENOMENA is common when there are irregular singularities. Monodromy data of your system corresponds to matrices of STOKES constants.



Resurgence computes "jumps" from the study of  $\hat{\phi}$  in Borel plane  
 $\hookrightarrow$  singularities  
 $\hookrightarrow$  analytic continuation

eg.  $\hat{\phi} = \frac{1}{1-y}$   $\hookrightarrow$  singularity  
 Jump:  $e^{-z \cdot 1} \cdot \frac{2\pi i}{\text{Res } \hat{\phi}}$   $\checkmark$  (analytic continuation) monodromy

Def:  $\tilde{\Phi} \in z^{-1} \mathbb{C}[[z^{-1}]]_1$  is resurgent series if  $\hat{\phi} := \mathcal{B} \tilde{\Phi}$  has endless analytic continuation away from  $\Omega \subseteq \mathbb{C}$ .

$\forall L > 0 \exists \Omega_L \subseteq \mathbb{C}$  finite s.t.  $\tilde{\phi}$  can be analytically continued along any paths starting near  $y=0$  and of length  $\leq L$  in  $\mathbb{C} \setminus \Omega_L$ .

$\Omega$  can be countable,

NOT RESURGENT:  $\sum_{n=0}^{\infty} a_n n! z^{-n-1} \quad a_n = \begin{cases} -1 & n=2^k \\ 1 & \text{otherwise} \end{cases} \quad \sum_{k=0}^{\infty} y^{2^k} \quad \gamma_n = \frac{2\pi i}{2^m} \quad |s| < 1 \quad 3/$

# Simple resurgent function

$\Omega$  set of singularities,  $w \in \Omega$

$$\hat{\phi}(s+w) = \frac{C_w}{2\pi i s} + \frac{S_w}{2\pi i} \log(s) \tilde{\phi}_w(s) + h.f. \quad \text{near } s=0$$

$$C_w, S_w \in \mathbb{C}, \quad \tilde{\phi}_w(s) \in \mathbb{C}\langle s \rangle$$

— what is the relevant information?

e.g. ①  $\hat{\Phi} =$  asymptotics of  $\frac{\Gamma(z) e^z z^{-z+\frac{1}{2}}}{\sqrt{\pi}}$  as  $\text{Re } z \rightarrow \infty$

$$\Omega = 2\pi i \mathbb{Z}, \quad \hat{\Phi}(s+2\pi i) = \frac{1}{2\pi i} \log(s) \hat{\Phi}_{2\pi i}(s) + h.f.$$

same as  $\beta \hat{\Phi}$

(see Maxim's talk from this summer)

$$\textcircled{2} \quad \left[ \partial_z^2 - 1 + \frac{1}{z} \partial_z - \frac{1}{9z^2} \right] \phi = 0$$

$$\hookrightarrow \text{formal solutions } \tilde{\Phi}_{\pm 1}(z) = e^{\pm z} z^{-3/2} \sum_{n=0}^{\infty} \frac{\left(\frac{1}{6}\right)_n \left(\frac{5}{6}\right)_n}{2^n n!} (\mp z)^{-n-1}$$

$$\hat{\Phi}_2(s) = {}_2F_1(a, b, c; -\frac{s}{2}) \Rightarrow \hat{\Phi}_2(s+2) = \frac{1}{2\pi i} \log(s) \tilde{\Phi}_{-1}(s) + h.f.$$

$$\hat{\Phi}_{-2}(s) = {}_2F_1(a, b, c; \frac{s}{2}) \Rightarrow \hat{\Phi}_{-2}(s-2) = \frac{(-1)}{2\pi i} \log(s) \tilde{\Phi}_1(s) + h.f.$$

new Series!

idea of resurgence: functions near singularities reproduce themselves, "they resurge near their singularities" like in ①

Écalle introduced Alien calculus, and formalism of singularities to compute  $(C_w, S_w, \tilde{\phi}_w)$

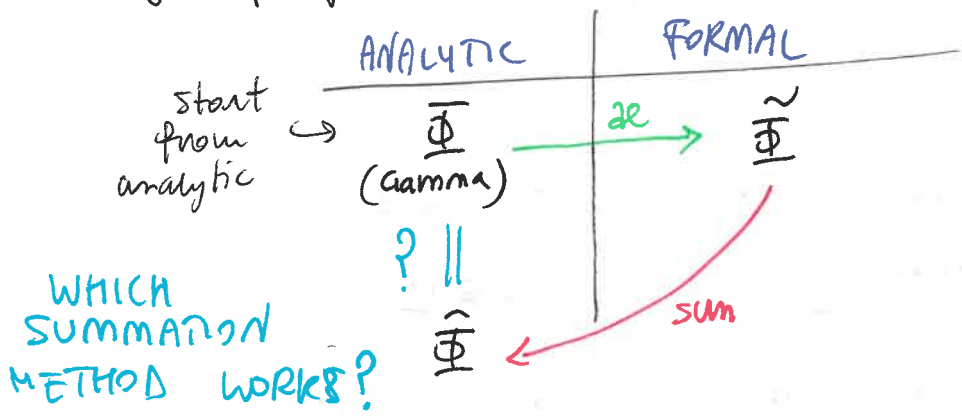
We refer to  $(C_w, S_w, \tilde{\phi}_w \mid w \in \Omega)$  as resurgent ~~data~~ structure

How do you check something is resurgent?

- knowing a close formula for the coefficients
  - ↳ know the sum (like hypergeometric)
  - ↳ HADAMARD PRODUCT to compute the sum of  $\tilde{\Phi}$
- knowing only  $a_1, a_2, \dots, a_N$ , no close formula
  - ↳ Borel + Padé numerically, at least to see the position of the singularities.
- the formal series comes from solving ODEs, PDEs, difference eq. you might know it's resurgent because of the problem you're solving (Borel, Laplace behaves well w/  $\partial_z, T_c$ , products, ...)

Back to summability

- Borel-Laplace  $\tilde{\Phi} \in \mathcal{E}' \subset \mathbb{C}[[z^{-1}]]_1 \rightarrow \hat{\Phi}$  analytic
- change perspective

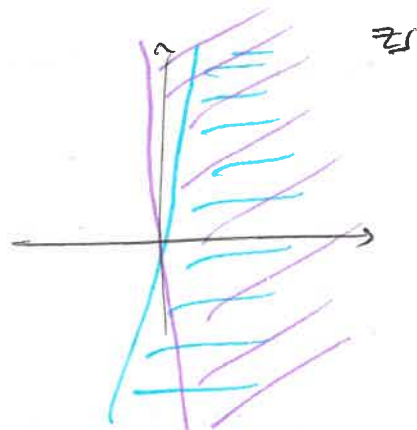
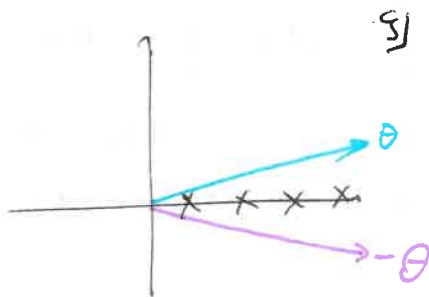


(for  $\Gamma(z)$ , Borel-Laplace works)

THEM (Nevanlinna) If  $\Phi$  is uniform Gevrey, asymptotics ~~sum~~ in a sector of opening angle  $\pi$ , to  $\tilde{\Phi}$ , then  $\Phi = L \circ B \tilde{\Phi}$ .

# Def (median renormalization)

$$S^{\text{med}} = \frac{L^\theta + L^{-\theta}}{2}$$



think of  $S^{\text{med}}$  as ~~an approximation~~ an approximation of  $L$  along ray of singularities.

$S^{\text{med}}$  is defined for  $\text{Re } z > 0$ .

rewrite

$$S^{\text{med}} = \begin{cases} L^\theta - \frac{\text{disc}}{2} \\ L^{-\theta} + \frac{\text{disc}}{2} \end{cases}$$

$$= \begin{cases} L^{\pi/2} - \frac{\text{disc}}{2} \\ L^{-\pi/2} + \frac{\text{disc}}{2} \end{cases}$$

$$\text{disc} := L^\theta - L^{-\theta}$$

deform the contour below, and collect all the singularities

## Conj [Costin-Garnufalidis]

If  $\Phi$  is WRT invariant of closed 3-manif. / Kashaev inv. for knots

then  $\Phi(z) = S^{\text{med}} \tilde{\Phi}(-\frac{1}{z})$  at  $z \in \frac{\mathbb{Q}}{2\pi i}$

Thm [Costin-Garnufalidis] true for trefoil knot.

Remark:  $\Phi(z) \neq L^{\pi/2} \Phi$   $\text{Im } z > 0$

$$\Phi(z) = L^{\pi/2} \Phi - \left| \frac{\text{disc}}{2} \right| \text{ at } z \in \frac{\mathbb{Q}}{2\pi i} // \text{rationals}$$

corrections coming from resurgent structure

Here there is something interesting going on, I think it's related with QUANTUM MODULARITY...